

Measurements of Quantum Interference of Two Identical Particles with respect to the Event Plane in Au+Au Collisions at $\sqrt{s_{NN}}=200$ GeV at RHIC-PHENIX

(RHIC-PHENIX実験 200 GeV 金+金衝突における
同種2粒子を用いた量子力学的干渉効果の反応平面依存性の測定)

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博士論文公開発表会(本審査) **10/24/2013**

outline

■ Introduction

- ✧ Hanbury-Brown&Twiss Interferometry (HBT)
- ✧ Motivation

■ Experiment/Analysis

- ✧ PHENIX Detectors
- ✧ Data selection
- ✧ Analysis Method for HBT

■ Results/Discussion

- ✧ HBT measurement with respect to the event plane
- ✧ Blast-wave model
- ✧ Monte-Carlo simulation

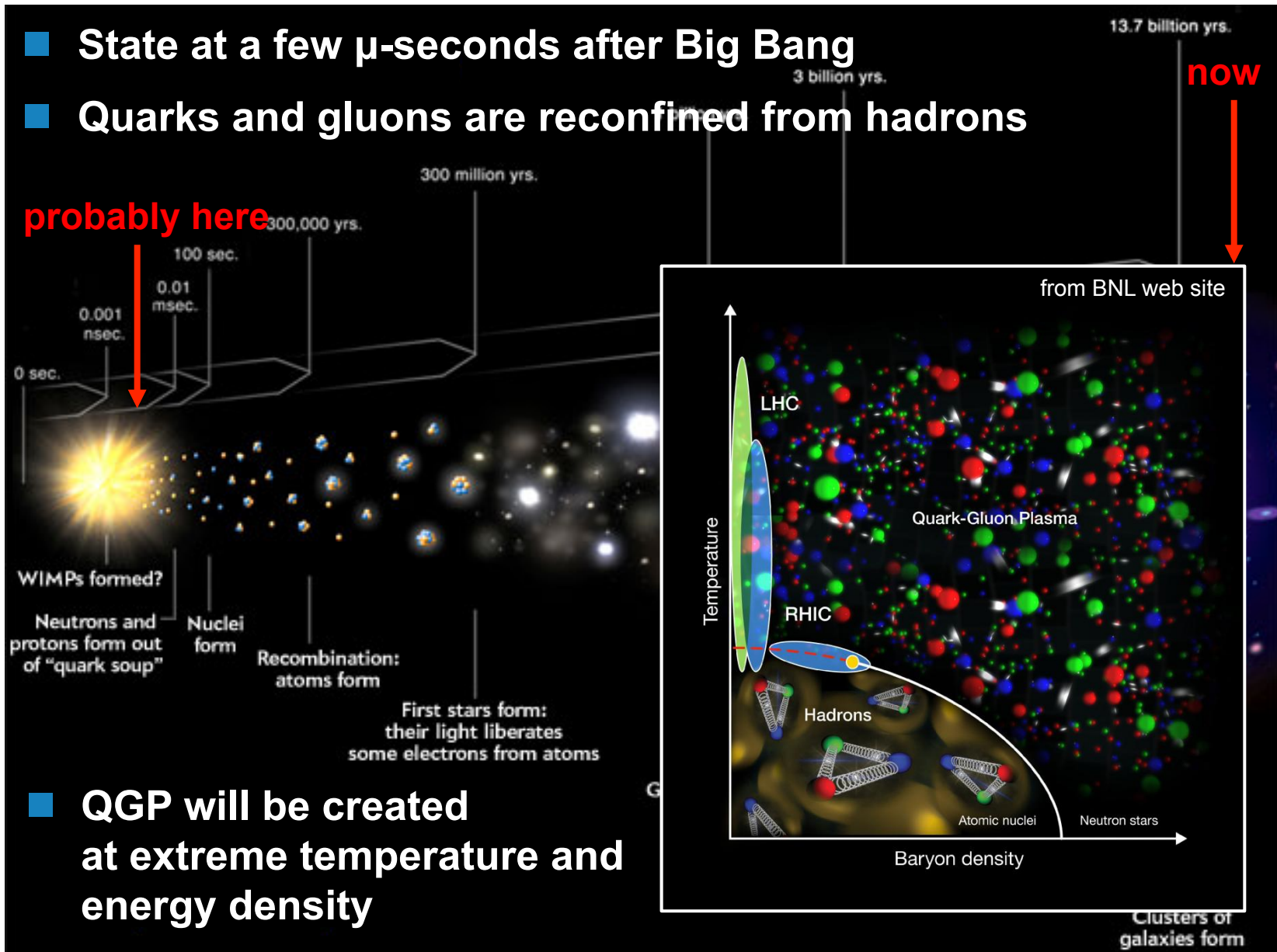
■ Summary/Conclusions

Introduction

Quark Gluon Plasma (QGP)

<http://www.scientificamerican.com/>

- State at a few μ -seconds after Big Bang
- Quarks and gluons are reconfined from hadrons

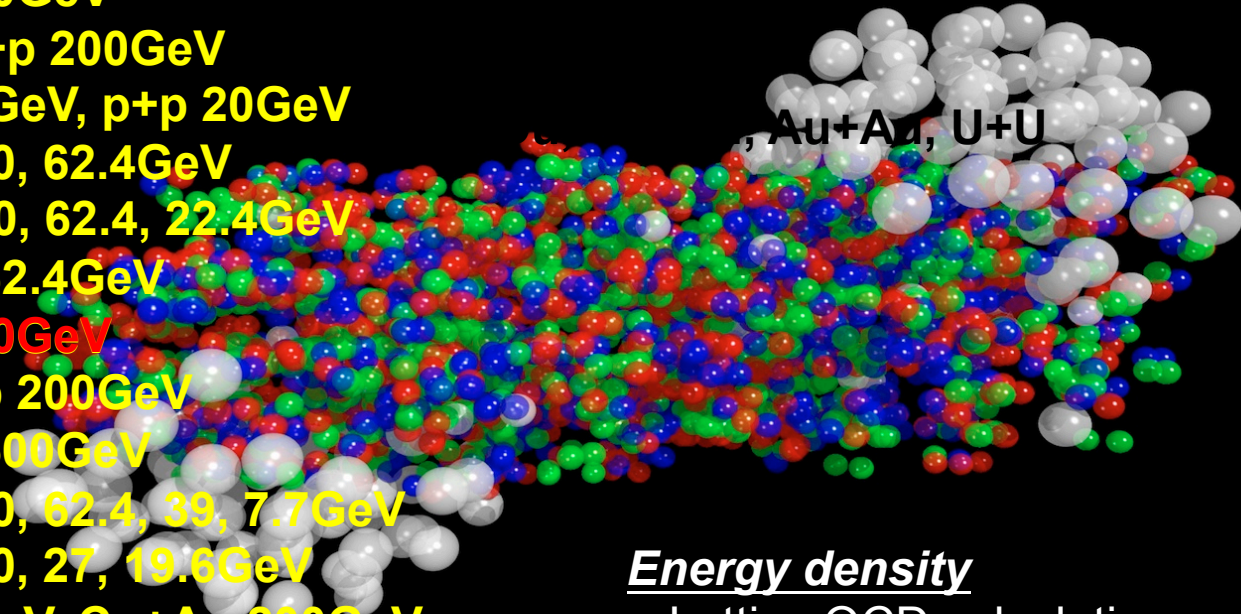


- QGP will be created at extreme temperature and energy density

Relativistic Heavy Ion Collisions

- Relativistic Heavy Ion Collider is an unique tool to create QGP.

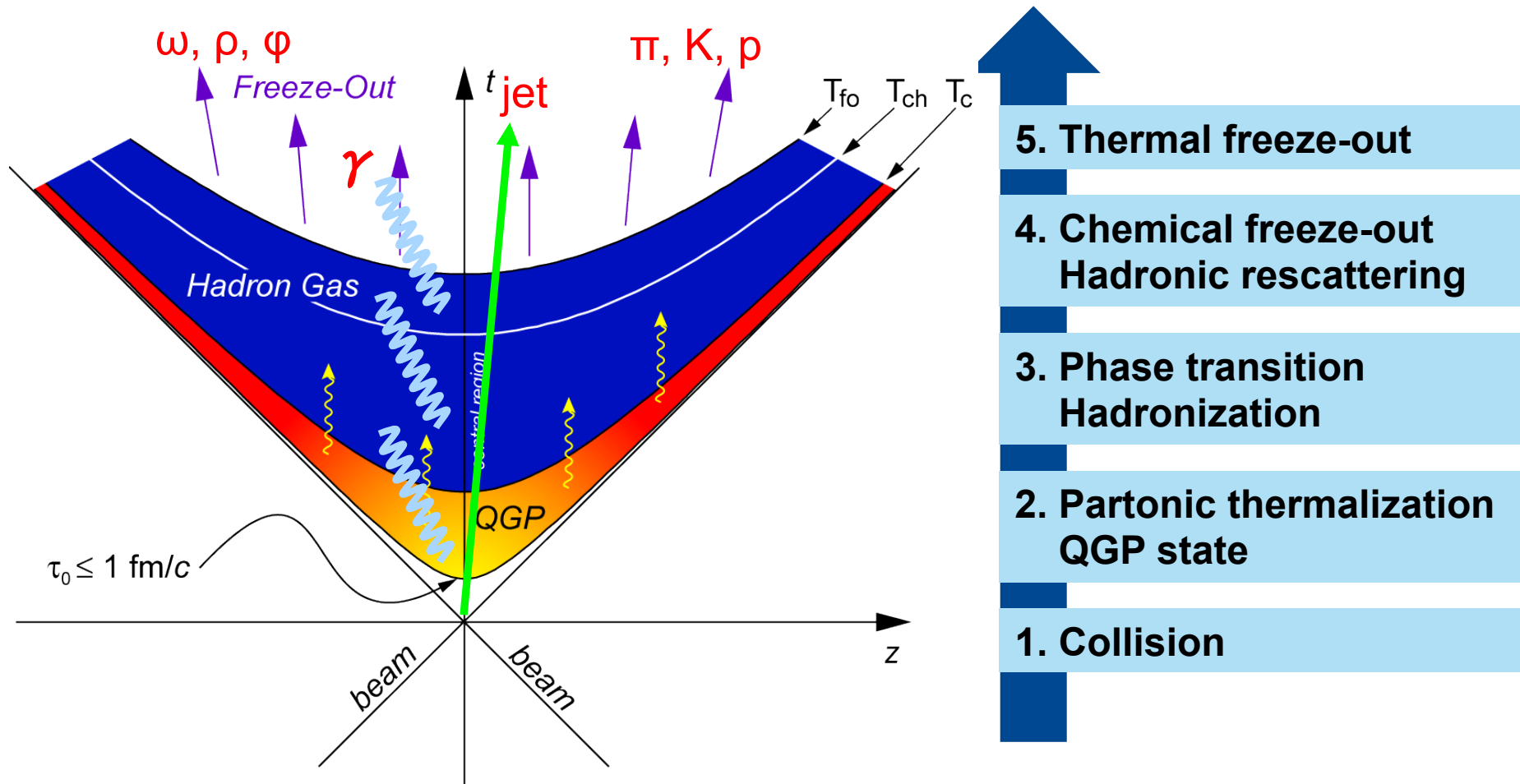
Year	Species/Energy
2001	Au+Au 130GeV
2002	Au+Au, p+p 200GeV
2003	d+Au 200GeV, p+p 20GeV
2004	Au+Au 200, 62.4GeV
2005	Cu+Cu 200, 62.4, 22.4GeV
2006	p+p 200, 62.4GeV
2007	Au+Au 200GeV
2008	d+Au, p+p 200GeV
2009	p+p 200, 500GeV
2010	Au+Au 200, 62.4, 39, 7.7GeV
2011	Au+Au 200, 27, 19.6GeV
2012	U+U 193GeV, Cu+Au 200GeV



Energy density

- Lattice QCD calculation
 - $T_c \sim 170 \text{ MeV}$
 - $\epsilon_c \sim 1 \text{ GeV/fm}^3$
- Au+Au 200GeV @RHIC
 - $\epsilon_{Bj} \sim 5 \text{ GeV/fm}^3 > \epsilon_c$

Space-Time Evolution

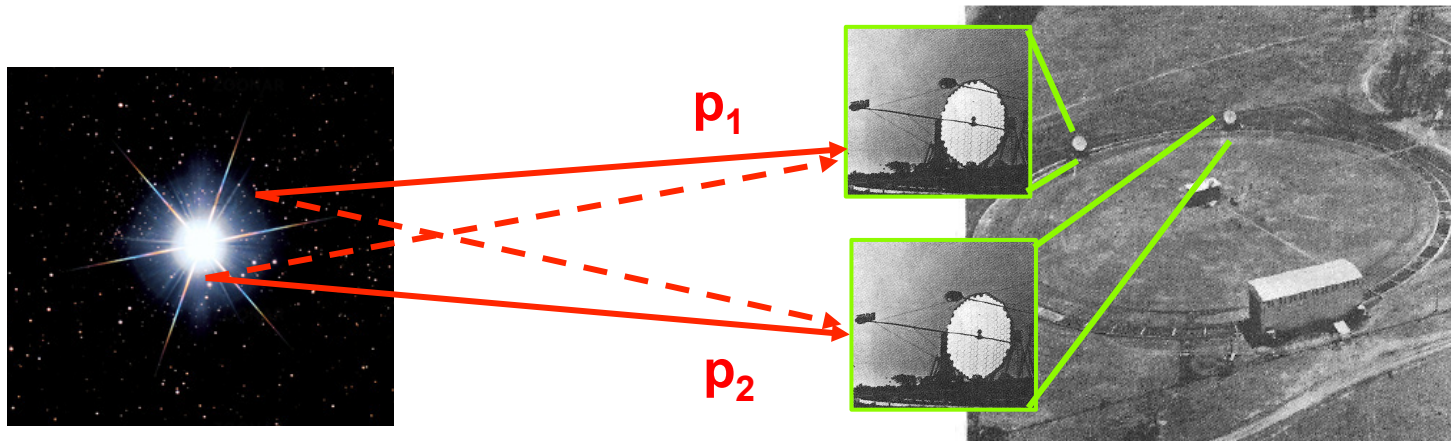


How fast the system evolves?
 How much the system size and shape?
 Study a detailed space-time picture of evolution of the QGP

HBT Interferometry

■ R. Hanbury Brown and R. Twiss

- ✧ In 1956, the angular diameter of Sirius was measured.
- ✧ HBT effect is *quantum interference* between two identical particles.



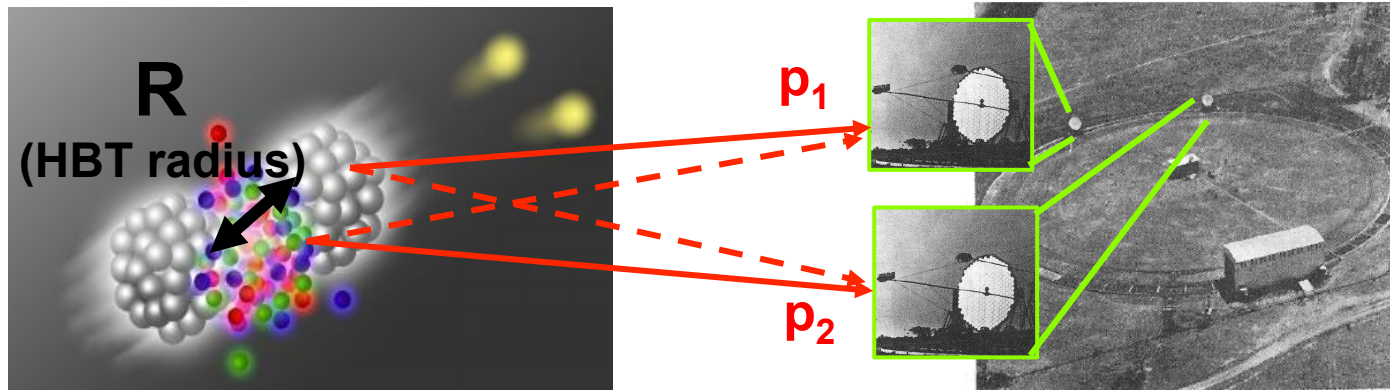
G. Goldhaber, Proc. Int. Workshop on Correlations and Multiparticle production(1991)

wave function of 2 bosons(fermions) : $\Psi_{12} = \frac{1}{\sqrt{2}}[\Psi_1(p_1)\Psi_2(p_2) \pm \Psi_2(p_1)\Psi_1(p_2)]$

Correlation function C_2 : Ratio of probabilities to detect 2 particles and 1 particle

$$C_2 = \frac{P_{12}}{P_1 P_2}$$

HBT Interferometry



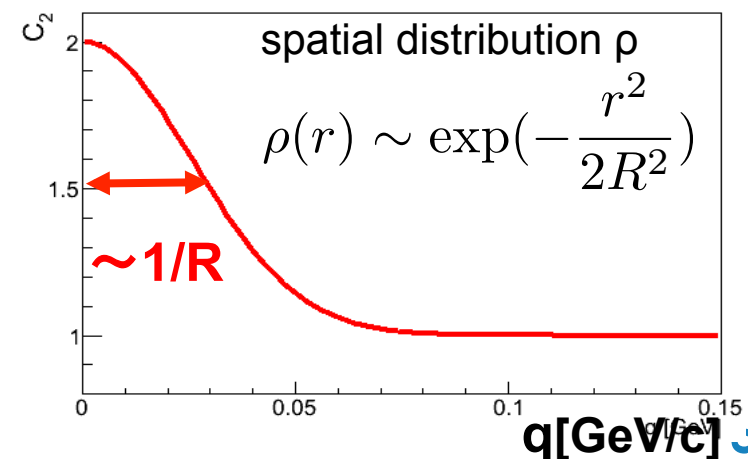
(assuming plane wave)

$$\begin{aligned}
 P_{12} &\propto [e^{-ip_1 \cdot (x_A - x_1)} e^{-ip_2 \cdot (x_B - x_2)} + e^{-ip_2 \cdot (x_A - x_2)} e^{-ip_1 \cdot (x_B - x_1)}]^2 \\
 &= [2 \pm e^{-ix_1(p_2 - p_1)} e^{ix_2(p_2 - p_1)} \pm e^{ix_1(p_2 - p_1)} e^{-ix_2(p_2 - p_1)}] \\
 &\propto 1 \pm \cos[(x_2 - x_1) \cdot (p_2 - p_1)]
 \end{aligned}$$

→ **HBT correlation**

$$\begin{aligned}
 C_2 &= \frac{P_{12}}{P_1 P_2} = \frac{\int dx_1 dx_2 \rho(x_1) \rho(x_2) |\Psi_{12}|^2}{\int dx_1 \rho(x_1) |\Psi_1|^2 \int dx_2 \rho(x_2) |\Psi_2|^2} \\
 &= 1 \pm \left| \int dx \rho(x) e^{ix(p_2 - p_1)} \right|^2 \\
 &= 1 + |\tilde{\rho}(q)|^2 \\
 &\approx 1 + \exp(-R^2 q^2)
 \end{aligned}$$

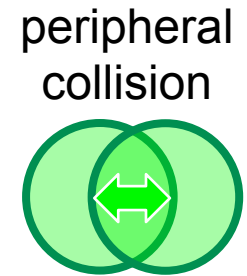
$\tilde{\rho}(q)$: Fourier transform of $\rho(r)$, $q = p_2 - p_1$



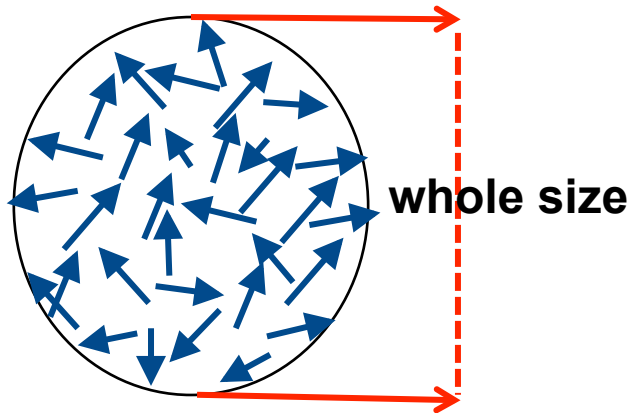
What do HBT radii depend on ?

- Centrality dependence
- Transverse momentum (k_T) dependence

- ✧ static source: whole source is measured.
- ✧ **expanding source**: Size of “**emission region**” is measured.
Collective expansion makes “x-p correlation”.

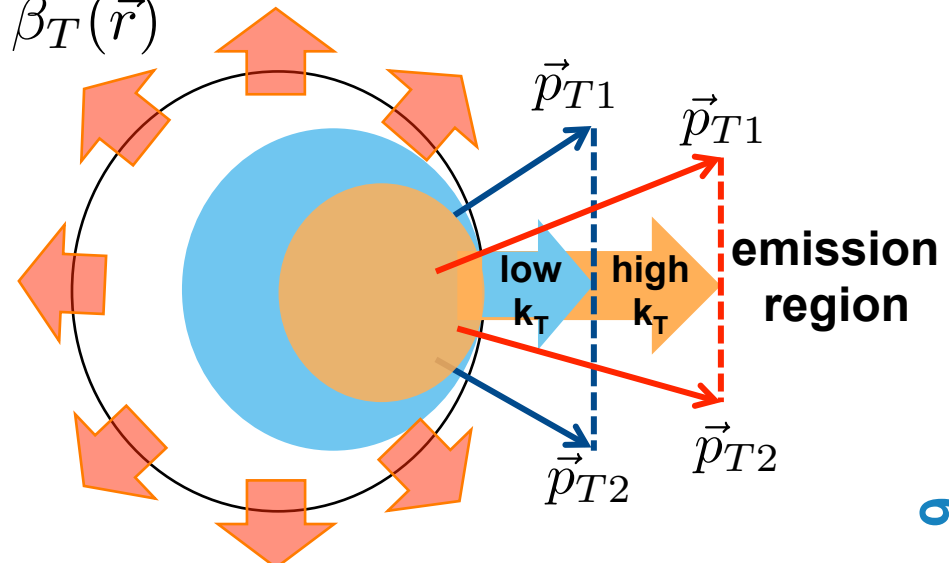


static



expanding

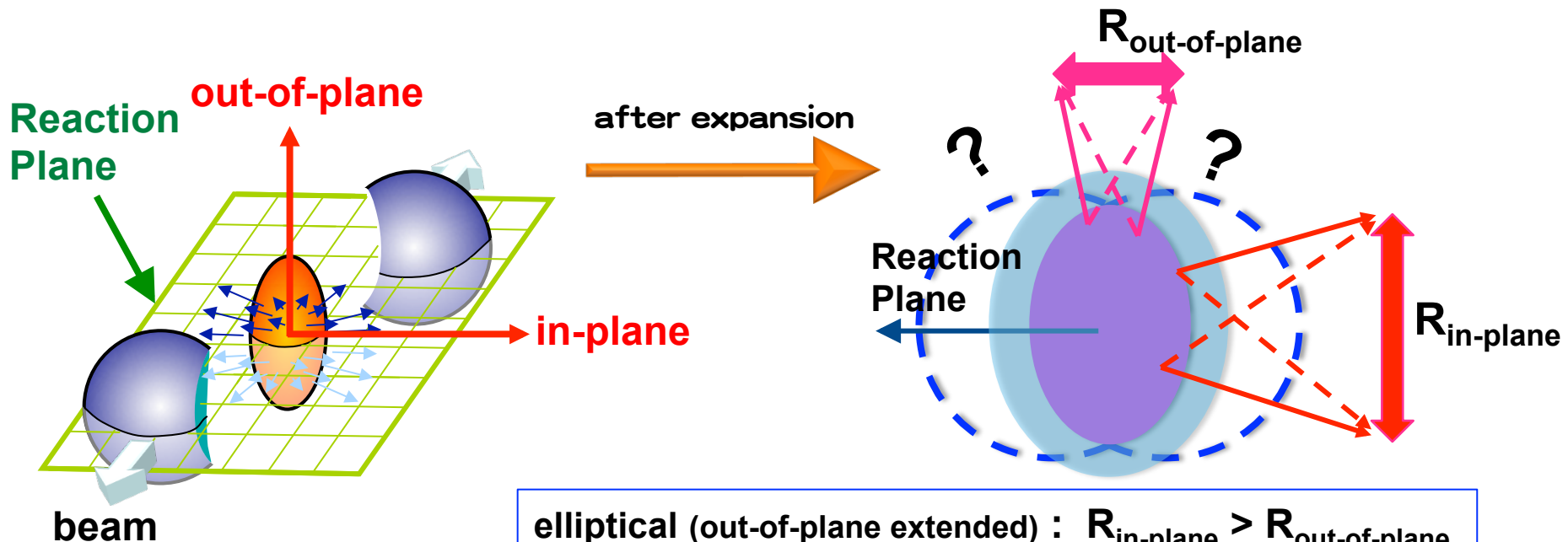
$$\vec{\beta}_T = \vec{\beta}_T(\vec{r})$$



$$\vec{k}_T = \frac{1}{2} (\vec{p}_{T1} + \vec{p}_{T2})$$

Azimuthal angle dependence

- Angle dependence of HBT radii w.r.t Reaction Plane reflects the source shape at freeze-out.
 - ✧ R.P is defined by beam axis and a vector between centers of colliding nuclei
- Initial spatial anisotropy causes momentum anisotropy (flow anisotropy)
- Final source eccentricity will be determined by **initial eccentricity**, **flow profile**, **expansion time**, and **viscosity** etc.

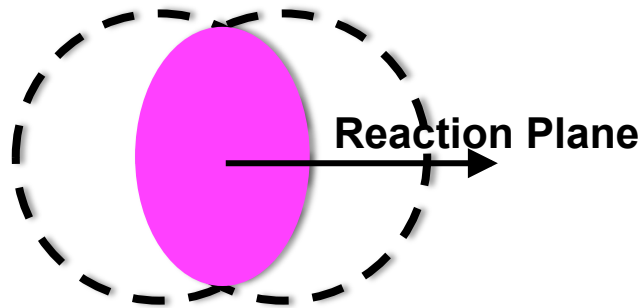


elliptical (out-of-plane extended)	:	$R_{\text{in-plane}} > R_{\text{out-of-plane}}$
spherical	:	$R_{\text{in-plane}} = R_{\text{out-of-plane}}$
elliptical (in-of-plane extended)	:	$R_{\text{in-plane}} < R_{\text{out-of-plane}}$

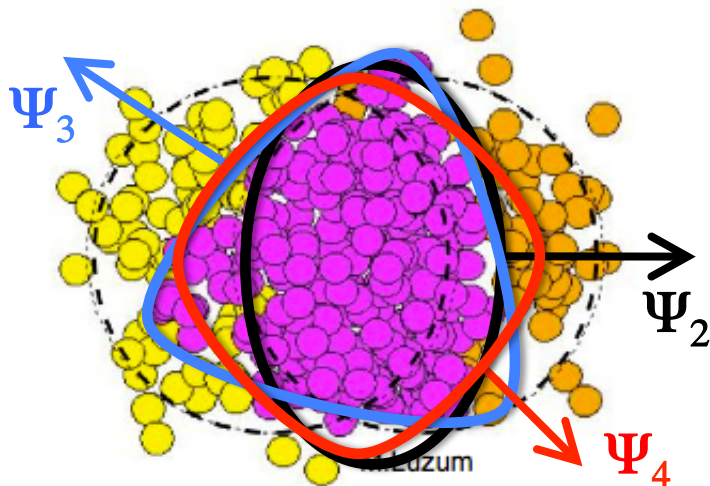
Higher Harmonic Flow and Event Plane

- Initial density fluctuations cause higher harmonic flow v_n
- Azimuthal distribution of emitted particles:

smooth picture



fluctuating picture



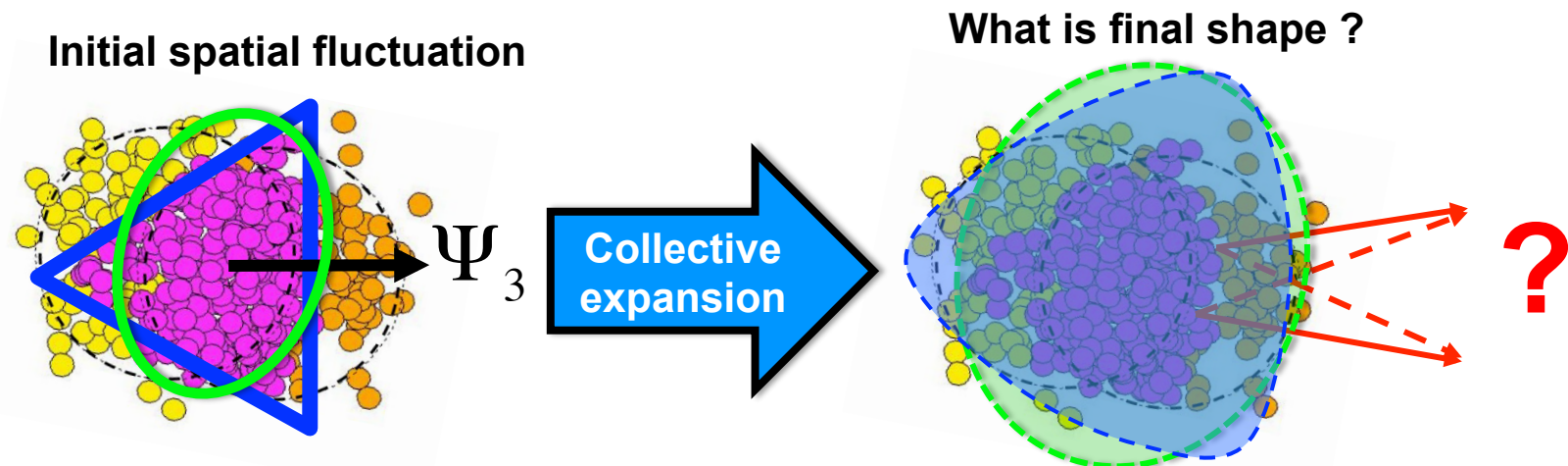
$$\frac{dN}{d\phi} \propto 1 + 2v_2 \cos 2(\phi - \Psi_2) + 2v_3 \cos 3(\phi - \Psi_3) + 2v_4 \cos 4(\phi - \Psi_4)$$

$$v_n = \langle \cos n(\phi - \Psi_n) \rangle$$

v_n : strength of higher harmonic flow
 Ψ_n : higher harmonic event plane
 ϕ : azimuthal angle of emitted particles

Motivation

- **Study the properties of space-time evolution in the heavy ion collisions via azimuthal HBT measurement.**
 - ✧ Measurement of charged pion HBT radii with respect to 2nd and 3rd-order event planes to reveal the detail of final state and space-time evolution of the system.
 - ✧ Study particle species dependence by comparison of charged pion and kaon HBT radii

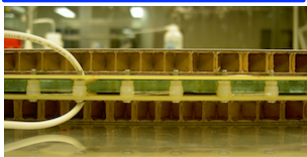


My Activities

3-years vacuum (It's mystery!)

2006(M1)

MRPC
construction



2007(M2)

RXNP
construction



Installed MRPC&RXNP

Azimuthal HBT analysis using Run4 data

2010(D1)

Di-jet Calorimeter
construction for ALICE

Summer Challenge (SC) @KEK

Shift taking @CERN

Start azimuthal HBT analysis using Run7 data

2011(D2)

Talk
WPCF2011

2012(D3)

2013(D4)

Talk
JPS fall

SC @KEK

Shift taking & Detector Expert
for Run11 @BNL

Talk
JPS spring

Talk
HIC in LHC

Talk
QM2012

SC @KEK

Shift taking & Detector Expert
for Run12 @BNL

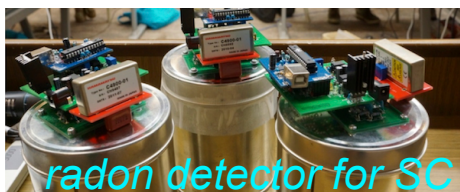
Poster
Radon Workshop

Talk
JPS spring

Work for Run13 @BNL

Talk
APPC12

SC @KEK



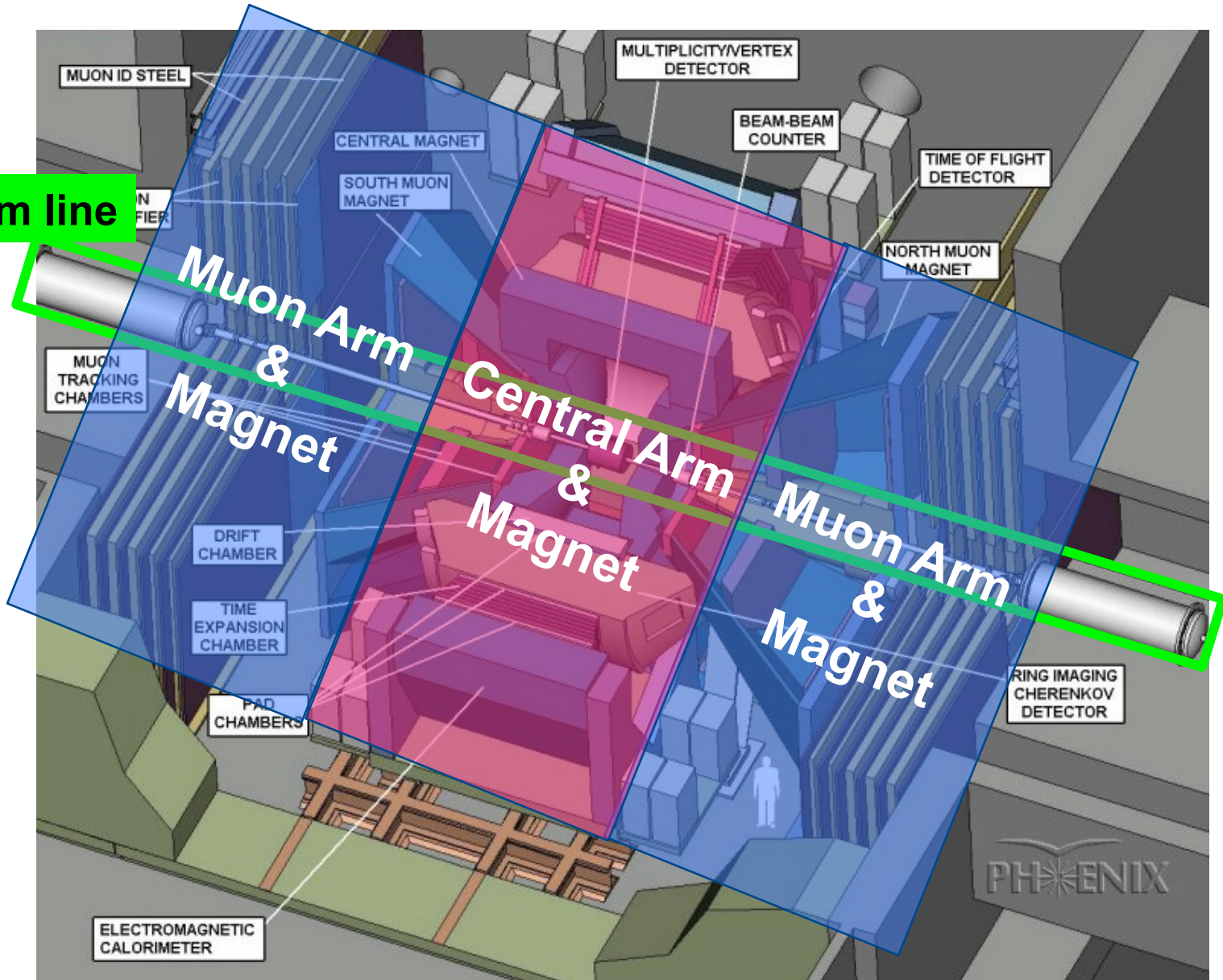
radon detector for SC

- ★ preliminary result for π/K HBT w.r.t Ψ_2
- ★ preliminary result for π/K HBT w.r.t Ψ_3
- ☆ Working for publication now

Experiment/ Analysis

PHENIX Experiment

Beam line



Roles of Detectors

Beam-Beam Counter ($|\eta|=3\sim 4$)

Quartz radiator+64PMTs

Zero Degree Calorimeter

Spectator neutron energy

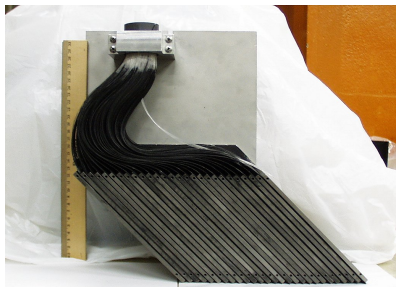
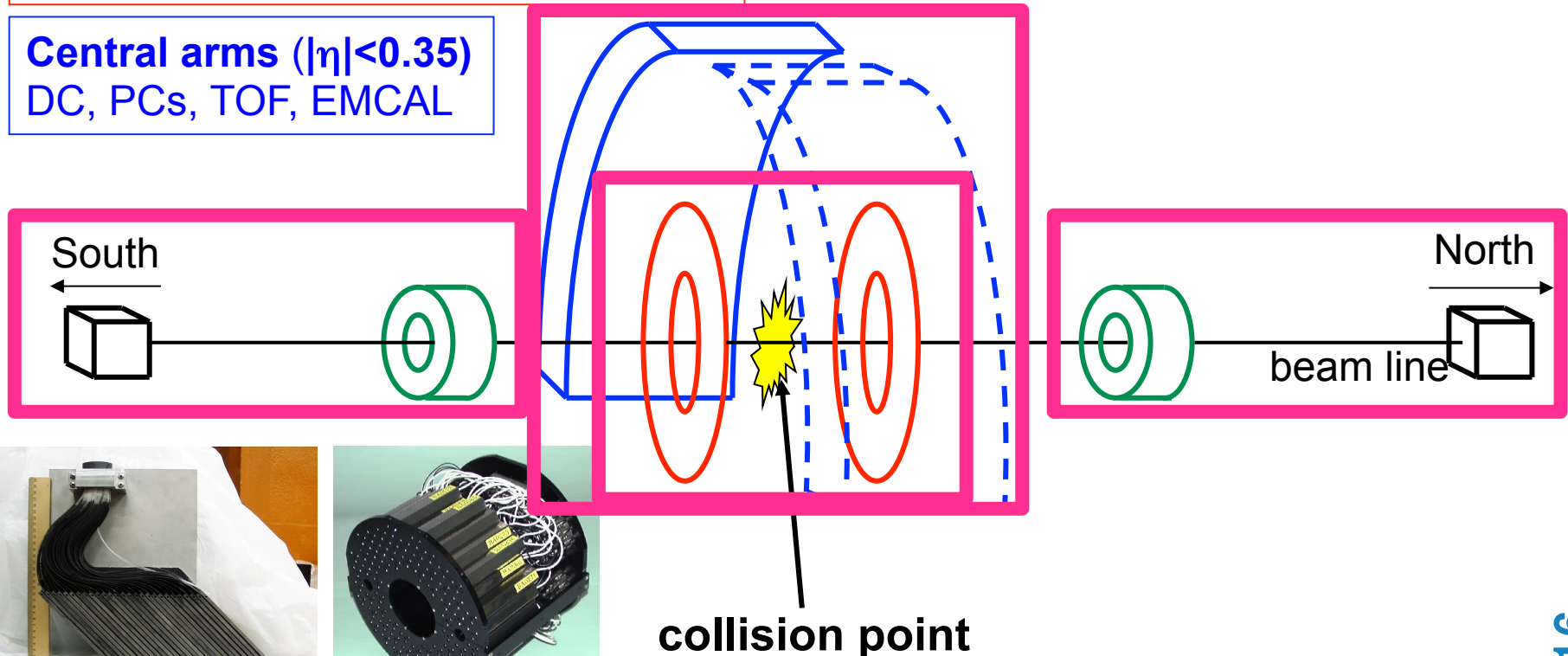
Reaction Plane Detector ($|\eta|=1\sim 2.8$)

2 rings of 24 scintillators

Central arms ($|\eta|<0.35$)

DC, PCs, TOF, EMCAL

- ⇒ Minimum Bias Trigger
- ⇒ Start time
- ⇒ Collision z-position
- ⇒ Centrality
- ⇒ Event Planes
- ⇒ Tracking, Momentum
- ⇒ Particle Identification



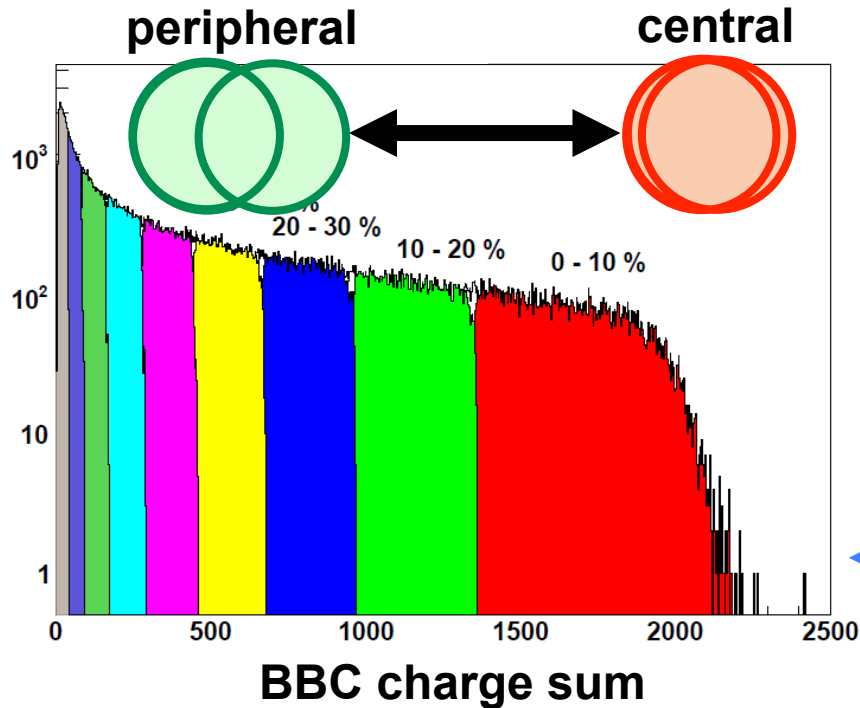
Collision Centrality

- Centrality is used to classify events instead of impact parameter.

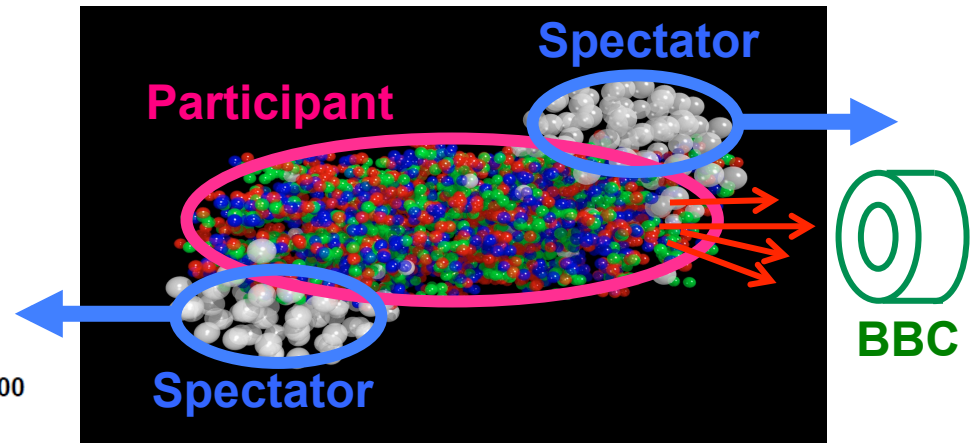
✧ impact parameter \propto multiplicity \propto charge sum at BBC

✧ 0% : head-head collision

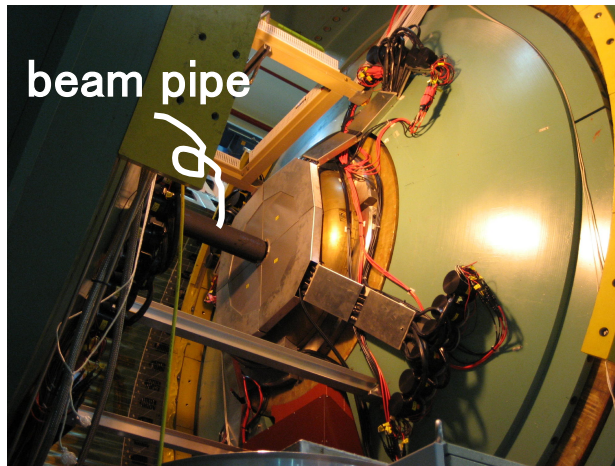
92% : most peripheral collisions



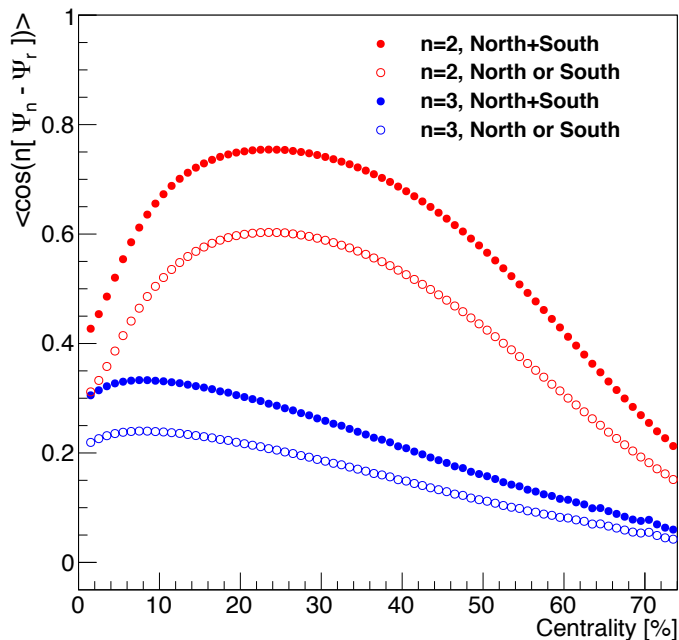
quartz radiator
+64PMTs



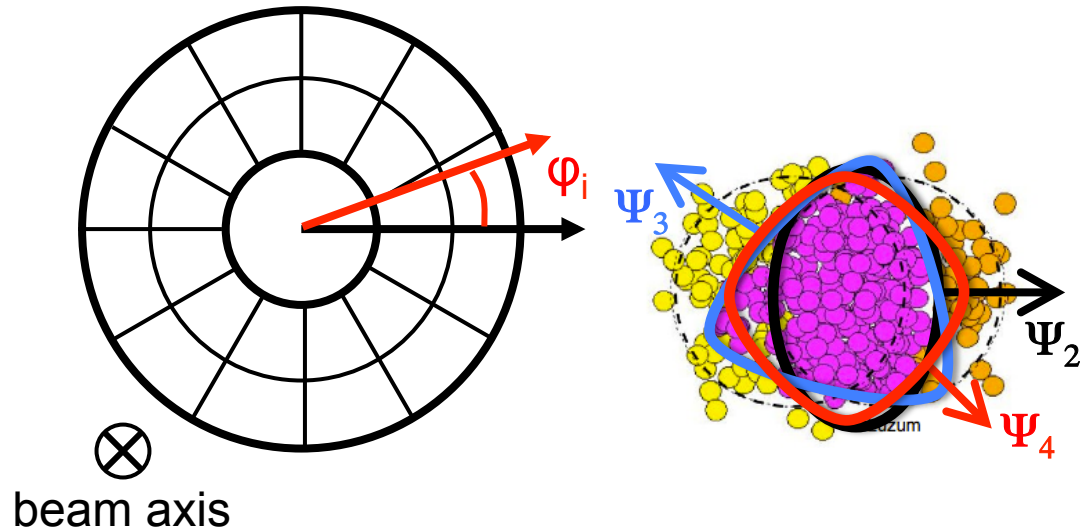
Event Plane



Resolution of Event planes



24 scintillator segments



$$\Psi_n = \frac{1}{n} \tan^{-1} \left(\frac{\sum w_i \cos(n\phi_i)}{\sum w_i \sin(n\phi_i)} \right)$$

■ Event plane was determined by Reaction Plane Detector (RXNP)

✧ Resolution: $\langle \cos(n(\Psi_n - \Psi_{\text{real}})) \rangle$

n=2 : ~ 0.75

n=3 : ~ 0.34

■ Determined by anisotropic flow itself

Particle Identification

- **EMC-PbSc is used.**

- ✧ timing resolution ~ 600 ps

- **Time-Of-Flight method**

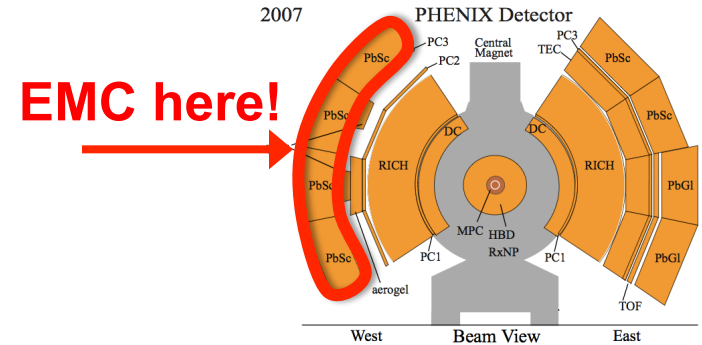
$$m^2 = p^2 \left(\left(\frac{ct}{L} \right)^2 - 1 \right)$$

p: momentum L: flight path length
t: time of flight

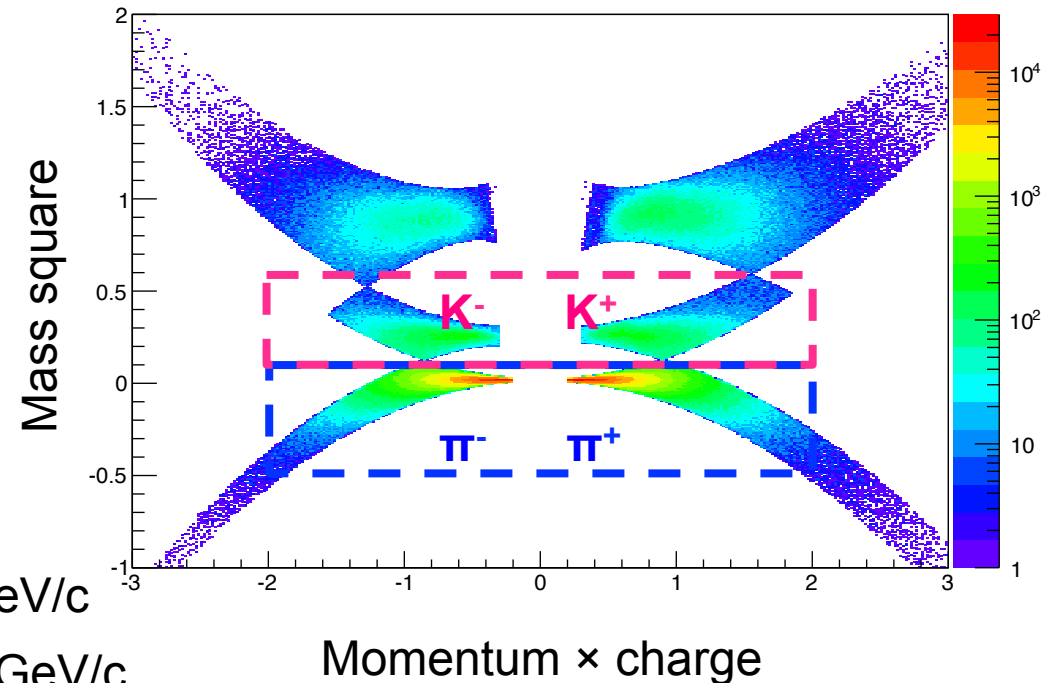
- **Charged π/K within 2σ**

- ✧ π/K separation up to ~ 1 GeV/c

- ✧ K/p separation up to ~ 1.6 GeV/c



Particle Identification by PbSc-EMC



Correlation Function

■ Experimental Correlation Function C_2 is defined as:

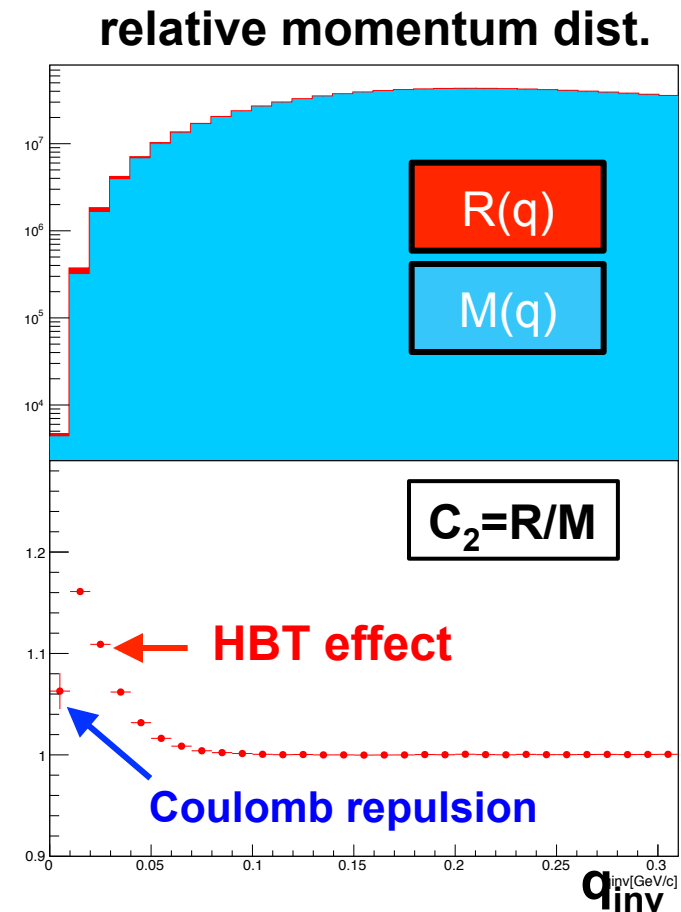
- ✧ $R(q)$: **R**ead pairs at the same event.
- ✧ $M(q)$: **M**ixed pairs selected from different events.

Event mixing was performed using events with similar z-vertex, centrality, E.P.

$$C_2 = \frac{R(\mathbf{q})}{M(\mathbf{q})}$$

$$\mathbf{q} = \mathbf{p}_1 - \mathbf{p}_2$$

- ✧ Real pairs include HBT effects, Coulomb interaction and detector inefficient effect. Mixed pairs doesn't include HBT and Coulomb effects.



3D-Analysis

■ “Out-Side-Long” frame

- ✧ Bertsch-Pratt parameterization
- ✧ Longitudinal Center of Mass System ($p_{z1}=p_{z2}$)

$C_2 = 1 + \lambda G$ λ : chaoticity R_μ : HBT radii

$$G = \exp(-\mathbf{R}^2 \mathbf{q}^2)$$

$$= \exp(-R_x^2 q_x^2 - R_y^2 q_y^2 - R_z^2 q_z^2 - \Delta\tau^2 q_0^2)$$

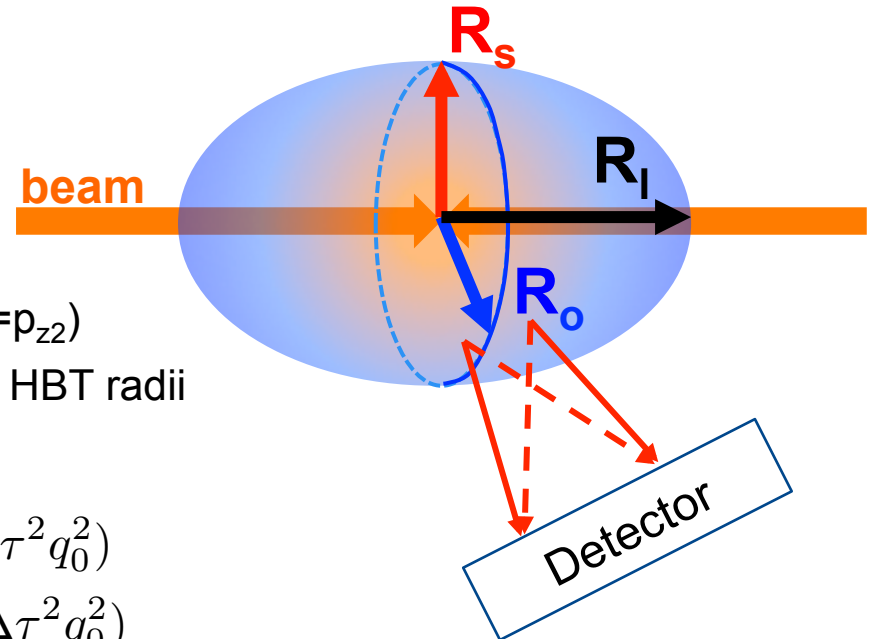
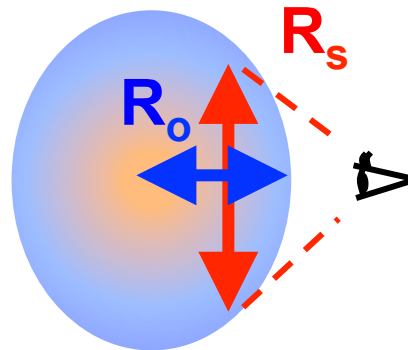
$$= \exp(-R_s^2 q_s^2 - R_o^{*2} q_o^2 - R_l^2 q_l^2 - \Delta\tau^2 q_0^2)$$

$$\stackrel{\text{LCMS}}{\approx} \exp(-R_s^2 q_s^2 - \underbrace{(R_o^{*2} + \beta_T \Delta\tau^2)}_{=\mathbf{R}_o^2} q_o^2 - R_l^2 q_l^2)$$

including cross term

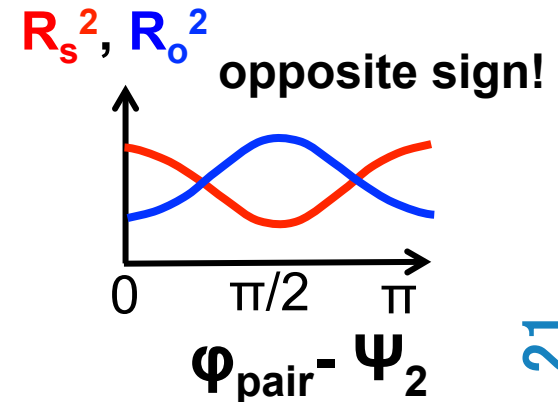
$$G = \exp(-R_s^2 q_s^2 - R_o^2 q_o^2 - R_l^2 q_l^2 - 2R_{os}^2 q_s q_o)$$

- R_s = “width” of source
- R_o = “thickness” of source



$$\vec{k}_T = \frac{1}{2}(\vec{p}_{T1} + \vec{p}_{T2})$$

$$\vec{q}_o \parallel \vec{k}_T, \quad \vec{q}_s \perp \vec{k}_T$$



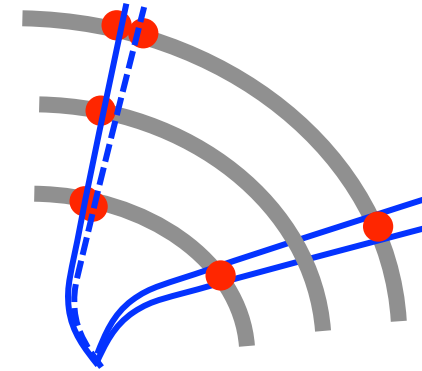
Pair Selection

■ Ghost Tracks

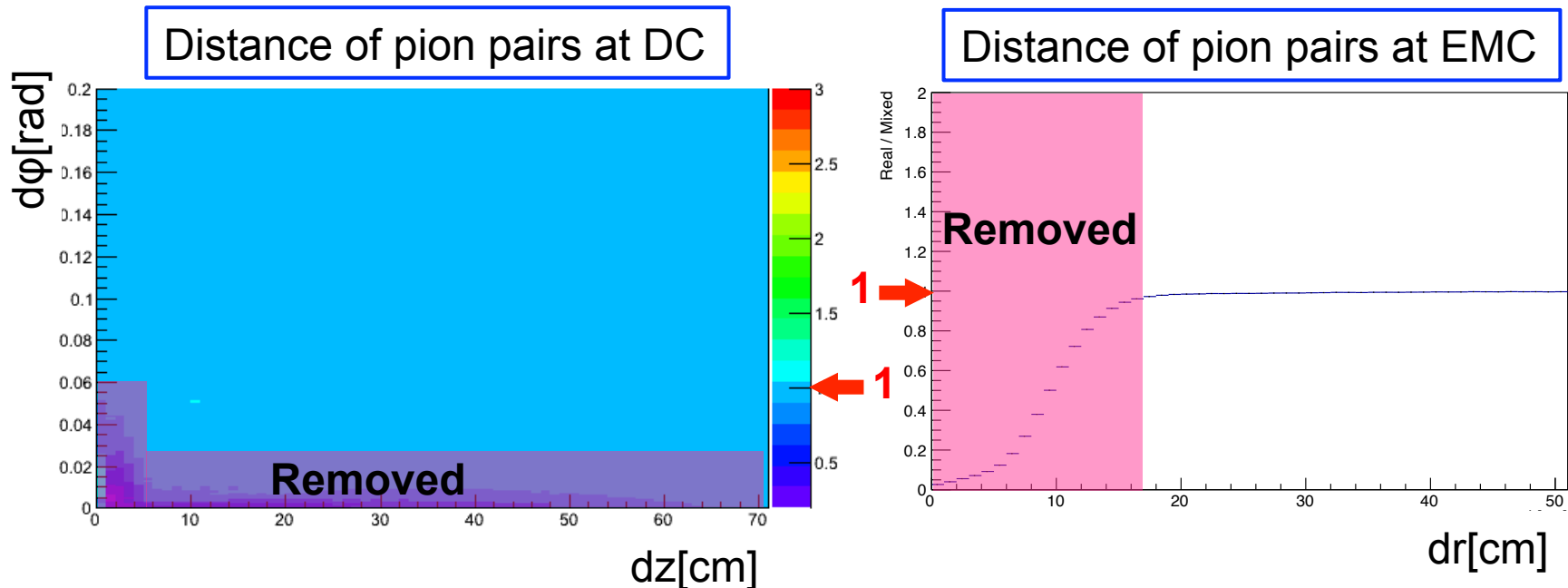
✧ A **single** particle is reconstructed as **two** tracks

■ Merged Tracks

✧ **Two** particles is reconstructed as a **single** track



- **Real/Mixed distribution of relative hit position should be unity in case of no mis-reconstruction and ideal detector efficiency.**



Coulomb Interaction

- Coulomb repulsion for like-sign pairs reduces pairs at low- q .

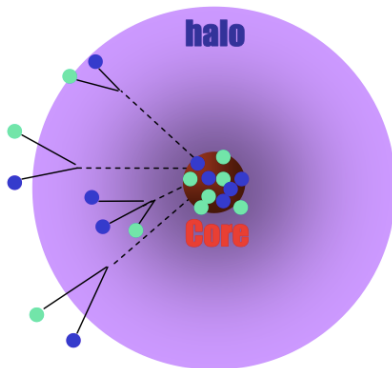
✧ Estimated by Coulomb wave function

$$\left[-\frac{\hbar^2 \nabla^2}{2\mu} + \frac{Z_1 Z_2 e^2}{r} \right] \Psi(r) = E \Psi(r) \quad \gamma = \frac{me^2}{\hbar^2 q} Z_1 Z_2$$

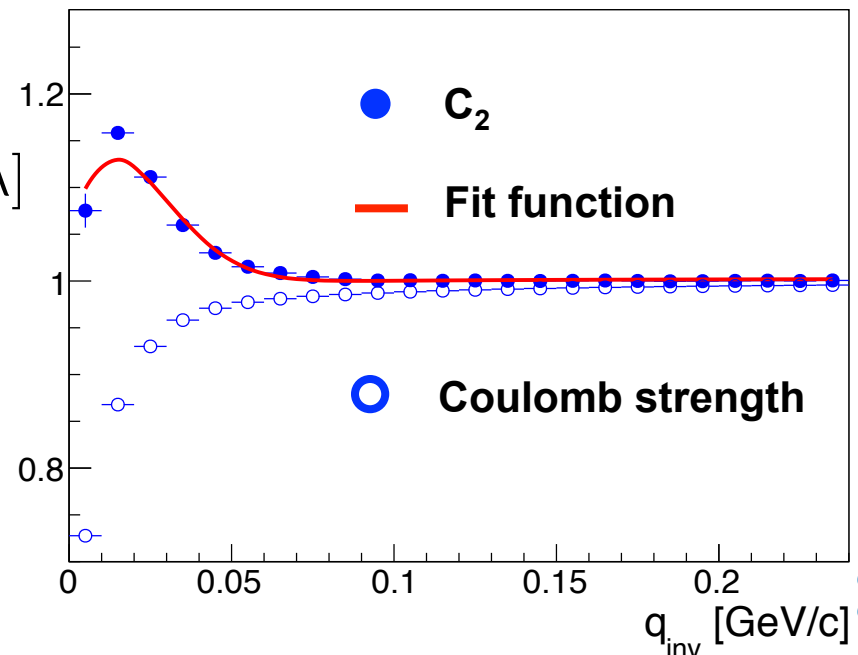
- The correction was applied in fit function for C_2

✧ Core-Halo model

$$\begin{aligned} C_2 &= C_2^{core} + C_2^{halo} \\ &= N[\lambda(1 + G)F_{coul}] + [1 - \lambda] \end{aligned}$$

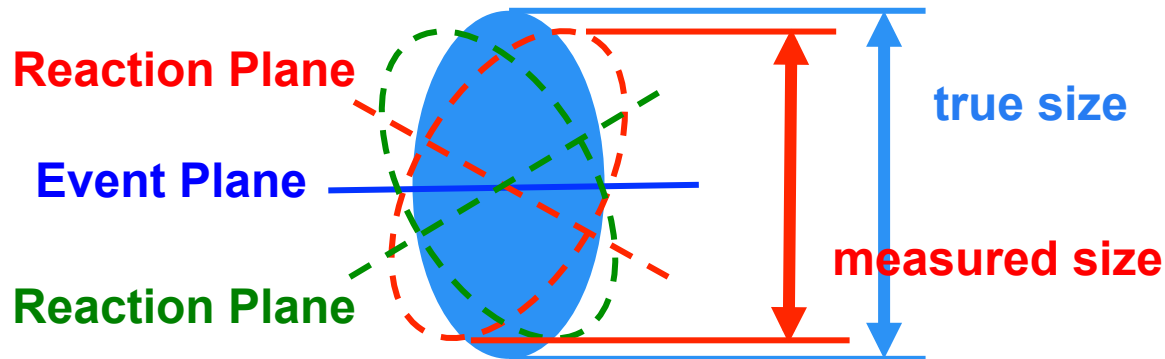


F_{coul} : Coulomb term
 G : Gaussian term



Correction of Event Plane Resolution

- Smearing effect by finite resolution of the event plane



Resolution correction

- correction for q-distribution

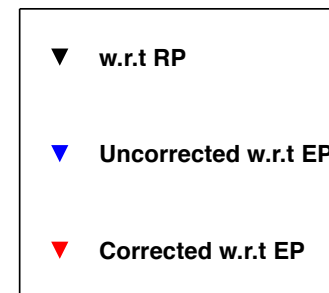
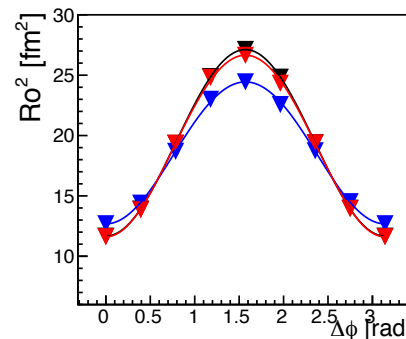
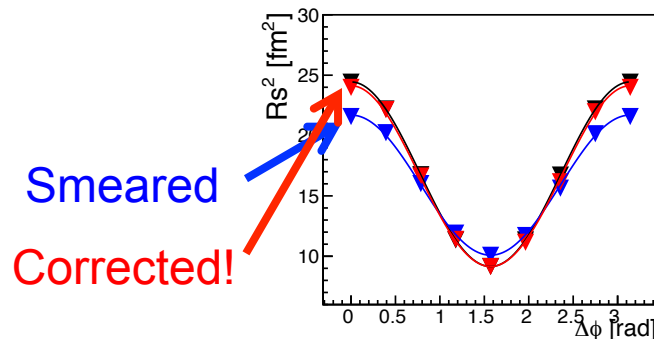
PRC.66, 044903(2002)

- Check by simulation

$$A_{corr}(q, \Phi_j) = A_{uncorr}(q, \Phi_j) + 2 \sum \zeta_{n,m} [A_c \cos(n\Phi_j) + A_s \sin(n\Phi_j)]$$

$$\zeta_{n,m} = \frac{n\Delta/2}{\sin(n\Delta/2) \langle \cos(n(\Psi_m - \Psi_{real})) \rangle}$$

event plane resolution



Systematic uncertainties

■ Track and Pair selection

- ✧ Track matching cut at Pad Chamber and EMCal (~5%)
- ✧ PID cut (~2%)
- ✧ Pair selection cut at Drift Chamber and EMCal (~6%)

■ Input source size for the Coulomb interaction (~2%)

■ Event plane determination (~4%)

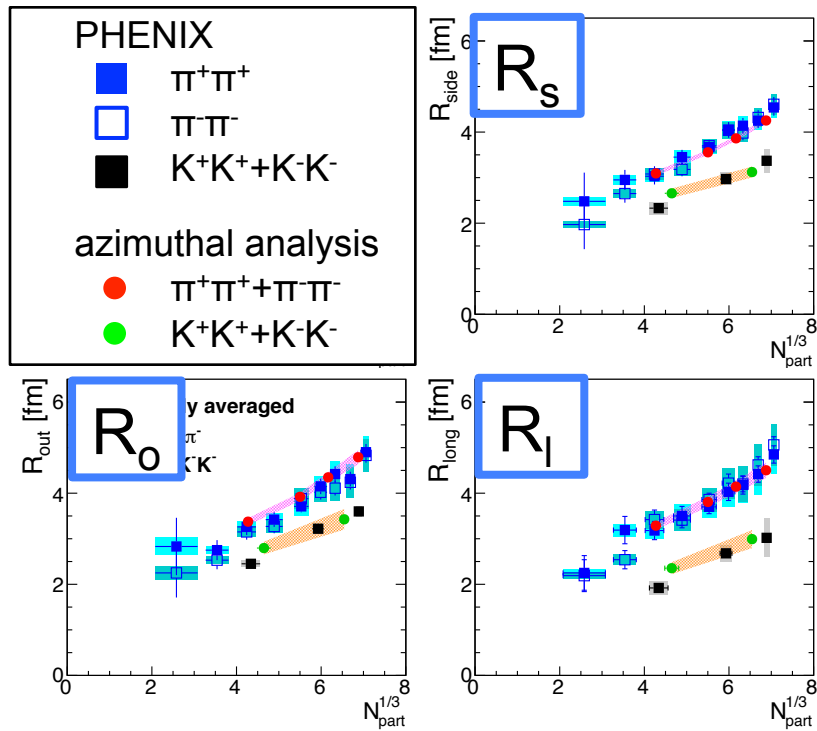
- ✧ Measured with north, south, and both combined RXNP

✧ The values within () are for R_{μ}^2 of pions

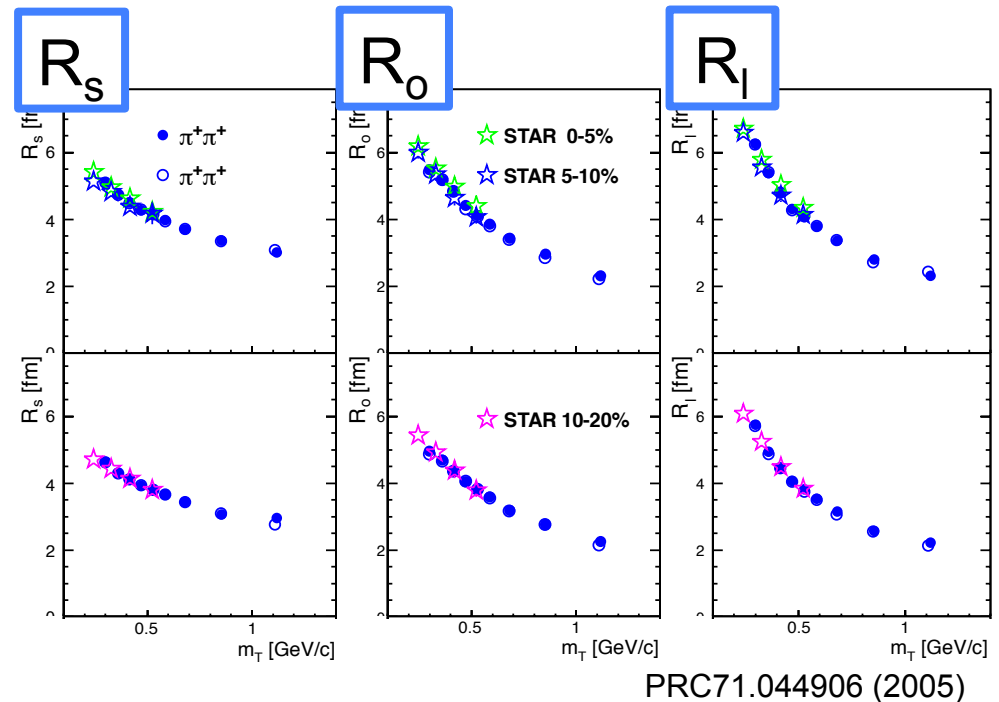
Consistency check with the previous results

- Extracted HBT radii are compared to the PHENIX and STAR results.

Centrality dependence



k_T dependence



PRC71.044906 (2005)

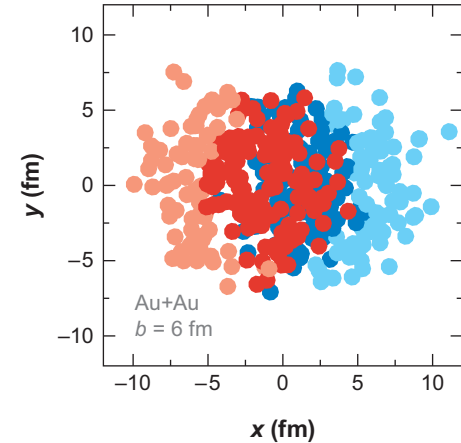
PRL93.152302(2004), PRL103.142301(2009)

Consistent with PHENIX and STAR results within systematic errors.

Initial spatial anisotropy

■ Estimated by Monte-Carlo Glauber simulation

- ✧ assuming a Woods-Saxon density
- ✧ collision takes place if $d < \sqrt{\sigma_{nn}/\pi}$
 - ✓ d : distance between nucleons
 - ✓ σ_{nn} : total cross section

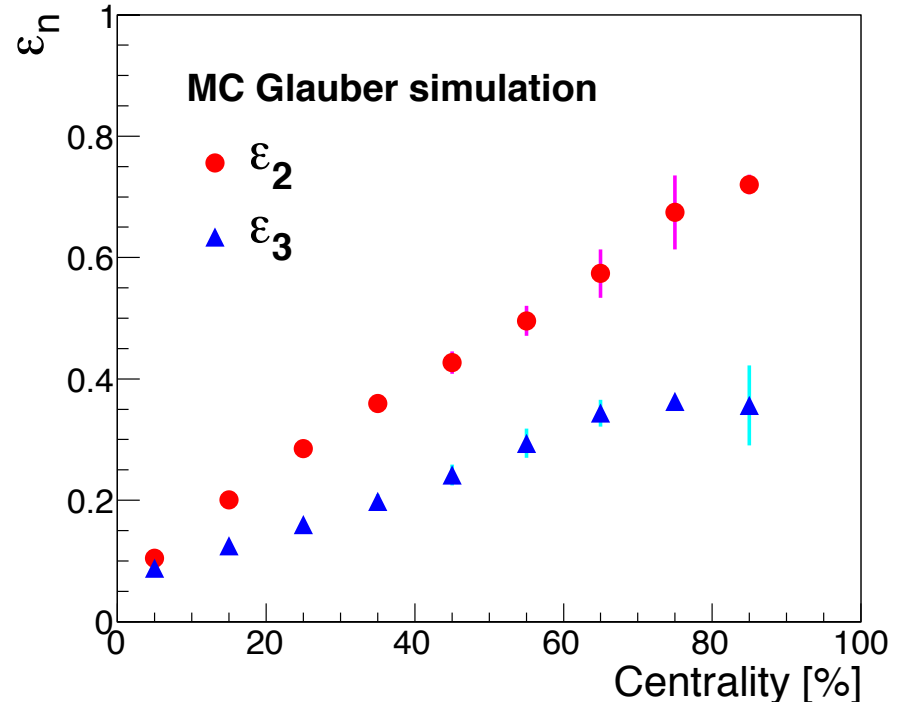


$$\varepsilon_n = \frac{\langle r^2 \cos[n(\phi - \Psi_n)] \rangle}{\langle r^2 \rangle}$$

$$\Psi_n = \frac{1}{n} \left[\tan^{-1} \left(\frac{\langle r^2 \sin(n\phi) \rangle}{\langle r^2 \cos(n\phi) \rangle} \right) + \pi \right]$$

■ Centrality dependence of initial ε_2 and ε_3 are seen.

- ✧ $\varepsilon_3 \sim \varepsilon_2/2$ at mid-central collision



Results/ Discussion

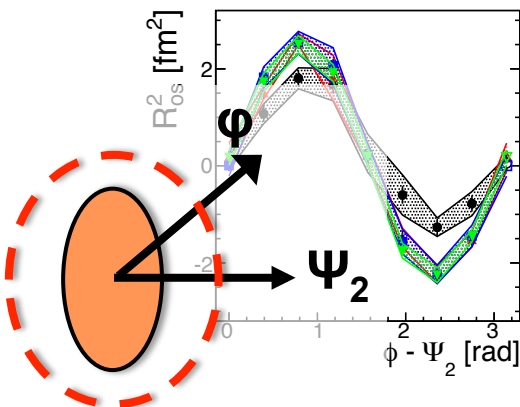
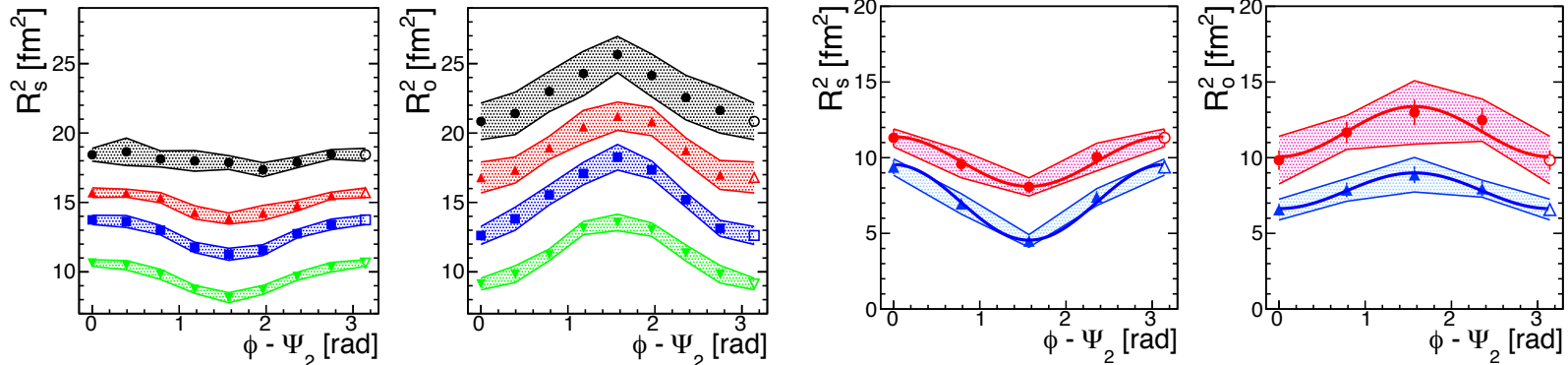
- Azimuthal angle dependence w.r.t 2nd-order event plane
- Comparison of π/K HBT radii
- Blast-wave model fit

- Azimuthal angle dependence w.r.t 3rd-order event plane
- Monte-Carlo simulation

Centrality dependence of pion HBT radii w.r.t Ψ_2

■ Oscillations of R_s , R_o , and R_{os} are seen for π/K

- ✧ Centrality dependence of R_s oscillation is seen.
- ✧ R_o has stronger oscillation than R_s in all centrality.



Note: R_o inherently depends on β_T & $\Delta\tau$

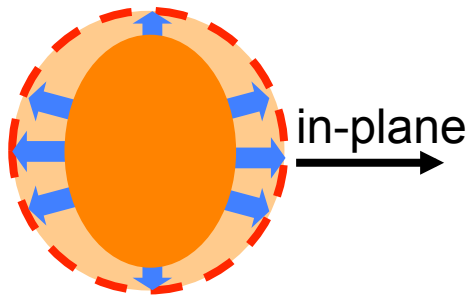
$$C_2 = 1 + \lambda \exp(-R_s^2 q_s^2 - R_o^2 q_o^2 - R_l^2 q_l^2 - 2R_{os}^2 q_o q_s)$$

$$R_o^2 = R_o^{*2} + \beta_T \Delta\tau^2$$

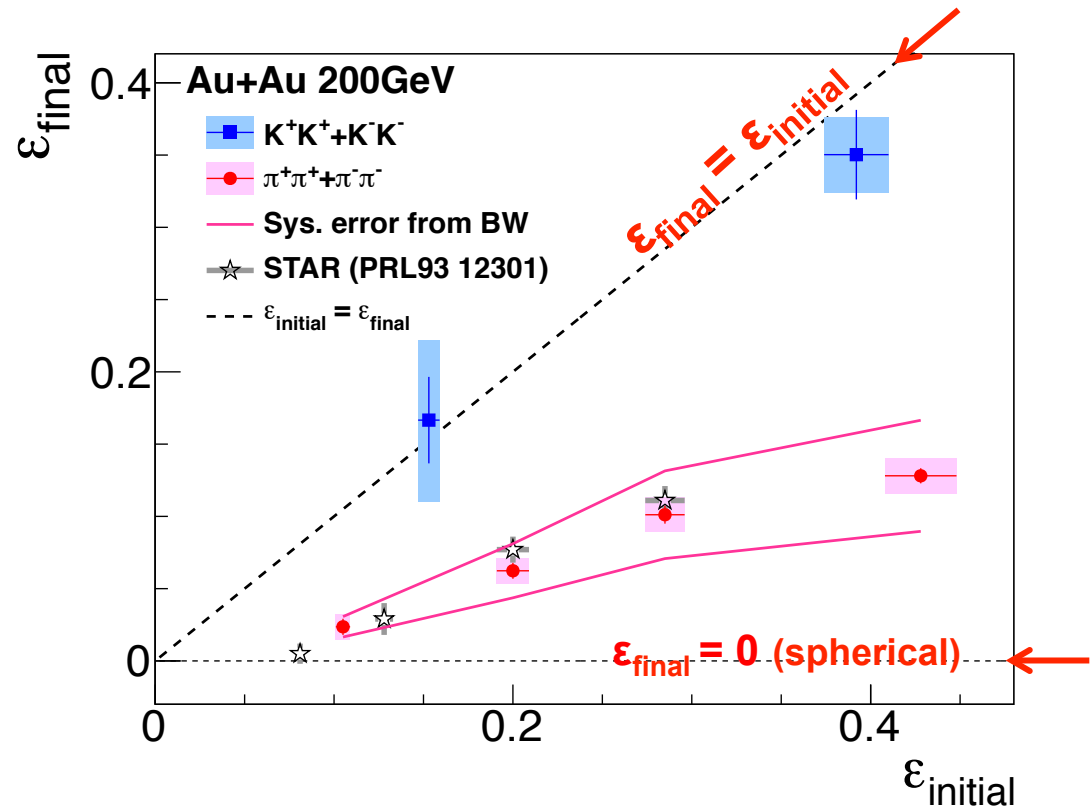
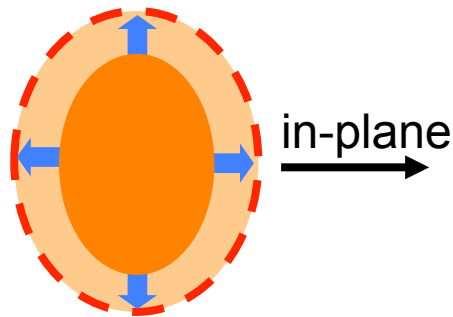
- effect of flow anisotropy?
- $\Delta\tau$ depend on azimuthal angle?
- difference of “width” and “thickness”?
- Study with Blast-wave model

Eccentricity at freeze-out

pion



kaon



■ $\epsilon_{\text{final}} \approx \epsilon_{\text{initial}}/2$ for pion

✧ Consistent with STAR result.

✧ This Indicates that source expands to in-plane direction, and still elliptical shape.

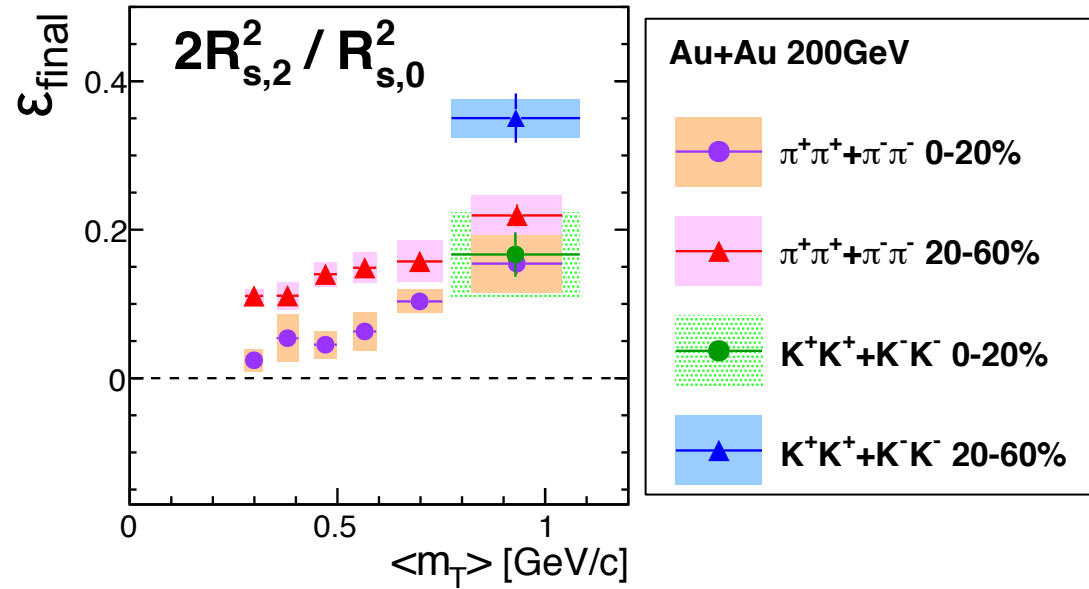
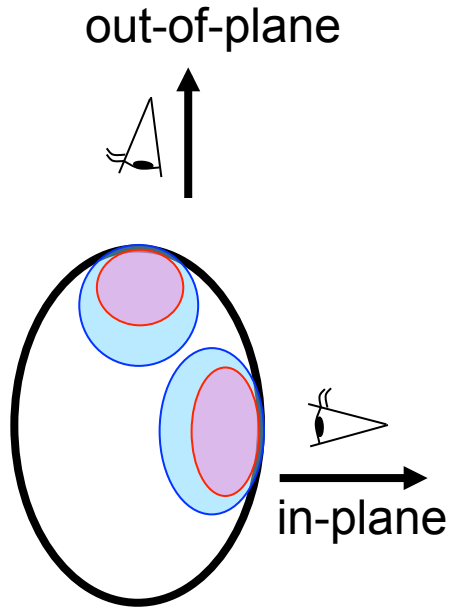
■ $\epsilon_{\text{final}} \approx \epsilon_{\text{initial}}$ for kaon

✧ Kaon may freeze-out sooner than pion due to less cross section ?

✧ Is the emission region different due to the different m_T ?

m_T dependence of ϵ_{final}

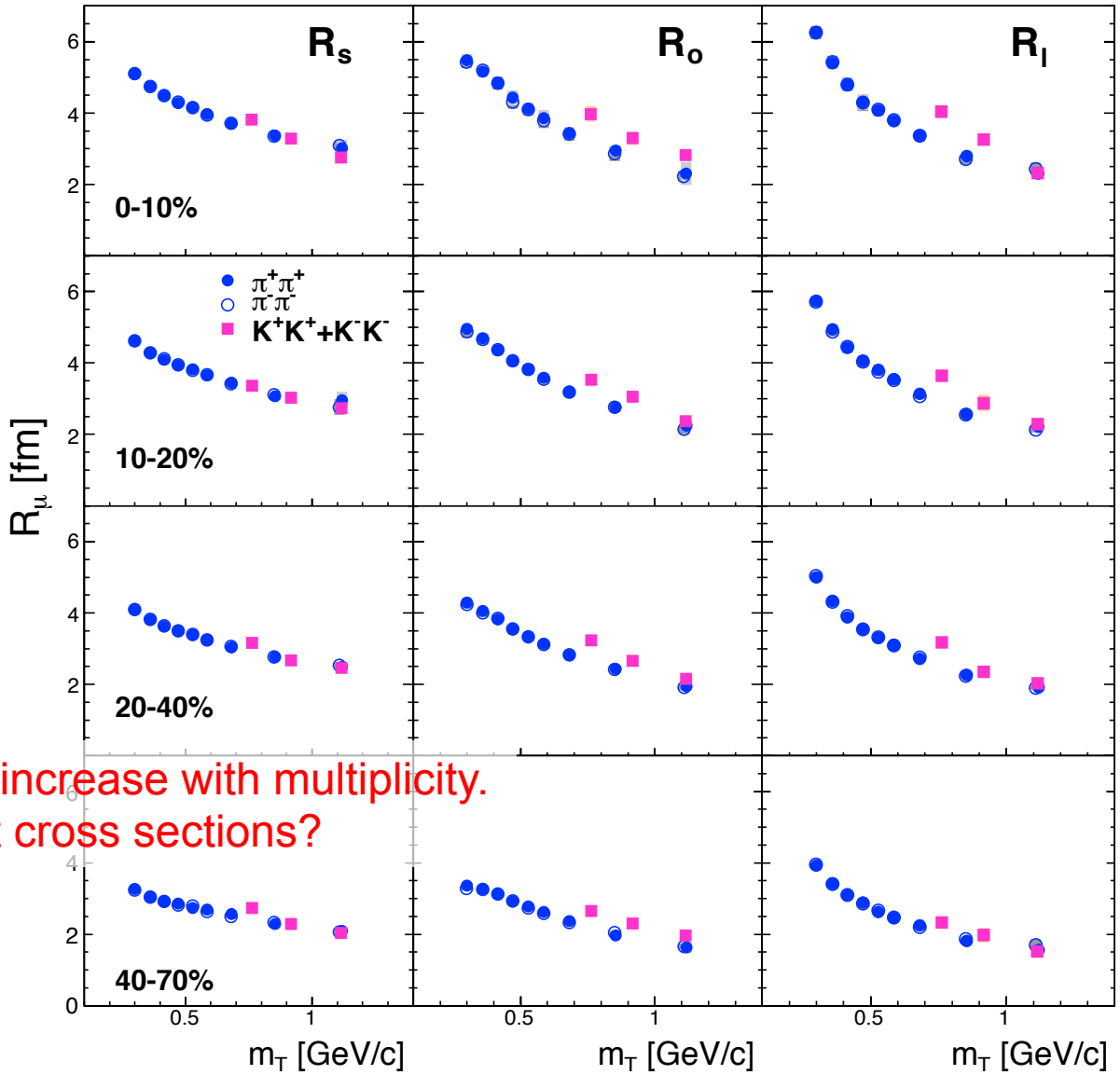
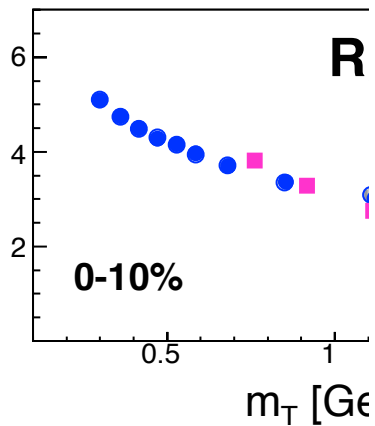
$$m_T = \sqrt{k_T^2 + m^2}$$



- ϵ_{final} of π increases with m_T
 - ✧ It would reflect the variation of the emission region.
- **Still difference between π/K in 20-60% even at the same m_T**
 - ✧ Indicates sooner freeze-out time of K than π ?
 - ✧ How about the average radii ?

Comparison of π/K HBT radii (1)

- Difference can be seen
- Charged kaons have



Difference seems to increase with multiplicity.
Effect of the different cross sections?

Comparison of π/K HBT radii (2)

■ Comparison of positive and negative kaons

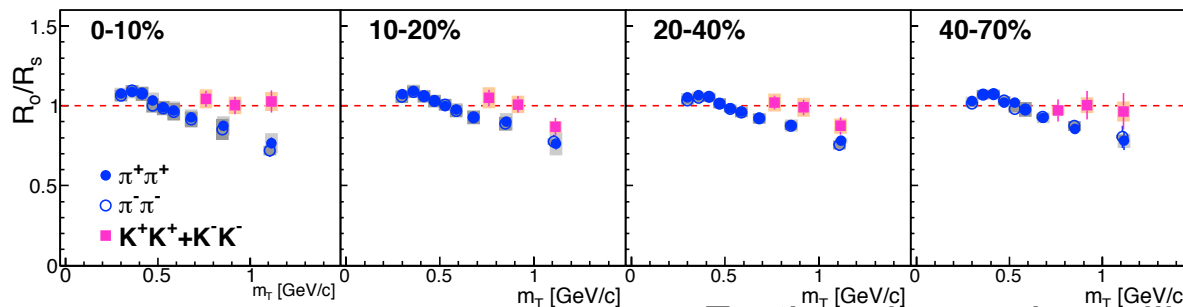
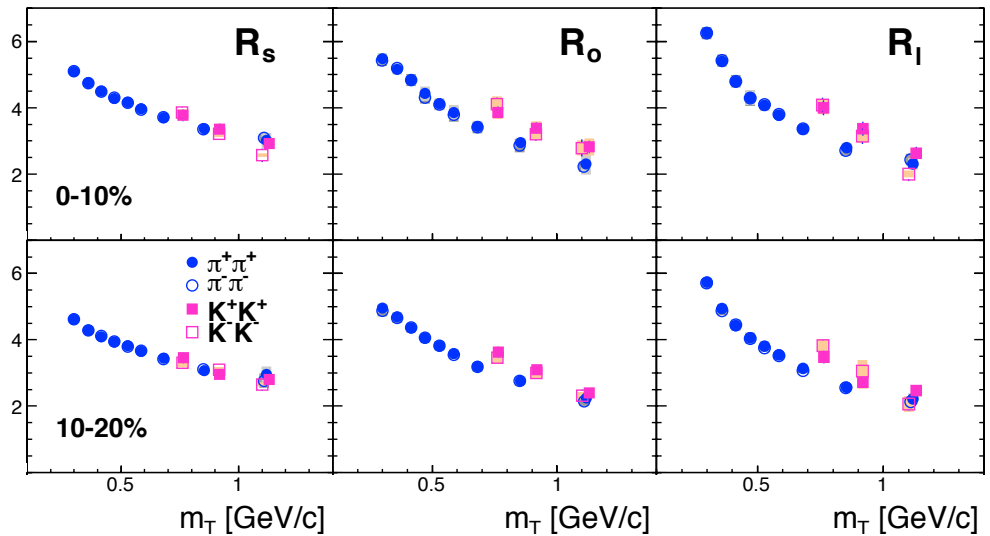
- ✧ K^+ have less σ_{total} on nucleons than K^-
 $\sigma(K^-)/\sigma(K^+) \approx 5$ at $p \approx 0.3$ [GeV/c] ($\sigma(K^-)/\sigma(\pi^-) \approx 2$)

■ No significant difference between both kaons

- ✧ π/K difference may not be explained by the effect of hadronic rescattering.

■ R_o/R_s for the emission duration

- ✧ No centrality and m_T dependence
- ✧ $K > \pi$: relatively longer emission duration of K ?



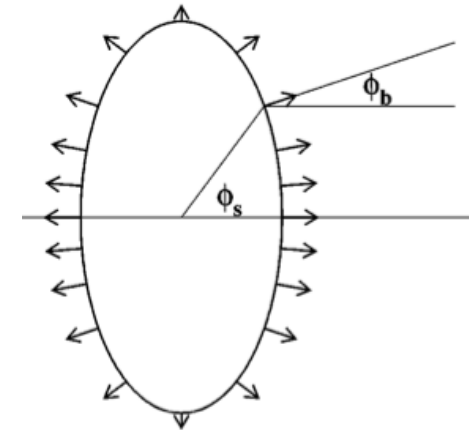
Further discussion will be done later.

Interpretation by Blast wave model

■ Hydrodynamic model parameterized with freeze-out conditions

- ✧ Thermal equilibrium + collective expansion
- ✧ Freeze-out takes place for all hadrons at the same time
- ✧ Well reproduced p_T spectra & elliptic flow at low p_T

PRC 70, 044907 (2004)



Free parameters which characterize freeze-out

- T_f : temperature at freeze-out
 ρ_0, ρ_2 : transverse rapidity
 R_x, R_y : transverse sizes (shape)
 $\tau_0, \Delta\tau$: freeze-out time, emission duration

* box profile is assumed as spatial density

$$\beta_T = \tanh(\rho)$$

$$\rho(r, \phi) = \tilde{r} [\rho_0 + \rho_2 \cos(2\phi)]$$

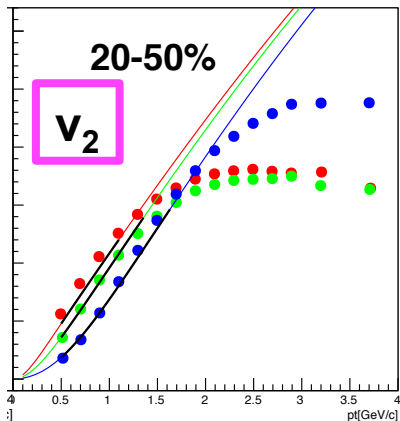
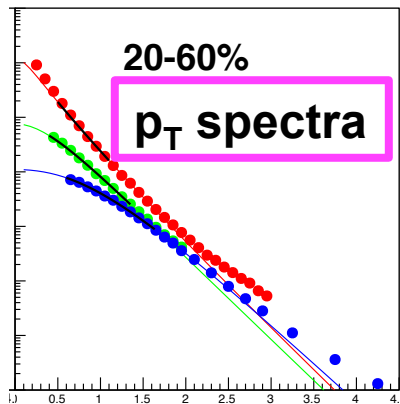
$$\tilde{r} = \sqrt{\left(\frac{r \cos(\phi)}{R_x}\right)^2 + \left(\frac{r \sin(\phi)}{R_y}\right)^2}$$

■ Motivation of BW study

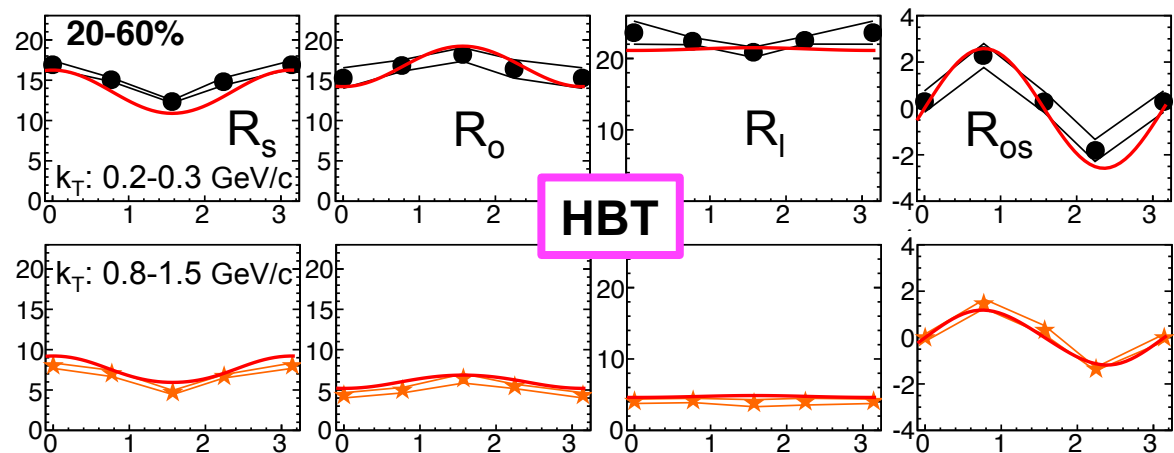
- ✧ Are π/K HBT results reproduced by the same freeze-out parameters?
- ✧ Extract the temporal information

Fit by Blast-wave model

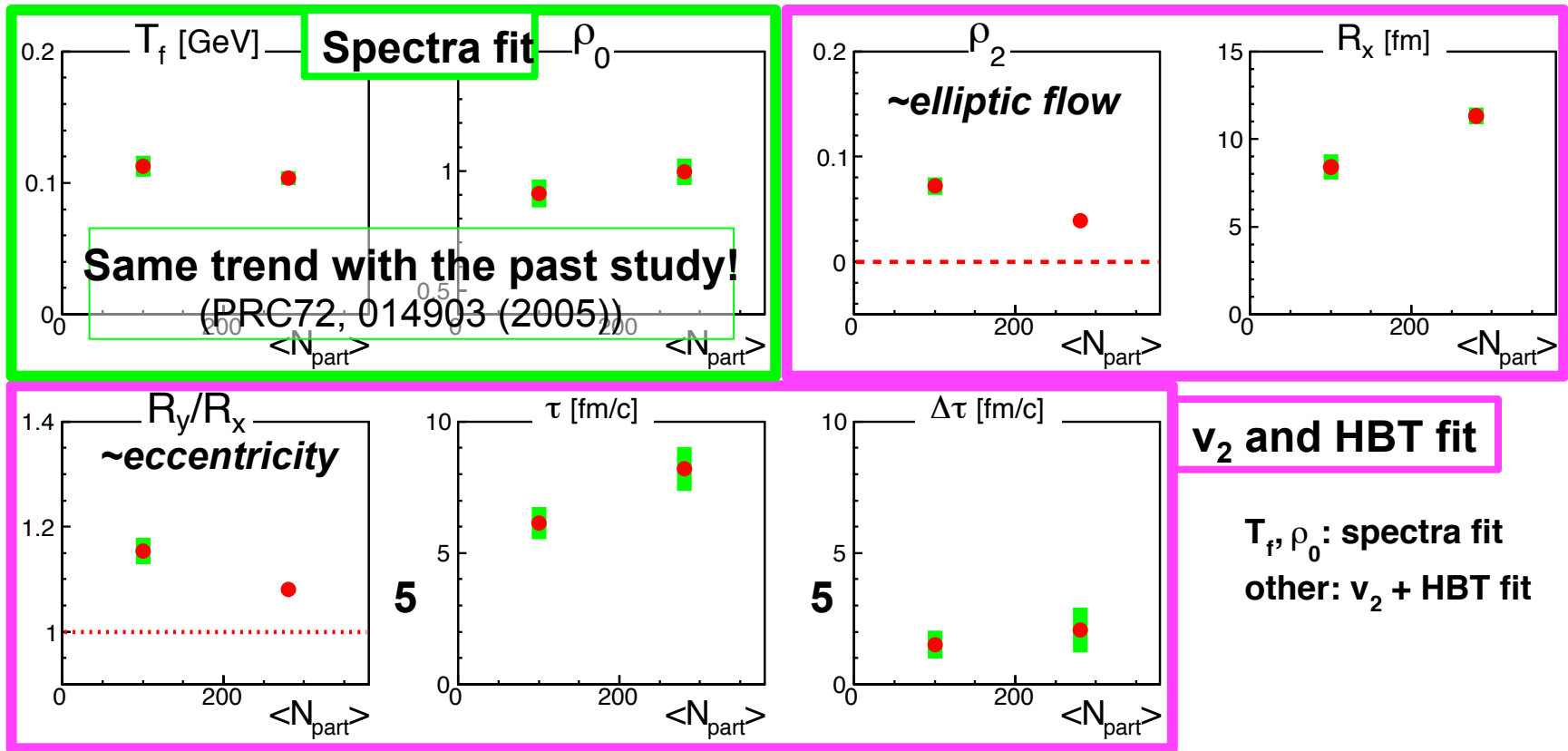
- Transverse momentum distribution (p_T spectra) and elliptic flow v_2 are used to effectively reduce the BW parameters.



1. Fit p_T spectra to obtain T_f and ρ_0
 - spectra data from PHENIX (PRC69,034909(2004))
2. Fit v_2 and HBT radii for all k_T simultaneously
 - $\rho_2, R_x, R_y, \tau_0, \Delta\tau$ are obtained.



Extracted freeze-out parameters



- ρ_2 and ellipticity(R_y/R_x) increase with decreasing N_{part}

✧ Reasonable in terms of v_2 and ϵ_f

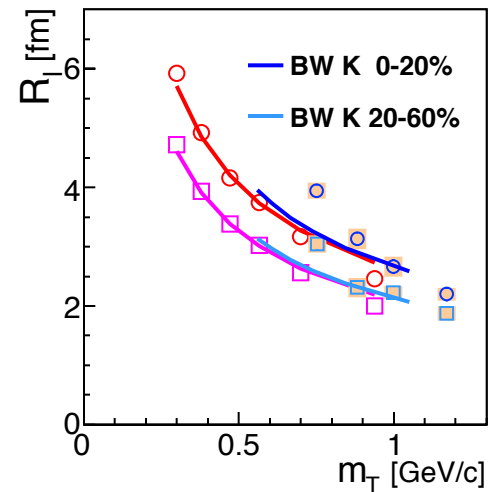
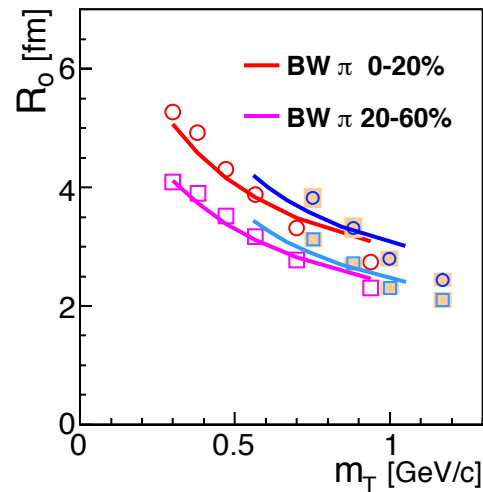
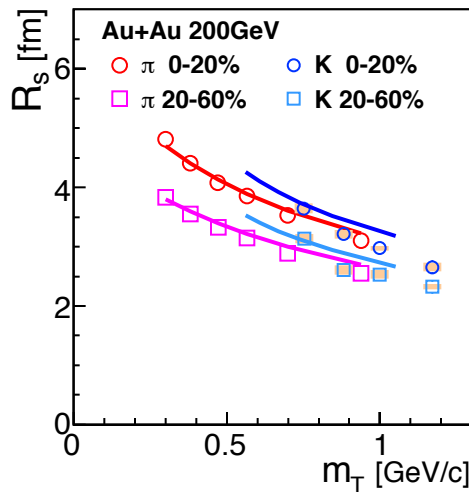
- τ and $\Delta\tau$ increases with increasing N_{part}

✧ Freeze-out time: 6~8 fm/c, Emission duration: 1.5~2fm/c

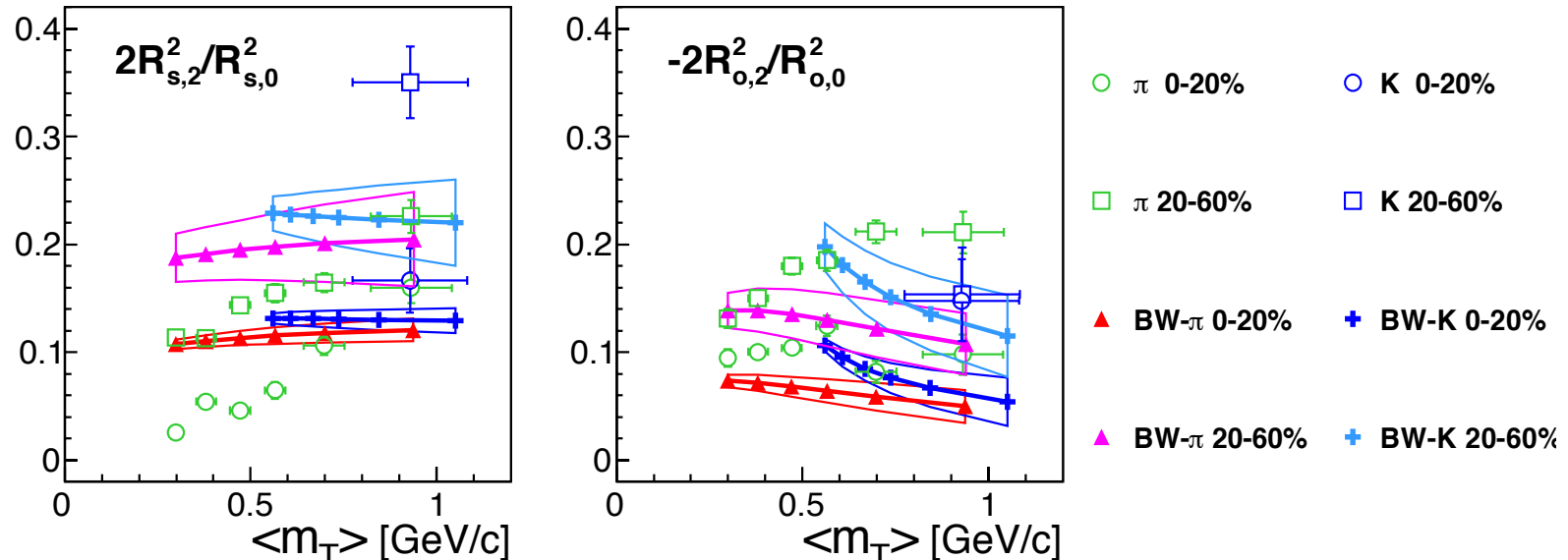
✓ Similar value ($\tau=8.6$, $\langle\Delta\tau\rangle=2$) in Source imaging+Therminator model(PRC100.232301)

Comparison of mean radii between π/K

- Mean HBT radii of π are reproduced well by BW, but not for K
 - ✧ Using fit parameters obtained by fitting pion HBT radii
 - ✧ BW also expects larger radii of K in R_s and R_l
- Indicate different freeze-out mechanism?
 - ✧ e.g Freeze-out time or/and emission duration are different?

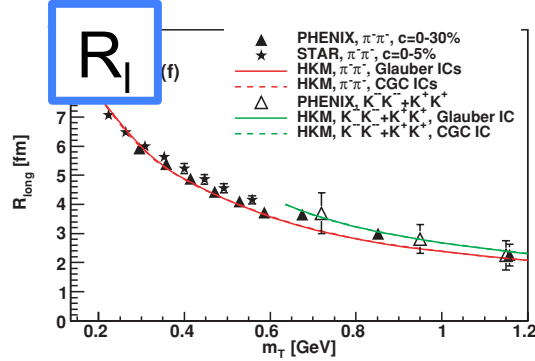
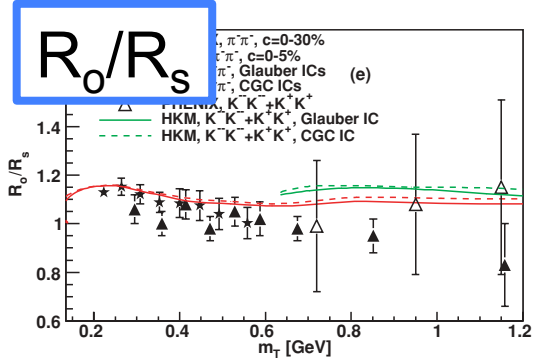
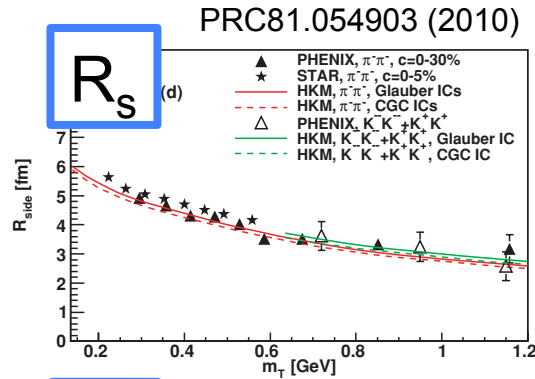
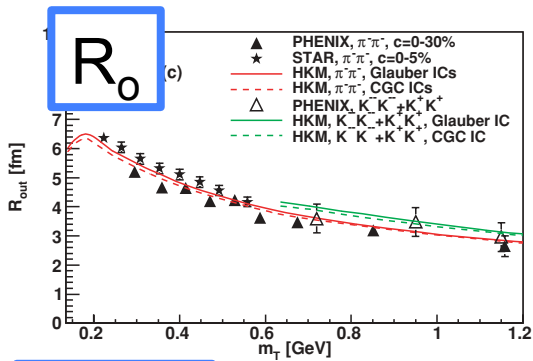


Comparison of oscillation amplitude between π/K



- BW doesn't also reproduce the oscillation amplitudes of R_s^2 quantitatively, and R_o^2 qualitatively.
- Underestimate of R_o oscillation in BW.
 - ✧ ϕ -dependent $\Delta\tau$?
 - ✧ Spatial density/flow profile is not appropriate?
- π/K difference in the oscillation cannot be explained by BW.

Recent theoretical result



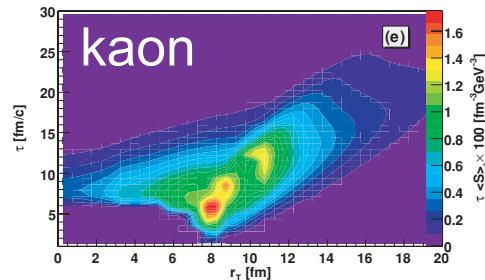
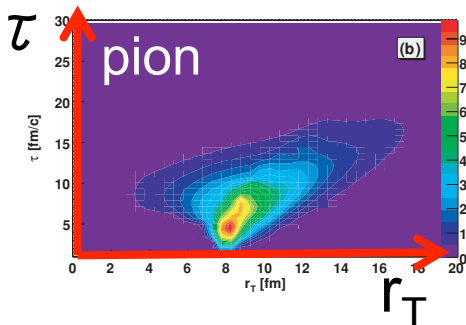
HydroKinetic Model

- Glauber or CGC
- crossover transition
- microscopic transport
- pre-thermal flow
- implicitly viscosity

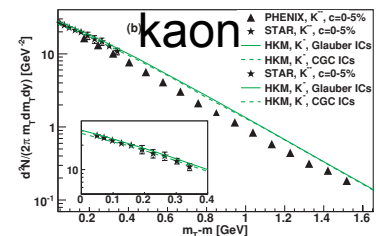
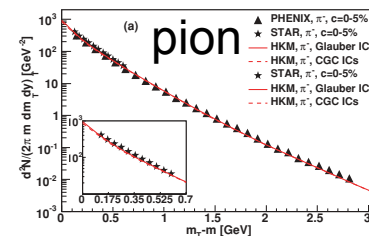
Larger R_0 and R_l of kaon

- ◇ Similar trend to the measured ones
- ◇ Majority of kaons leave system later than pions at the same m_T
- Longer emission duration ?

Emission function at $m_T=0.86\text{GeV}/c$



✳Note that kaon spectra is not reproduced well at high m_T



Summary-1

■ Azimuthal HBT radii w.r.t 2nd-order event plane

✧ pion ϵ_{final} is $\sim \epsilon_{\text{initial}}/2$

✓ Strong expansion to in-plane direction and still elliptical shape.

✧ kaon $\epsilon_{\text{final}} >$ pion ϵ_{final} even at the same m_T

✧ Stronger oscillation of R_o

■ Comparison of π/K HBT radii

✧ Difference of R_o and R_l in most central collisions

✧ No significant difference between K^+K^+ and K^-K^-

■ Study with Blast-wave model

✧ Freeze-out time ~ 8 fm/c, $\Delta\tau \sim 2$ fm/c

✓ Consistent with source imaging analysis

✧ π/K difference of mean radii is not explained well

✧ Oscillation amplitudes is not reproduced well

✓ Stronger R_o oscillation is not explained only by flow anisotropy within BW

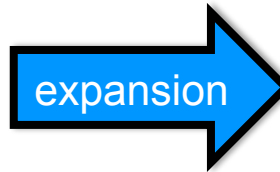
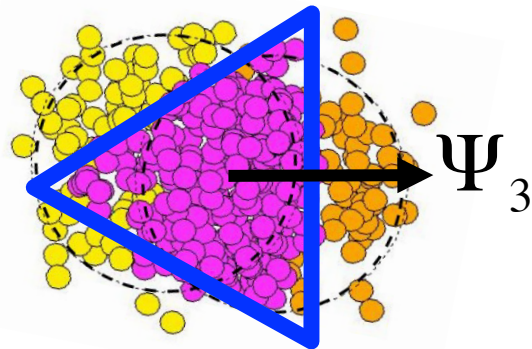
■ Comparison with HKM model

✧ Qualitatively consistent with larger R_o & R_l of K

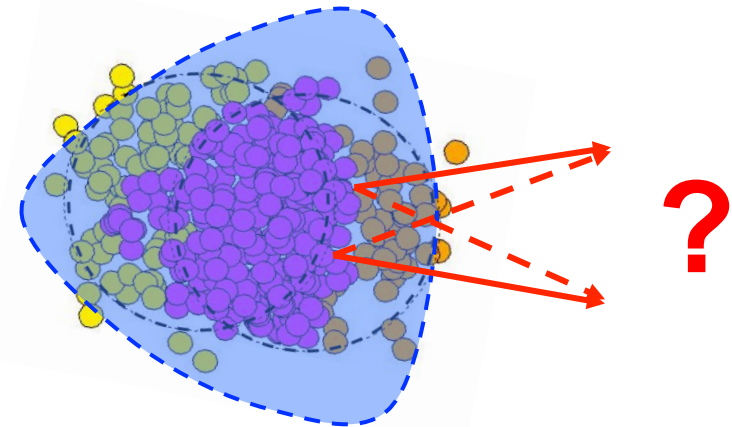
✧ Supporting longer emission duration of K from R_o/R_s

Azimuthal HBT w.r.t 3rd order event plane

Initial spatial fluctuation

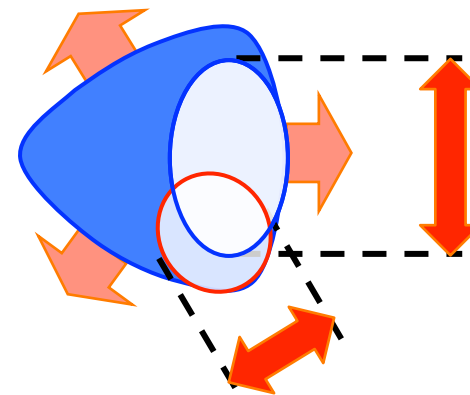
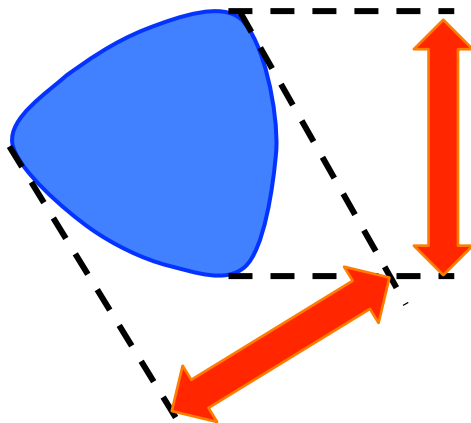


What is final shape ?

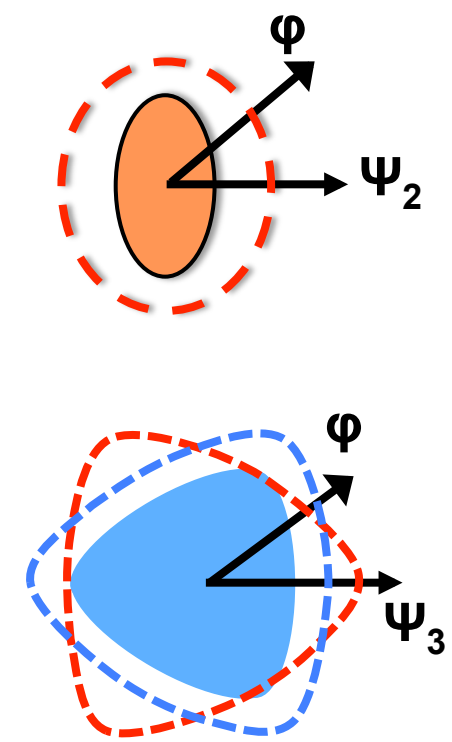
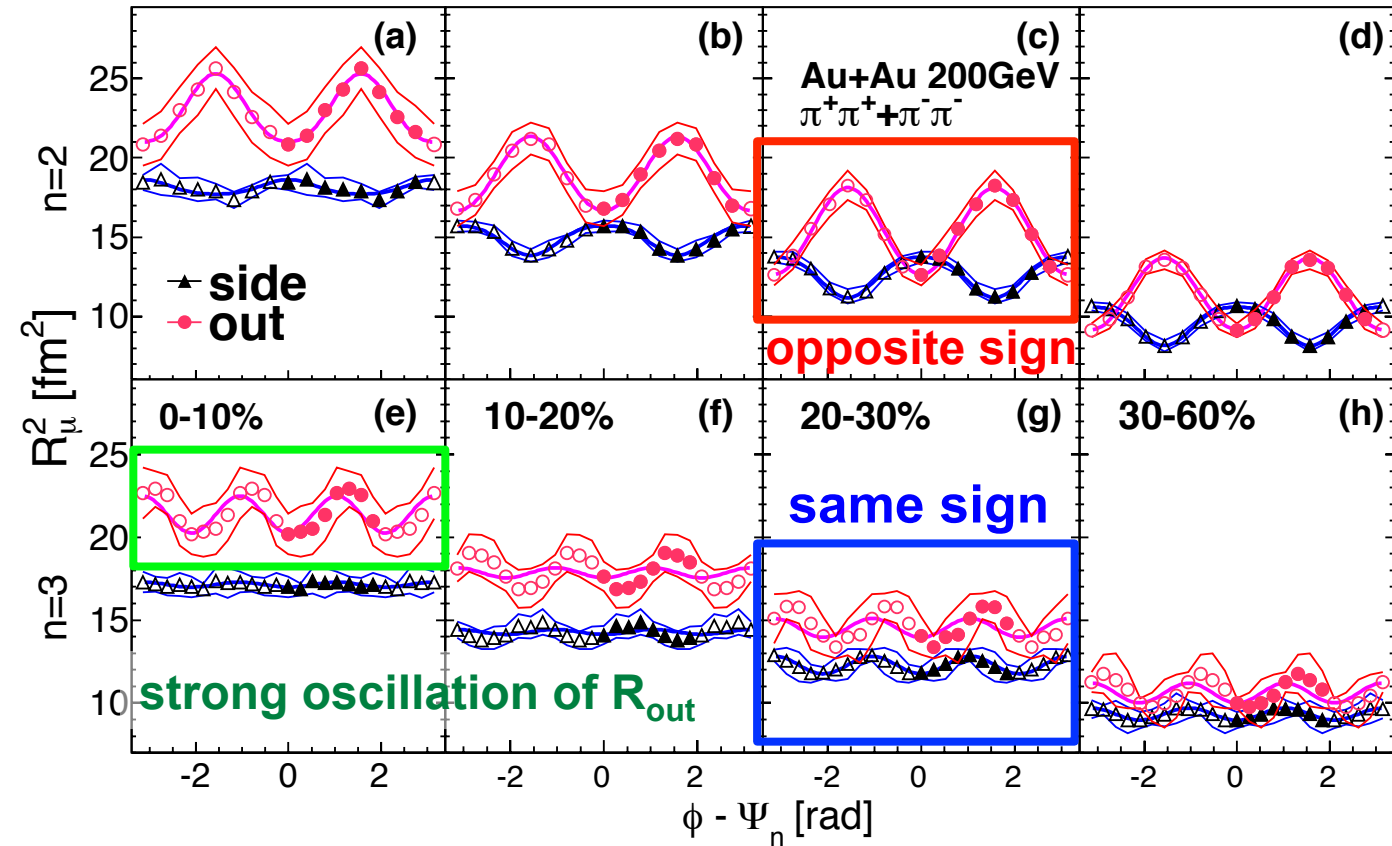


- Note that anisotropy is not observed in a static source.

measured size \approx extent of base of triangle

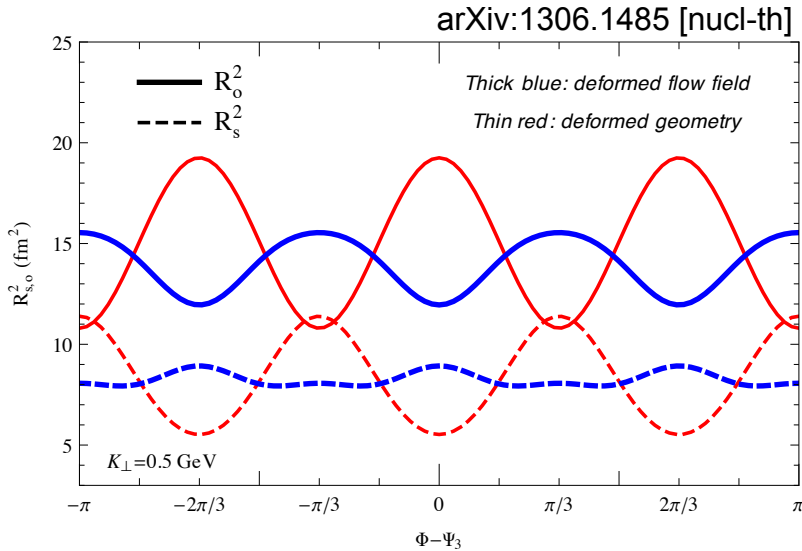


Pion HBT radii w.r.t Ψ_3



- In 0-10%, strong oscillation of R_o is seen as well as 2nd-order
- R_s slightly shows the same sign as R_o in 20-30% unlike 2nd-order

Possible explanation



- Qualitatively agreement with R_o oscillation and R_s flatness

- Comparison of m_T dependence

(Model calculations are scaled by 0.3)

- ✧ Similar trend of m_T dependence of $R_{o,3}^2$

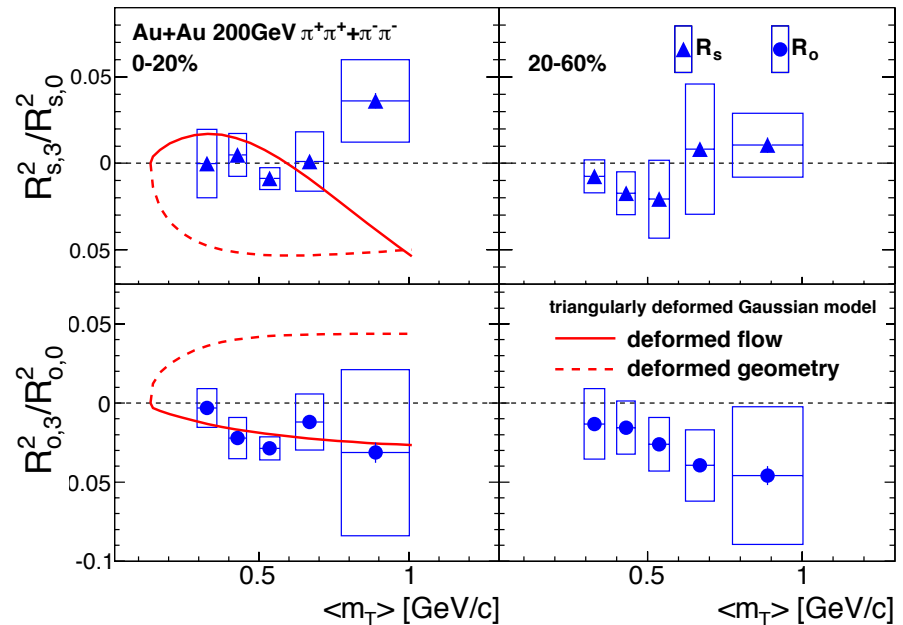
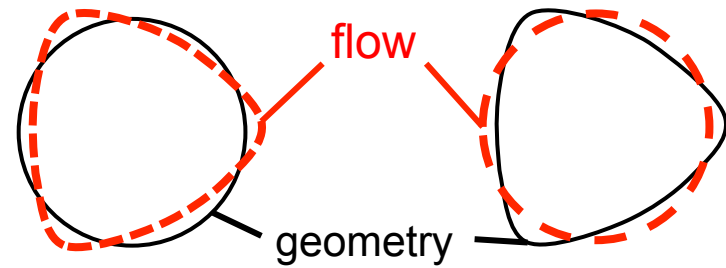
- ✧ $R_{s,3}^2$ seems to be opposite trend.

- ✓ negative value at low m_T , and goes up to positive value at higher m_T

- Gaussian model w/ and w/o deformed flow/geometry

deformed flow

deformed geometry



Interpretation by Monte-Carlo simulation

What does the negative sign of R_s (and the same sign as R_o) mean?
Need to disentangle both effects of the 3rd-order spatial and flow anisotropy!

■ Setup of simulation

- ✧ Similar to BW: thermal motion + transverse boost
- ✧ Spatial distribution
 - ✓ Assuming Woods-Saxon distribution
 - ✓ Spatial shape controlled by “ \mathbf{e}_3 ”
- ✧ Transverse flow
 - ✓ Radial flow with velocity β_0
 - ✓ Flow anisotropy controlled by “ β_3 ”
 - ✓ Boost to radial direction
- ✧ HBT correlation: $1 + \cos(\Delta \mathbf{r} \cdot \Delta \mathbf{p})$
- ✧ No Coulomb interaction, no opacity

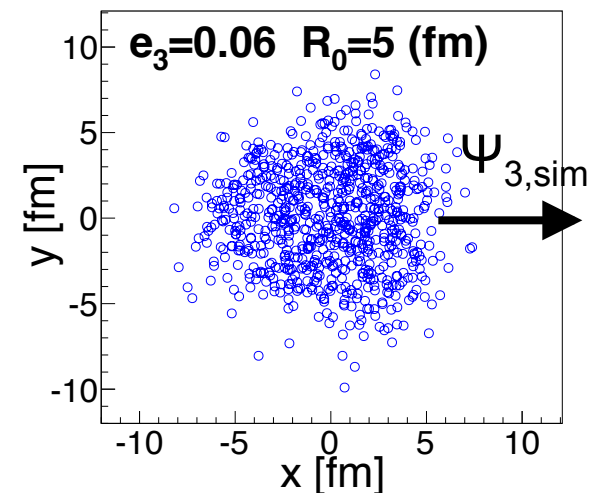
■ Other parameters

- ✧ thermal temperature T_f
- ✧ source size R_0
- ✧ strength of radial flow β_0
- ✧ emission duration $\Delta \tau$
- Tuned by pT spectra, average HBT radii

$$R = R_0(1 - e_3 \cos(3\Delta\phi))$$

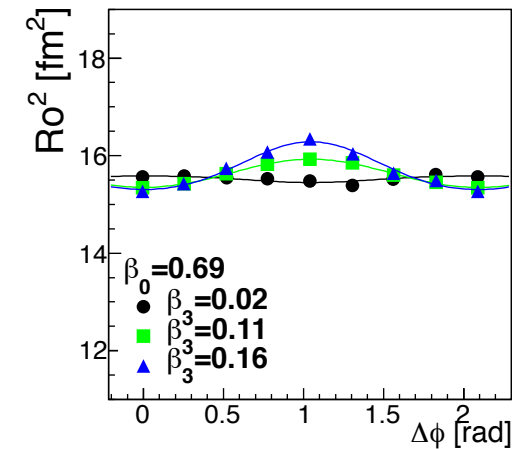
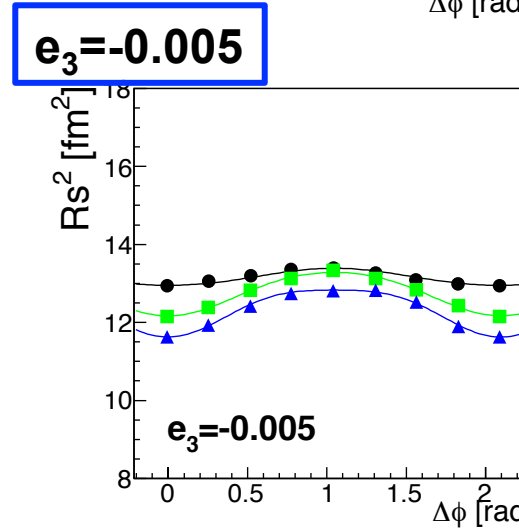
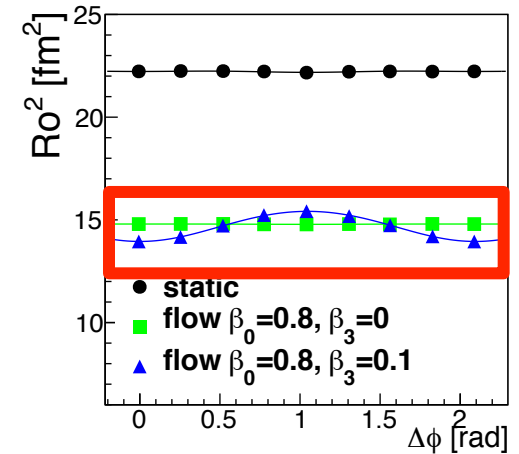
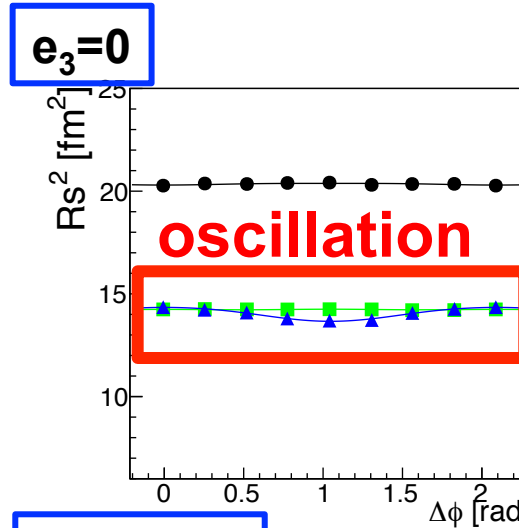
$$\beta_T = \tanh(\rho)$$

$$\rho = \tanh^{-1}[\beta_0 + \beta_3 \cos(3\Delta\phi)]\left(\frac{r}{R}\right)$$



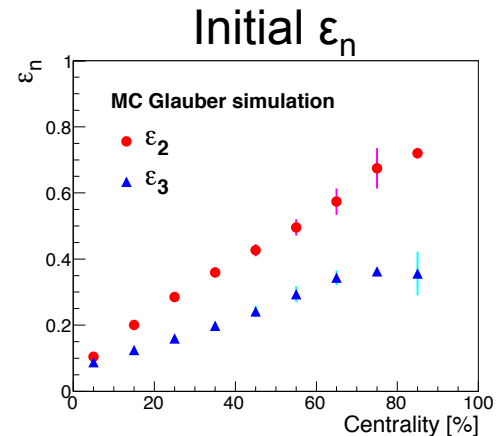
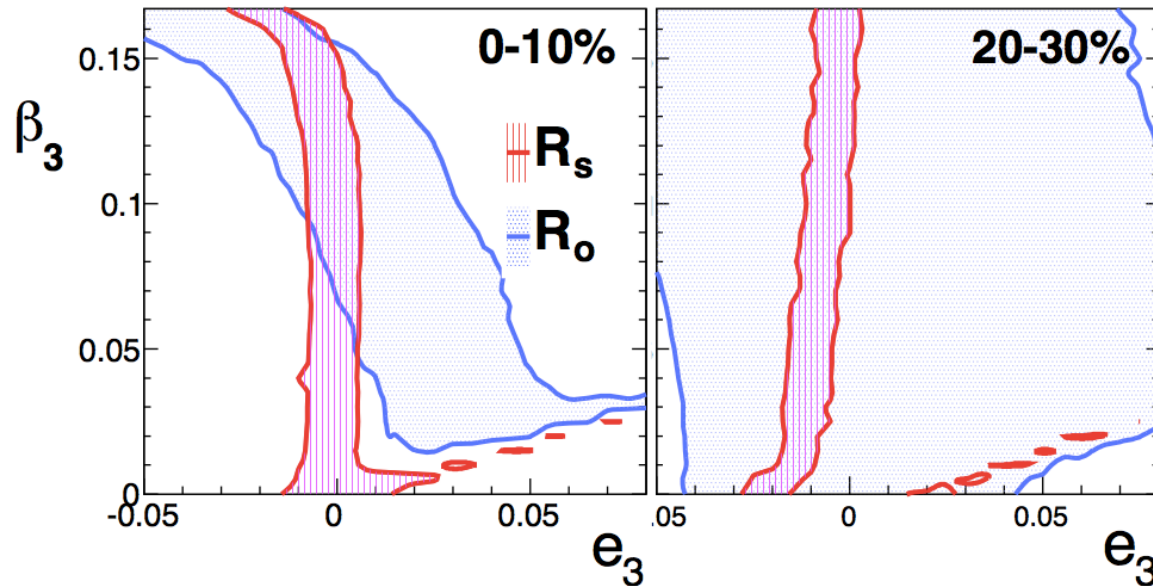
Oscillation of HBT radii in simulation

- Triangular flow makes oscillation even for spherical source.
- Amplitude and sign of the oscillation will be determined by the balance of the spatial and flow anisotropy.

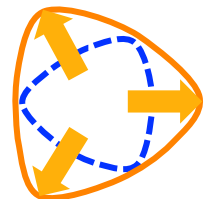


χ^2 minimization for e_3 and β_3

- χ^2 is calculated by difference of relative amplitudes of R_s^2 and R_o^2
- ✧ 1σ -contour lines are shown



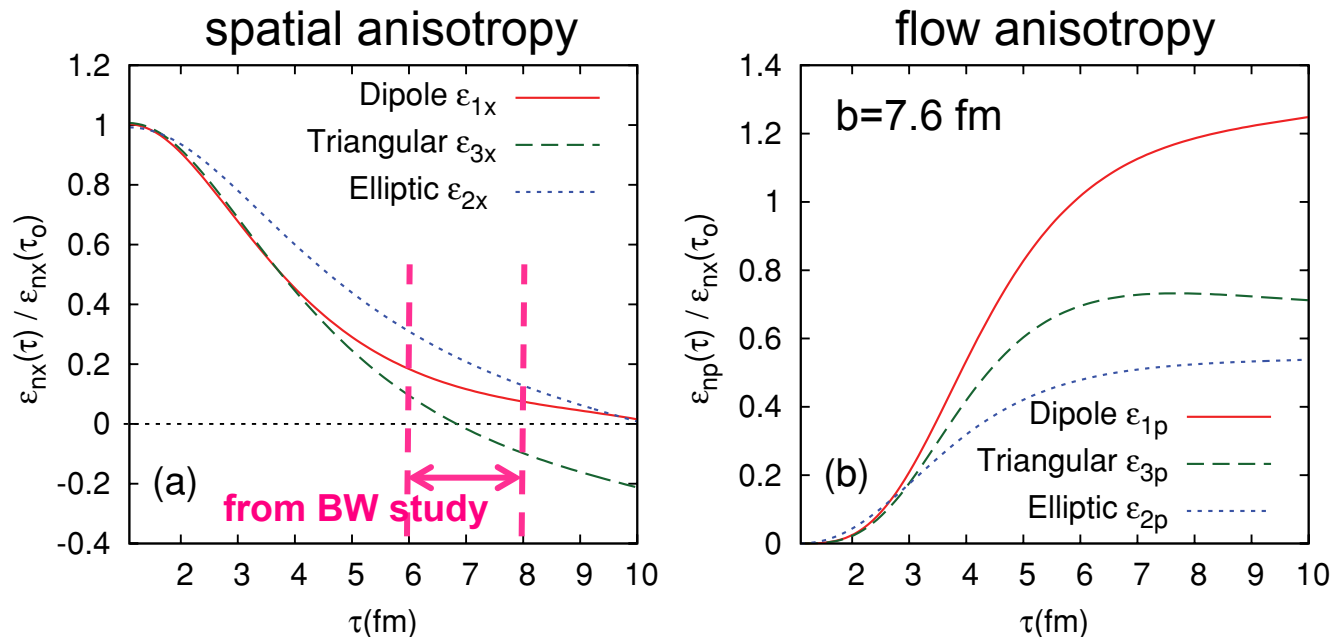
- e_3 is well constrained by R_s , (R_s is less affected by β_3)
- ✧ e_3 is close to zero in 0-10%, and slightly negative in 20-30%
- ✧ That indicates that initial ϵ_3 is much reduced, and potentially reversed by triangular flow under this simulation



Theoretical expectation for final triangularity

■ Initial Glauber for event-by-event cumulants + ideal hydrodynamics (PRC83.064904 (2011))

- ✧ For $\tau > 7$ at fm/c, only ε_3 shows negative value.
- ✧ This model also indicate the possibility of reversed triangularity at freeze-out



Summary-2

- **3rd-order oscillations of HBT radii were observed.**

- ✧ First measurement!
- ✧ Finite R_0 oscillation in all centralities
- ✧ Very small R_s oscillation in mid central

- **Comparison with Gaussian model with deformed flow/geometry**

- ✧ m_T dependence of R_0 oscillation can be qualitatively explained.
- ✧ R_0 oscillation will be mainly driven by triangular flow.

- **Monte-Carlo simulation**

- ✧ R_s oscillation is affected by geometry, not so much by flow.
- ✧ χ^2 minimization was performed
 - ✓ geometrical anisotropy e_3 shows zero ~ slightly negative

Conclusions

Azimuthal angle dependence of HBT radii was measured w.r.t 2nd and 3rd-order event plane in Au+Au collisions at $\sqrt{s_{NN}} = 200$ GeV

■ Final spatial distribution

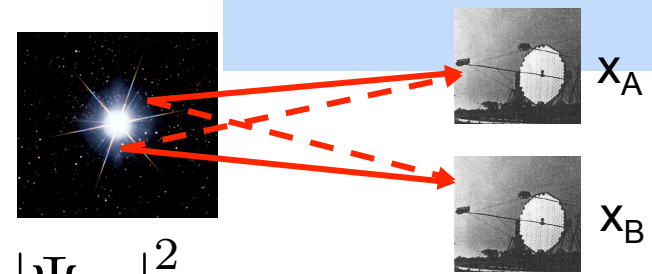
- ✧ Initial eccentricity is diluted, but still out-of-plane extended shape.
- ✧ Initial triangularity is significantly diluted, and potentially reversed.

■ Particle species

- ✧ Freeze-out mechanism may be different between pions and kaons.
- ✧ Result indicates longer emission duration of kaons, which is opposite to an intuitive expectation in terms of the cross section.

Back up

Correlation function C_2 is defined as:



$$C_2 = \frac{P_{12}}{P_1 P_2} = \frac{\int dx_1 dx_2 \rho(x_1) \rho(x_2) |\Psi_{12}|^2}{\int dx_1 \rho(x_1) |\Psi_1|^2 \int dx_2 \rho(x_2) |\Psi_2|^2}$$

where P_{12} is a possibility to measure two particles, and $P_1(P_2)$ is a possibility to measure one particle. $\rho(r)$ is the spatial distribution of particle emitting source. Assuming the plane wave for the wave function, C_2 is rewritten as:

$$\begin{aligned} P_{12} &\propto [e^{-ip_1 \cdot (x_A - x_1)} e^{-ip_2 \cdot (x_B - x_2)} + e^{-ip_2 \cdot (x_A - x_2)} e^{-ip_1 \cdot (x_B - x_1)}]^2 \\ &= [2 \pm e^{-ix_1(p_2 - p_1)} e^{ix_2(p_2 - p_1)} \pm e^{ix_1(p_2 - p_1)} e^{-ix_2(p_2 - p_1)}] \\ &\propto 1 \pm \cos[(x_2 - x_1) \cdot (p_2 - p_1)] \end{aligned}$$

correlation term

$$\begin{aligned} C_2 &= 1 \pm \left| \int dx \rho(x) e^{ix(p_2 - p_1)} \right|^2 \\ &= 1 \pm |\tilde{\rho}(q)|^2 \end{aligned}$$

$\tilde{\rho}(q)$: Fourier transform of $\rho(r)$
 $q = p_2 - p_1$

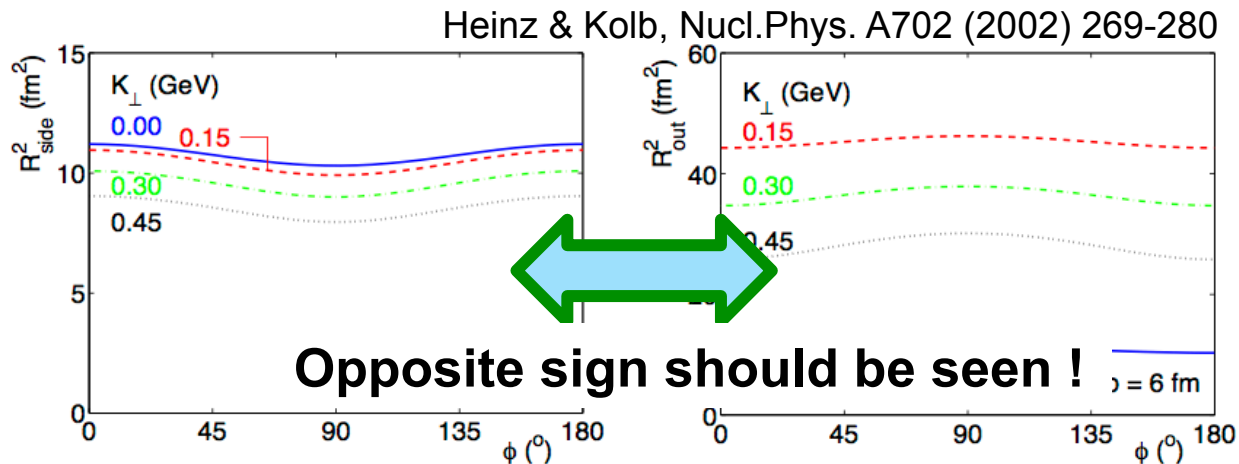
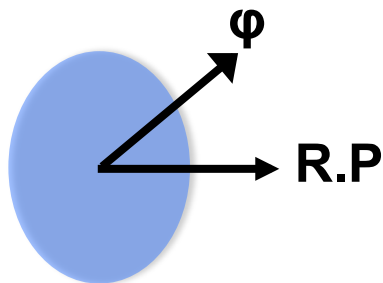
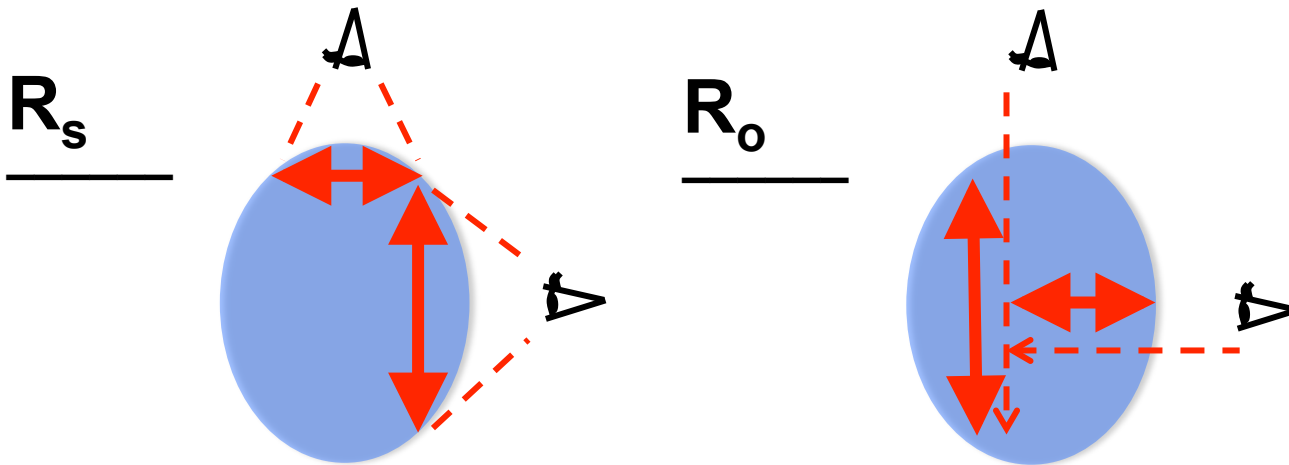
C_2 can be expressed as a function of relative momentum!

What do R_s and R_o represent ?

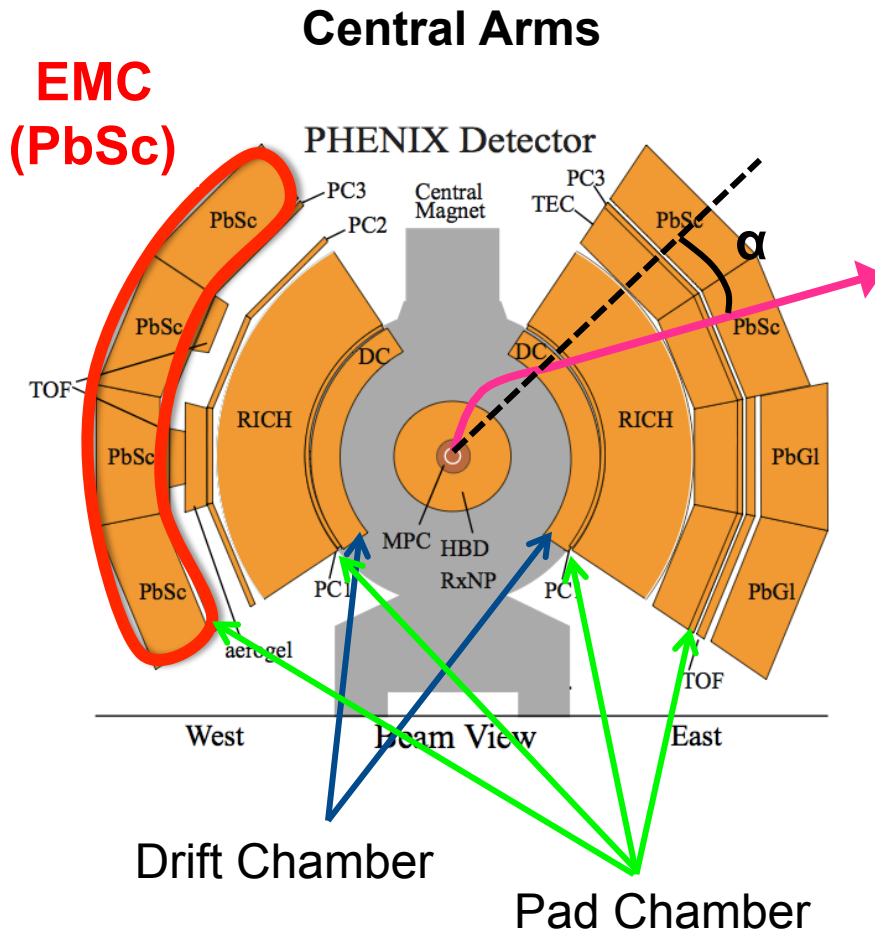
- R_s is “width” of source
- R_o is “thickness” of source

$$\vec{k}_T = \frac{1}{2}(\vec{p}_{T1} + \vec{p}_{T2})$$

$$\vec{q}_{out} \parallel \vec{k}_T, \vec{q}_{side} \perp \vec{k}_T$$



Track Reconstruction



■ Drift Chamber

- ✧ Momentum determination

$$p_T \simeq \frac{K}{\alpha} \quad \begin{array}{l} K: \text{field integral} \\ \alpha: \text{incident angle} \end{array}$$

■ Pad Chamber (PC1)

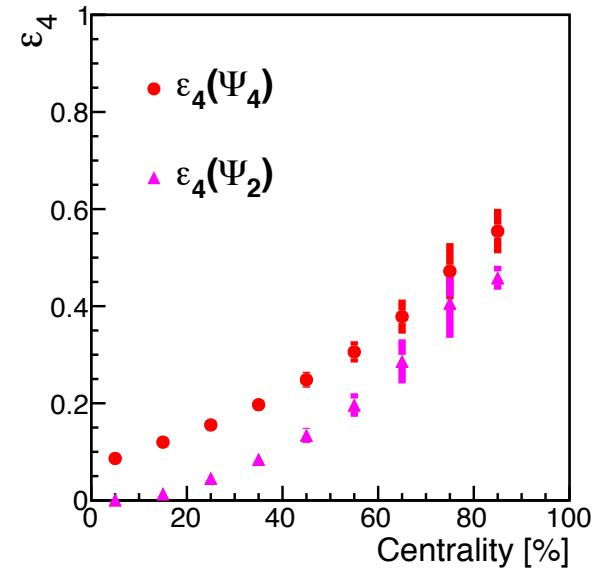
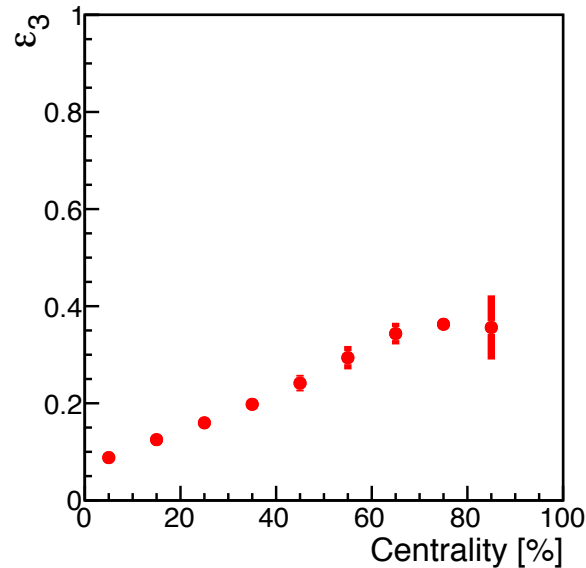
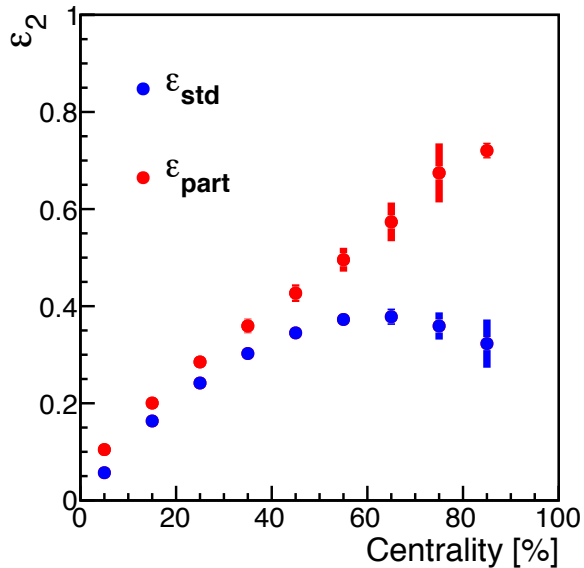
- ✧ Associate DC tracks with hit positions on PC1
 - ✓ p_z is determined

■ Outer detectors (PC3, TOF, EMCal)

- ✧ Extend the tracks to outer detectors

Initial spatial anisotropy by Glauber model

- Initial eccentricity and triangularity increase with centrality going from central to peripheral.



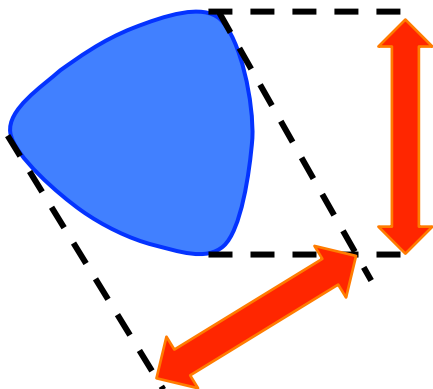
Simulation check

■ Static source

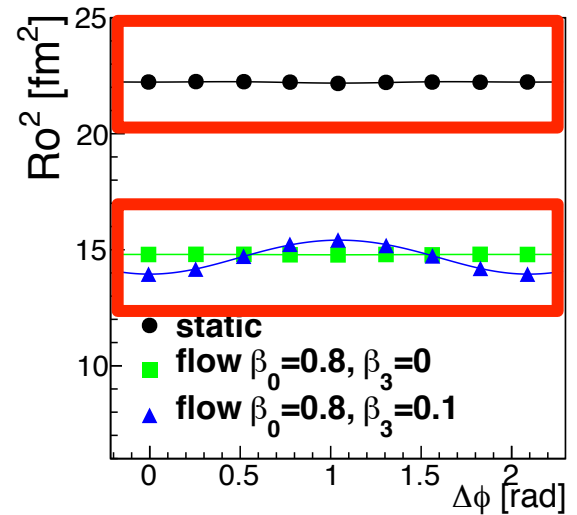
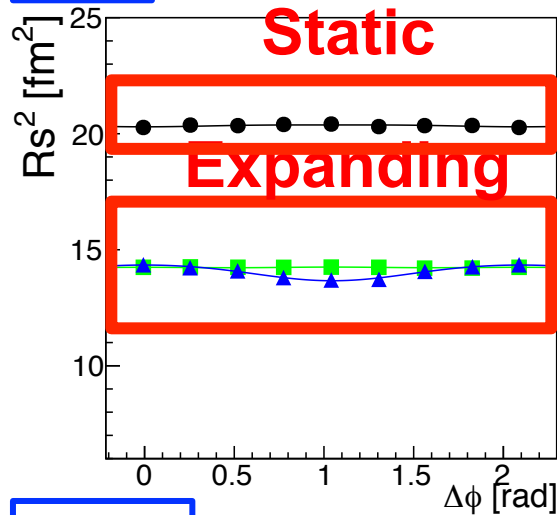
- ✧ No oscillation for triangular shape

■ Expanding source

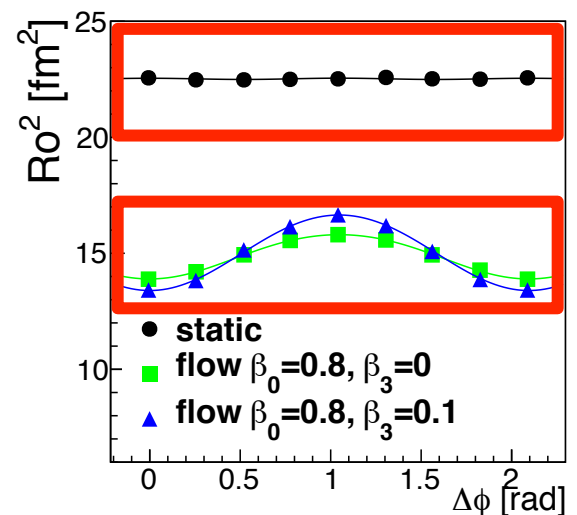
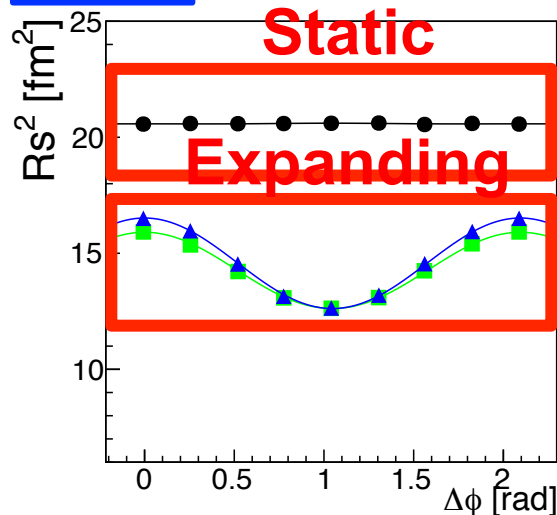
- ✧ Triangular shape
 - ✓ Oscillation appears !
- ✧ Spherical shape
 - ✓ β_3 makes oscillation !



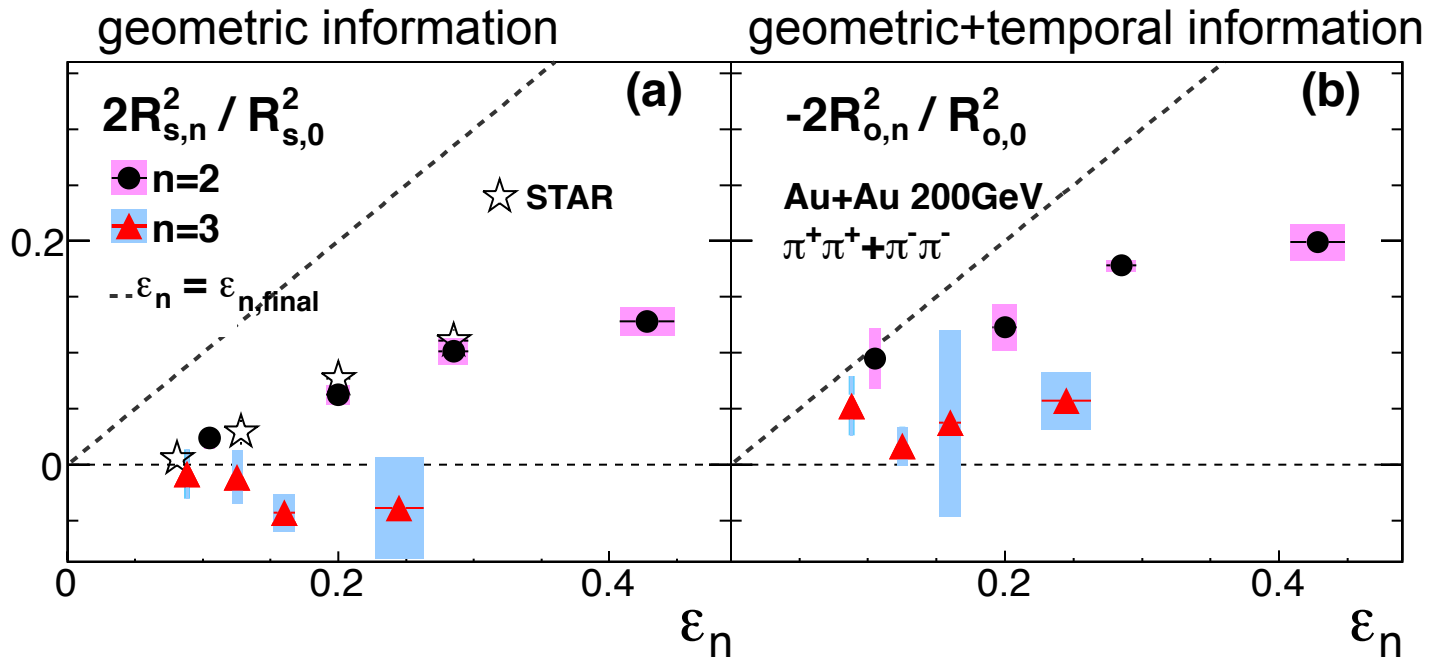
$e_3=0$



$e_3=0.1$



Initial vs Final source eccentricity



■ 2nd-order oscillation is sensitive to final eccentricity

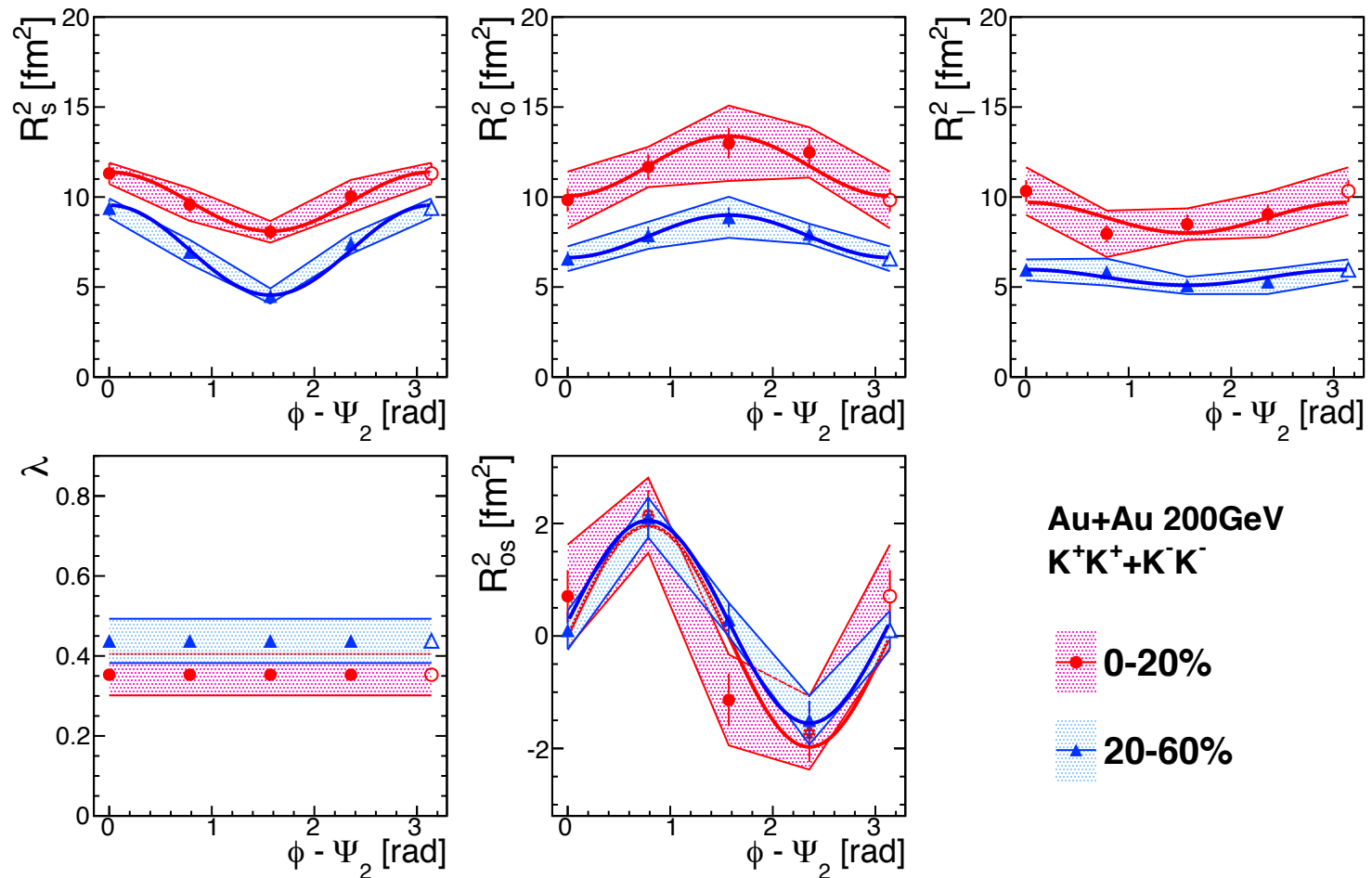
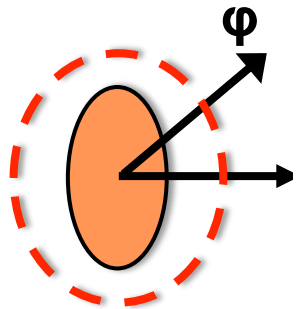
- ✧ final $\epsilon_2 \approx$ initial ϵ_2
- ✧ Strong expansion to in-plane direction, and instant emission

■ 3rd-order oscillation affected by spatial and flow anisotropy

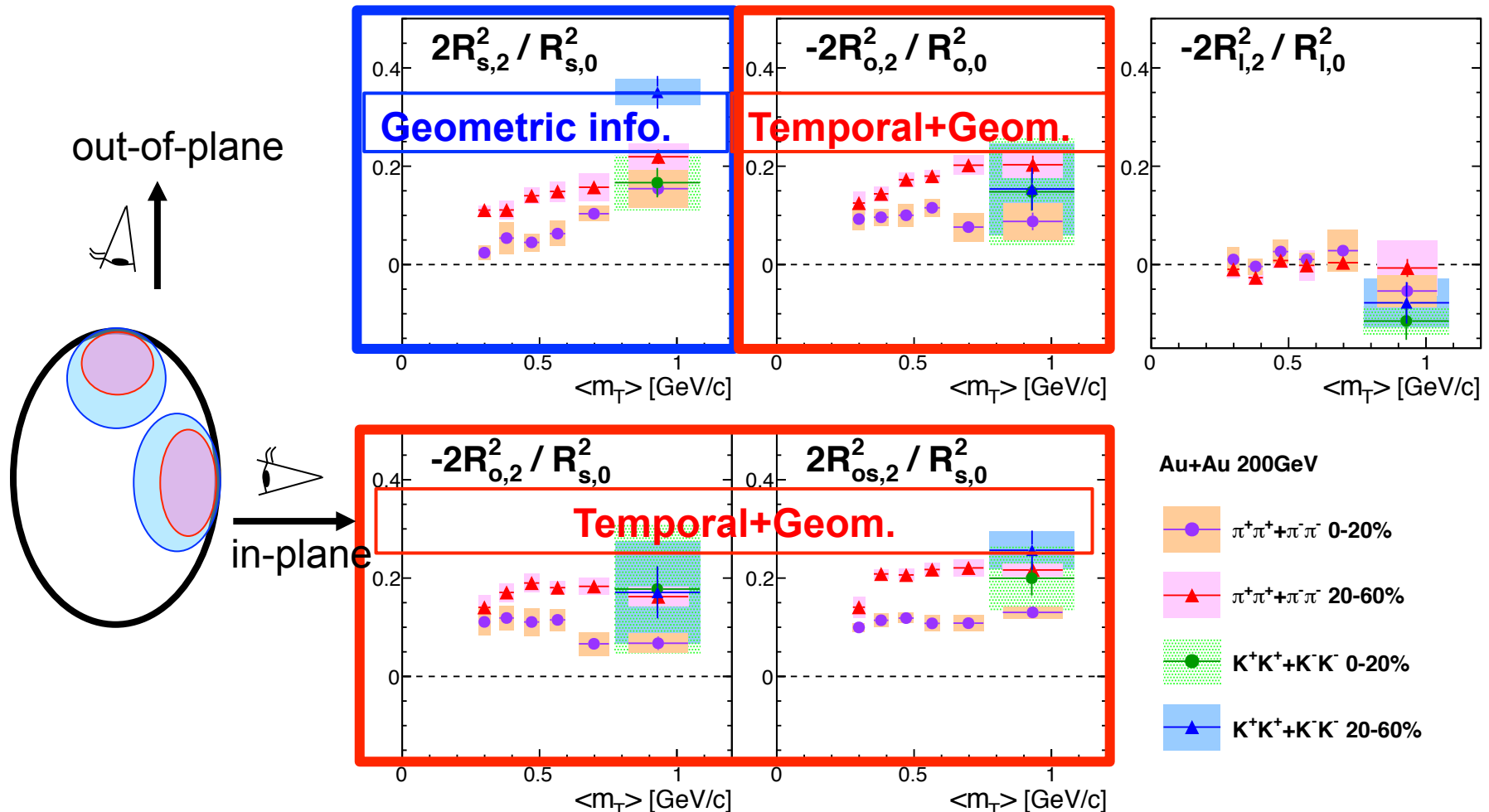
- ✧ Smaller than 2nd-order
- ✧ Initial triangular shape may be reduced by triangular flow

Centrality dependence of kaon HBT radii w.r.t Ψ_2

- Result of charged kaons show similar trends!



m_T dependence of relative amplitude



- Relative amplitude of R_{side} and R_{out} show m_T dependence

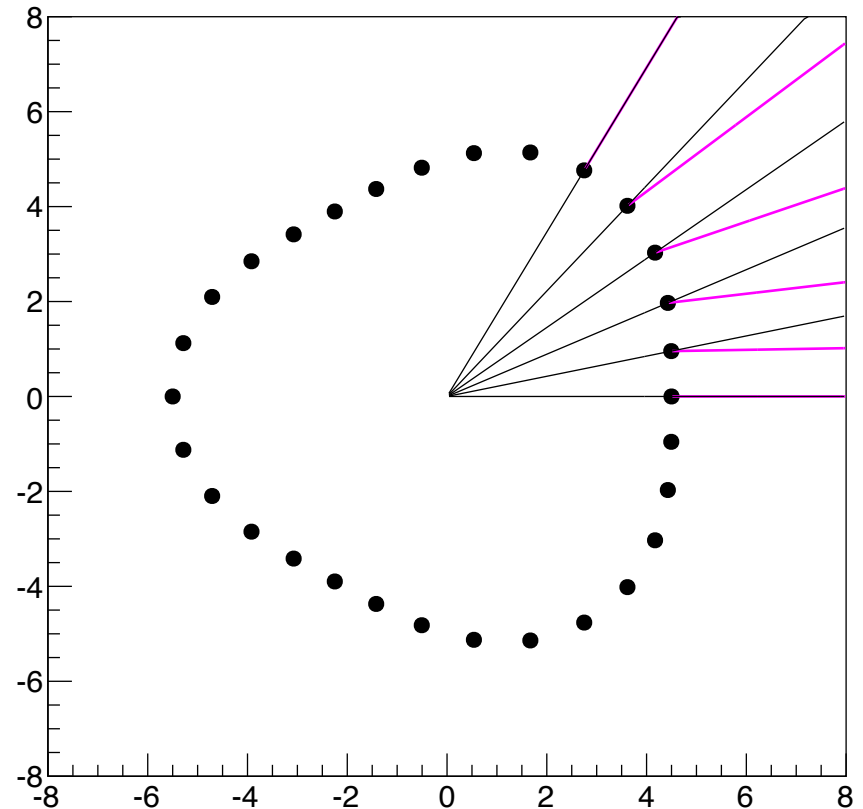
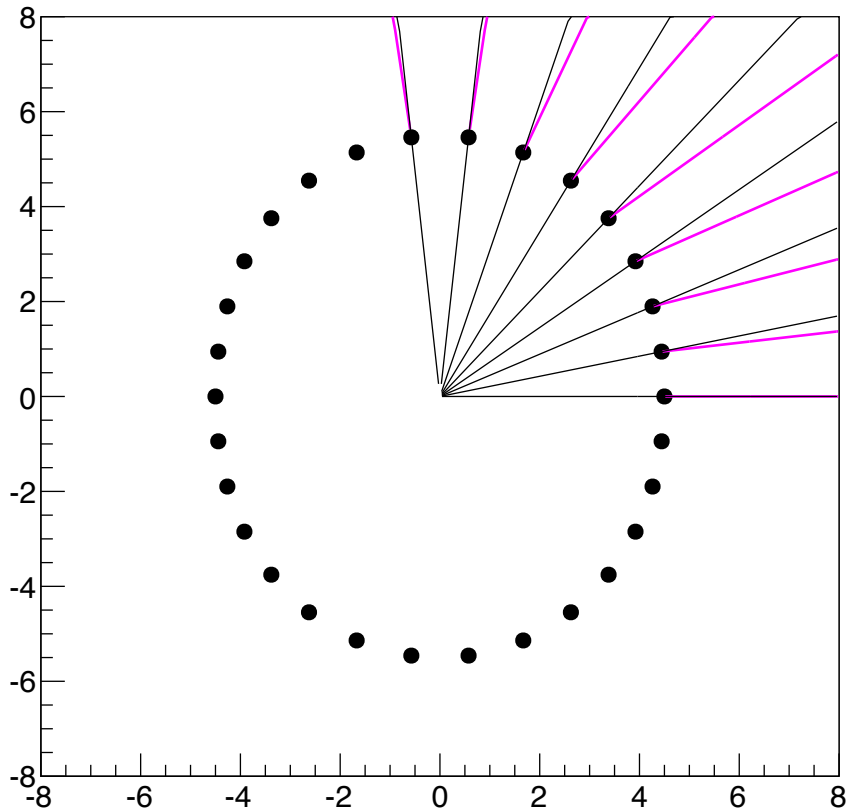
✧ R_o in 0-20% doesn't depend on m_T ?

- Difference between π/K is similar for other parameters

Boost angle

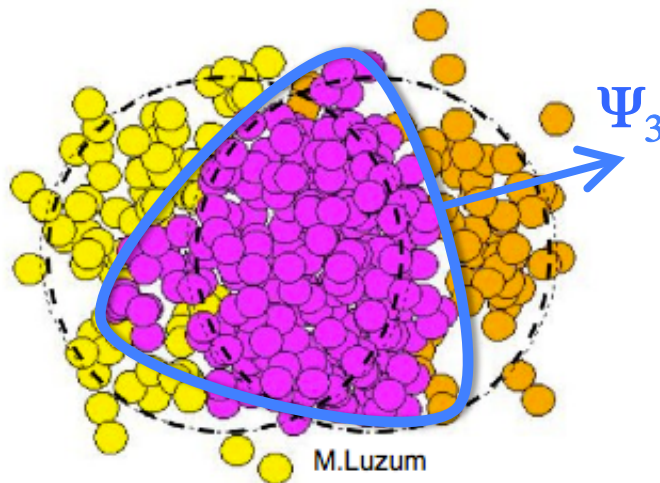
■ Boost angle is set to be :

- ✧ radial direction of the particle position (**radial boost**)
- ✧ perpendicular to the surface which is similar condition with BW (PRC70, 044907 (2004)) (**surface boost**)



HBT vs Higher Harmonic Event Plane

- The idea is to expand azimuthal HBT to higher harmonic event planes.
- ✧ may show the fluctuation of the shape at freeze-out.
- ✧ provide more constraints on theoretical models about the system evolution.



Hydrodynamic model calculation

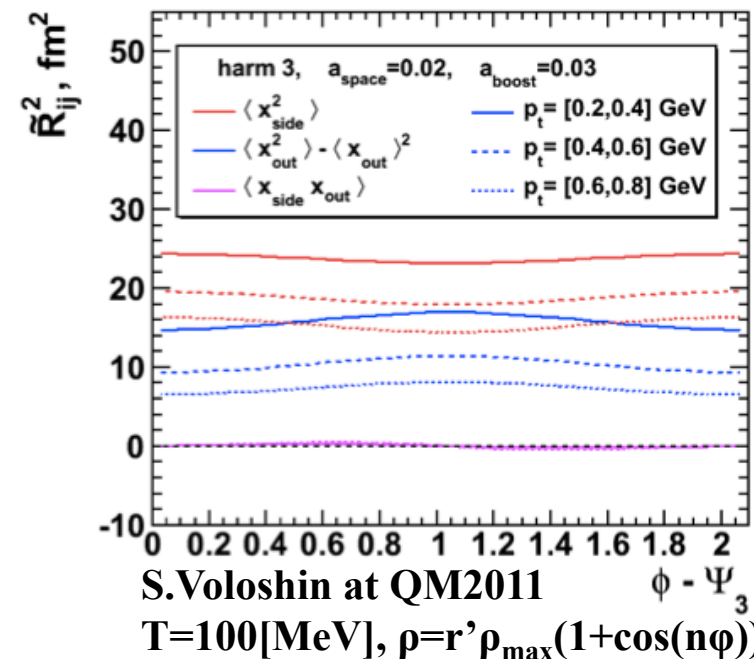
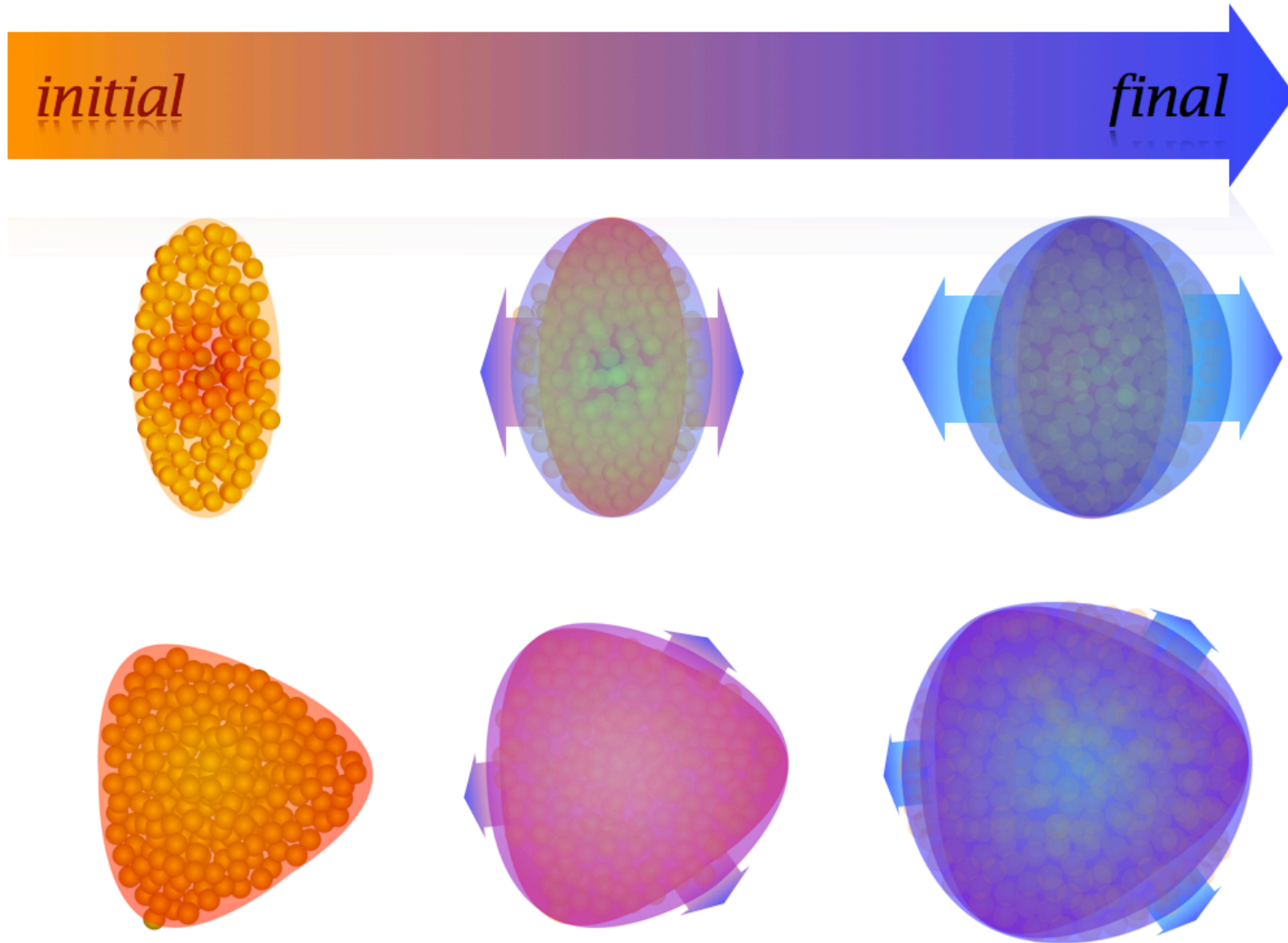
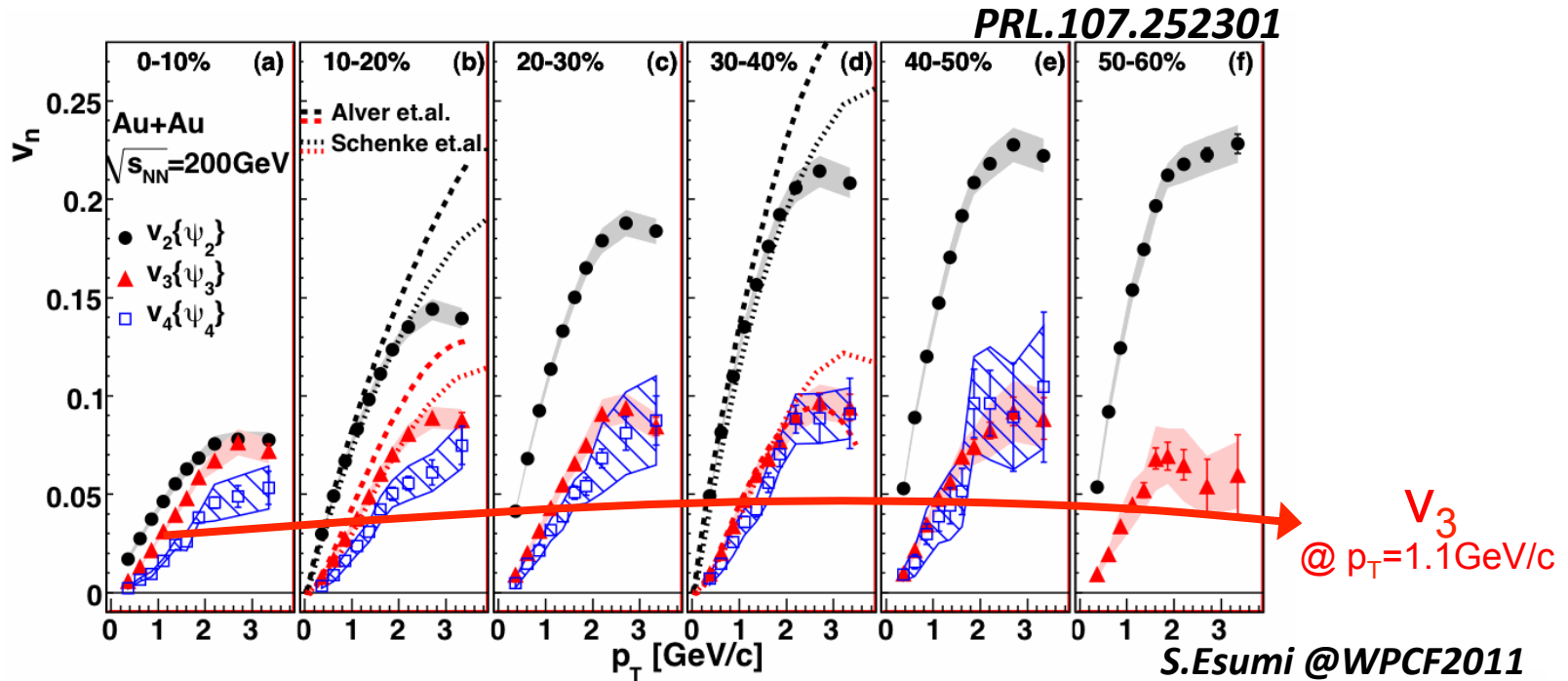


Image of initial/final source shape

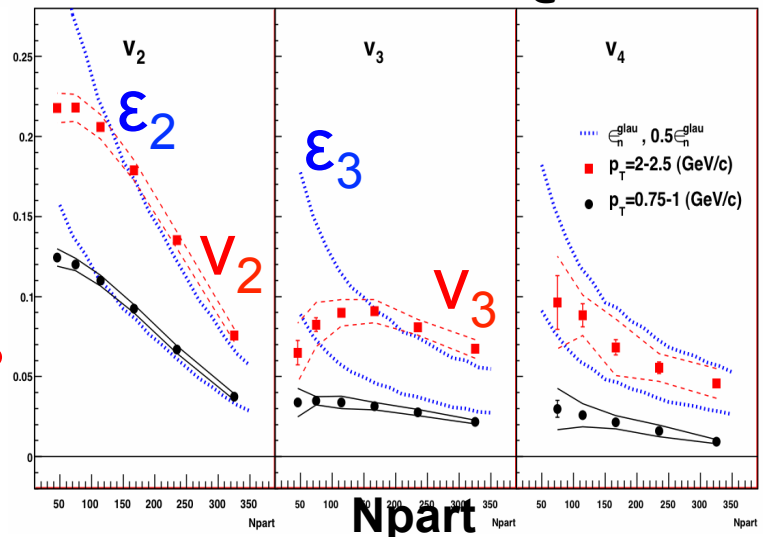


Centrality dependence of v_3 and ϵ_3



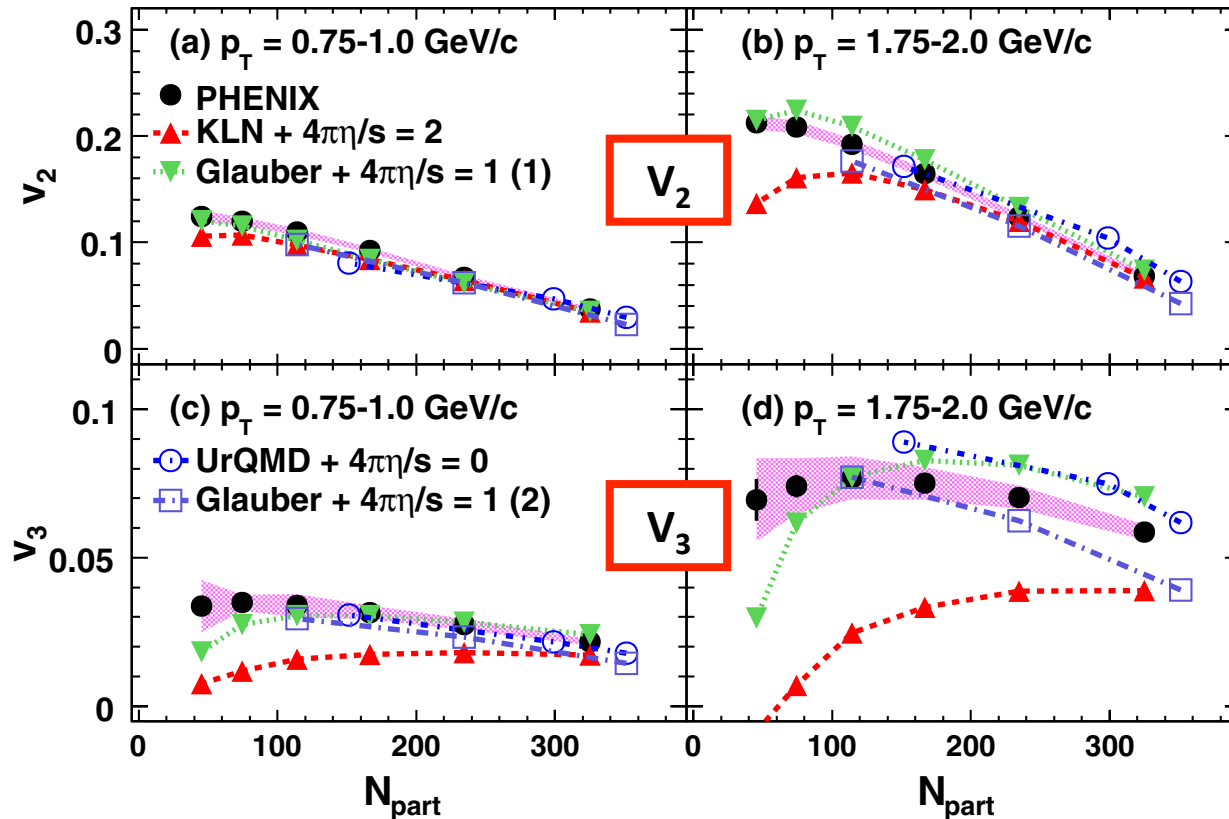
- Weak centrality dependence of v_3
- Initial ϵ_3 has centrality dependence

🍄 Final ϵ_3 has any centrality dependence?



v_3 breaks degeneracy

PRL.107.252301



- v_3 provides new constraint on hydro-model parameters

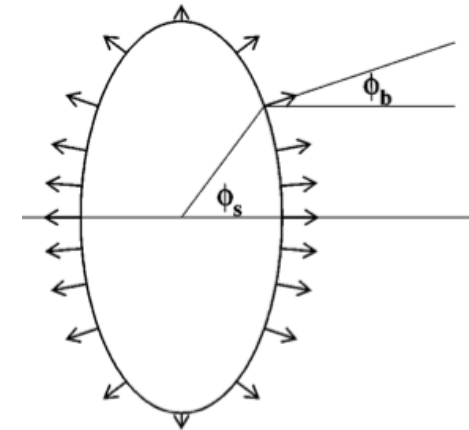
- ✧ Glauber & $4\pi\eta/s=1$: works better
- ✧ KLN & $4\pi\eta/s=2$: fails

Interpretation by Blast wave model

■ Hydrodynamic model assuming radial flow

- ✧ Well described p_T spectra & elliptic flow at low p_T
- ✧ Assuming freeze-out takes place for all hadrons at the same time
- ✧ Freeze-out condition is treated as free parameters.

PRC 70, 044907 (2004)



7 free parameters * box profile is assumed as spatial density

T_f : temperature at freeze-out

ρ_0, ρ_2 : transverse rapidity

R_x, R_y : transverse sizes (shape)

$\tau_0, \Delta\tau$: freeze-out time, emission duration

$$\beta_T = \tanh(\rho)$$

$$\rho(r, \phi) = \tilde{r}[\rho_0 + \rho_2 \cos(2\phi)]$$

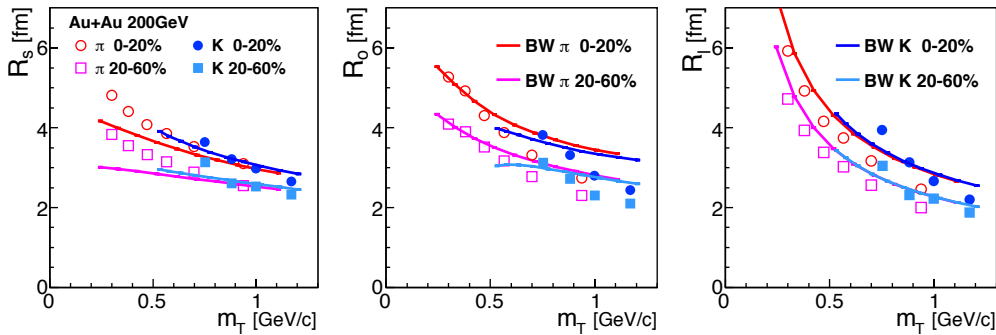
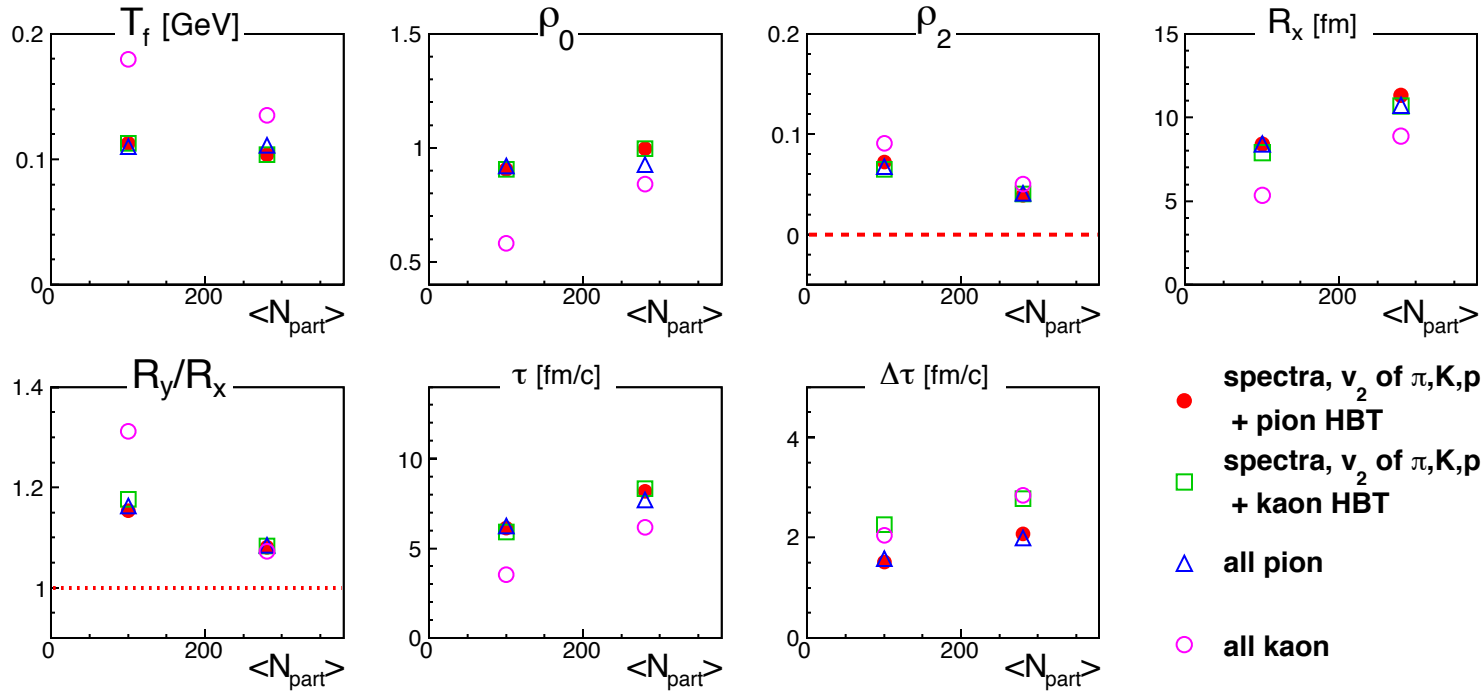
$$\tilde{r} = \sqrt{\left(\frac{r \cos(\phi)}{R_x}\right)^2 + \left(\frac{r \sin(\phi)}{R_y}\right)^2}$$

- ✧ Assuming a Gaussian distribution peaked at τ_0 and with a width $\Delta\tau$, and source size doesn't change with τ .

$$\frac{dN}{d\tau} \sim \exp\left(-\frac{(\tau - \tau_0)^2}{2\Delta\tau^2}\right)$$

- ✧ Spatial(R_y/R_x) and flow(ρ_2) anisotropy make HBT oscillation.

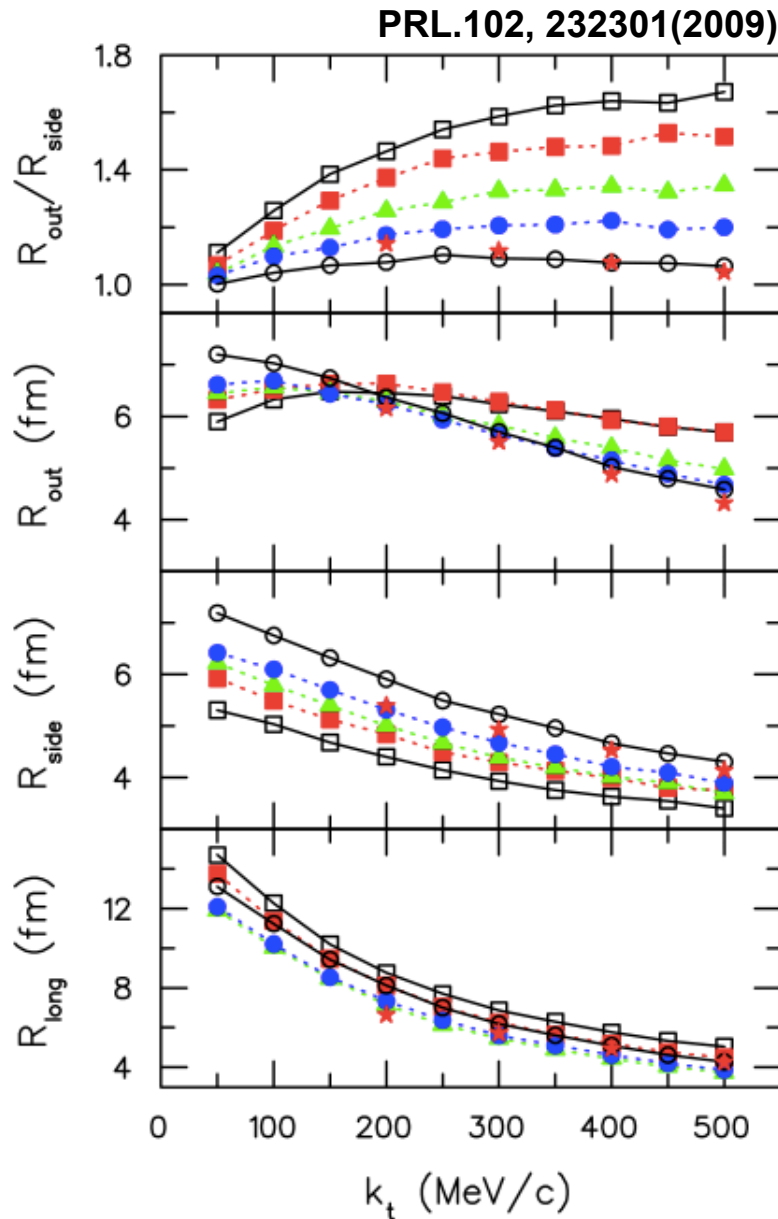
pion BW fit vs kaon BW fit



■ Fit only for kaon data shows longer $\Delta\tau$, while m_T dependence of mean radii is not reproduced.

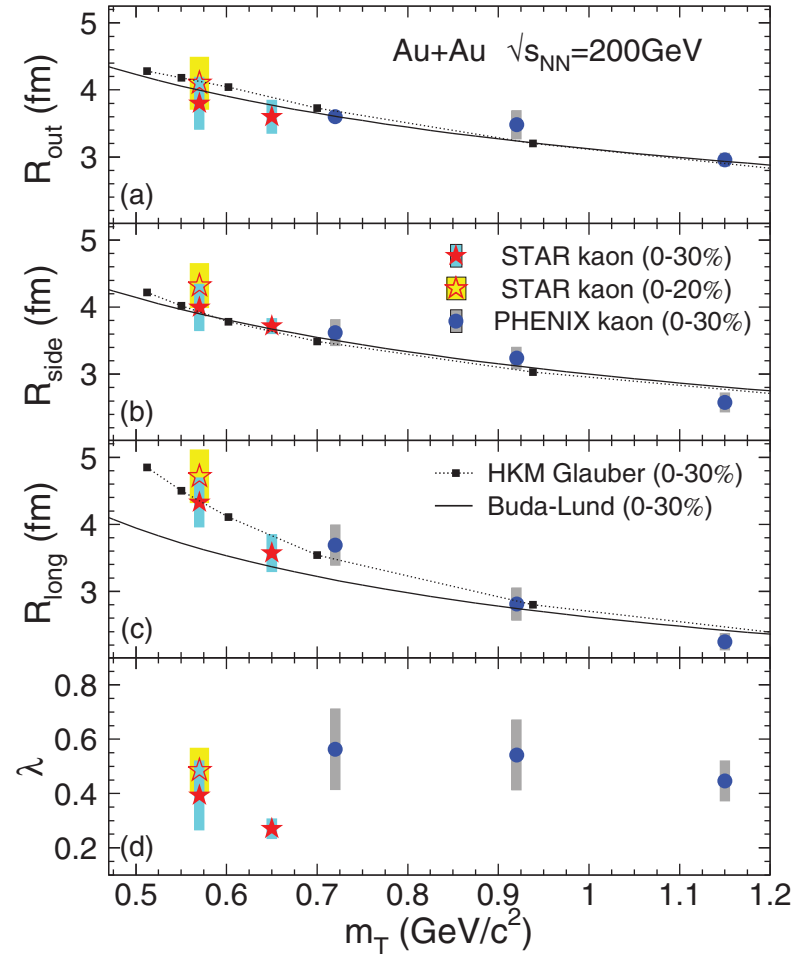
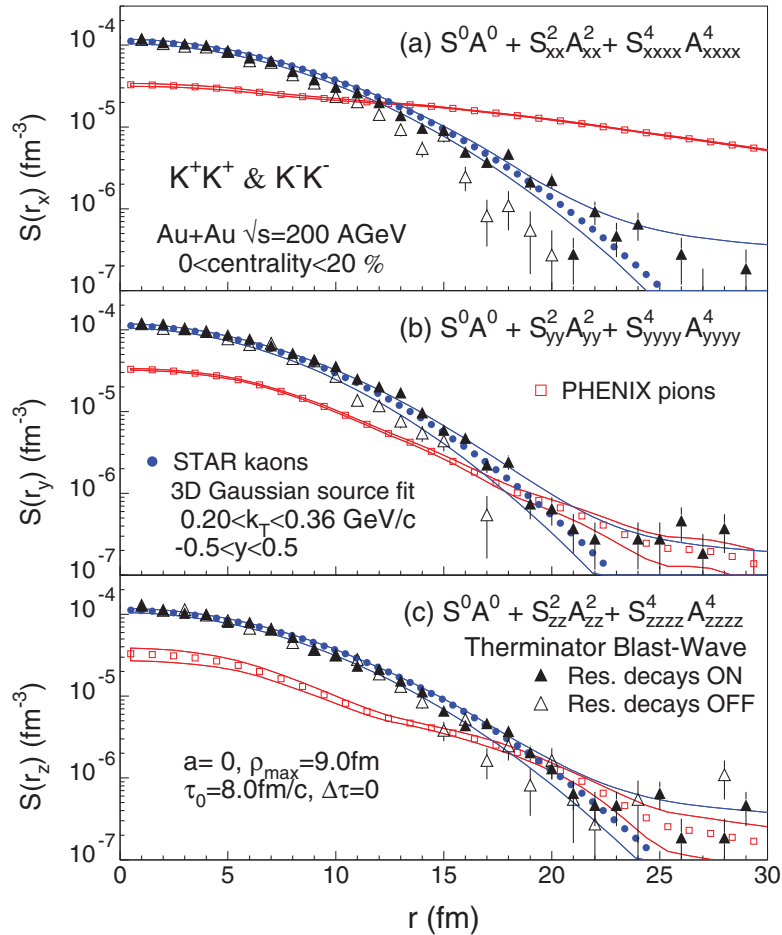
✧ azimuthal kaon data used in the fit has only a data point in k_T .

Comparison of models and the past HBT results



Hydrodynamic model can reproduce the past HBT result!
HBT can provide constraints on the model!

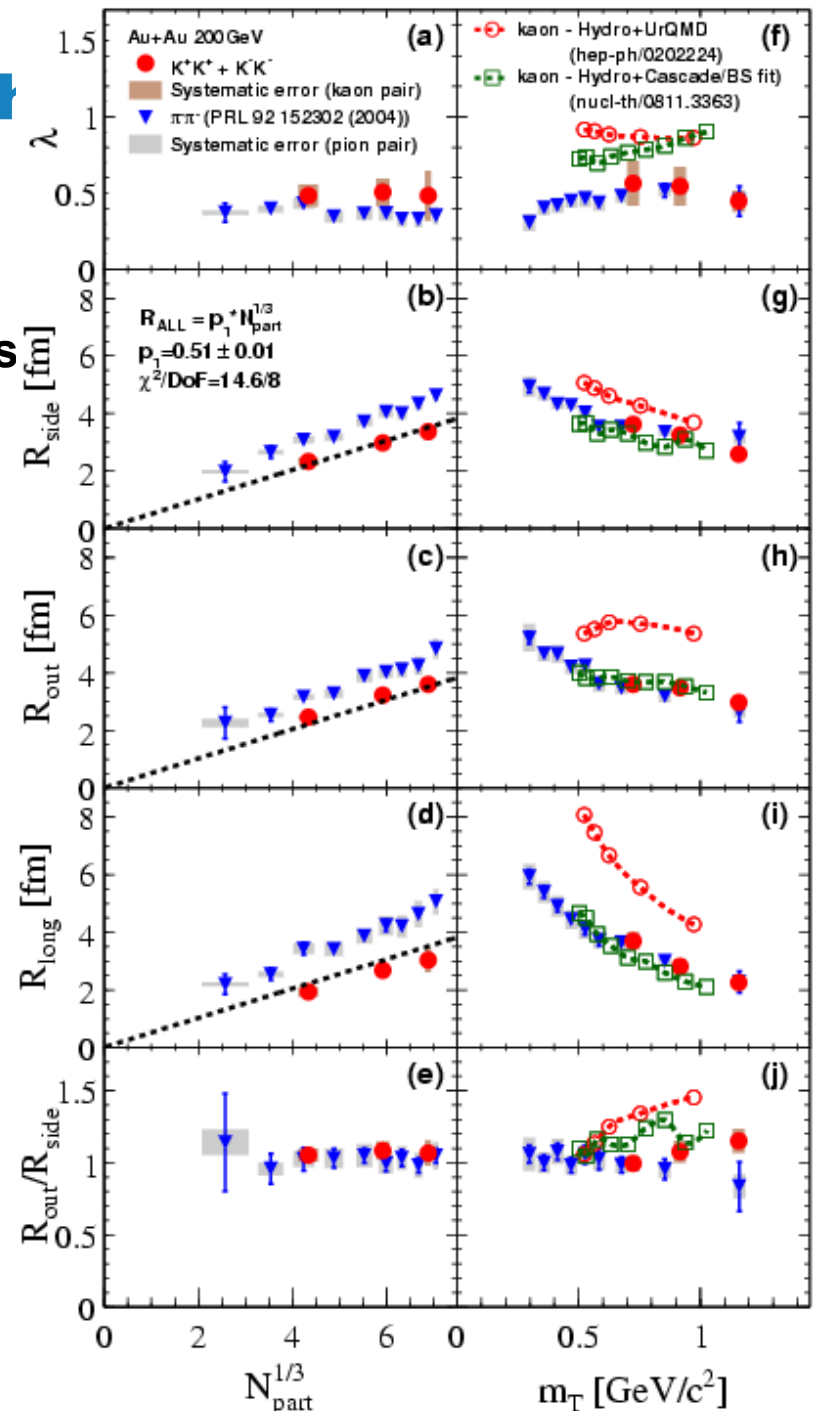
Imaging analysis from STAR



The past HBT Results for cl

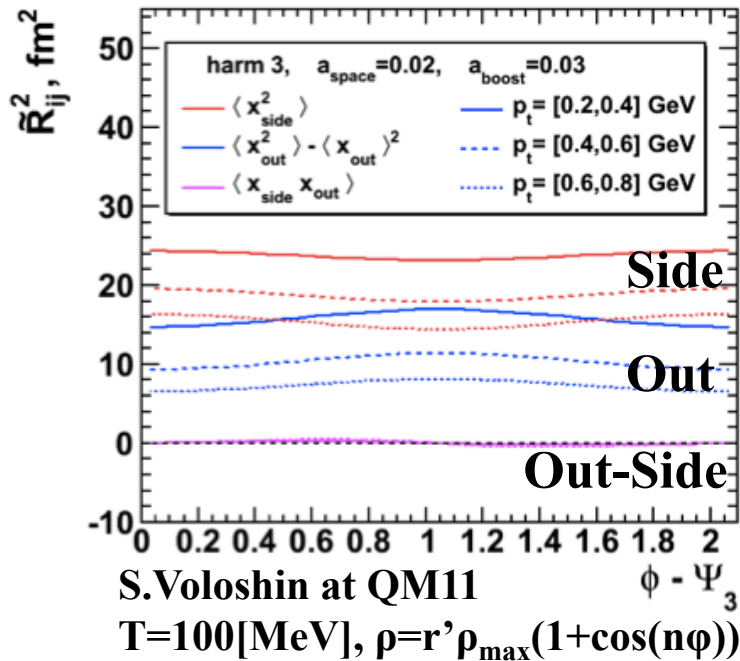
- Centrality / m_T dependence have been measured for pions and kaons

✧ No significant difference between both species

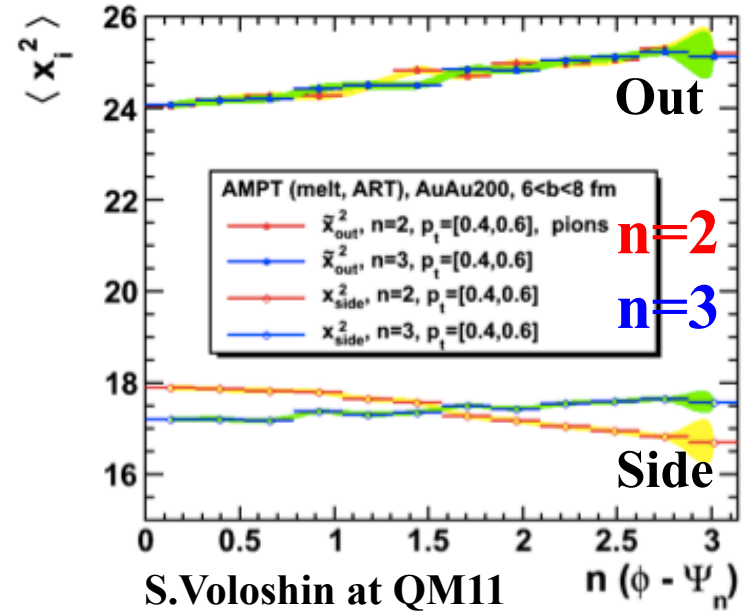


Model predictions

Blast-wave model



AMPT

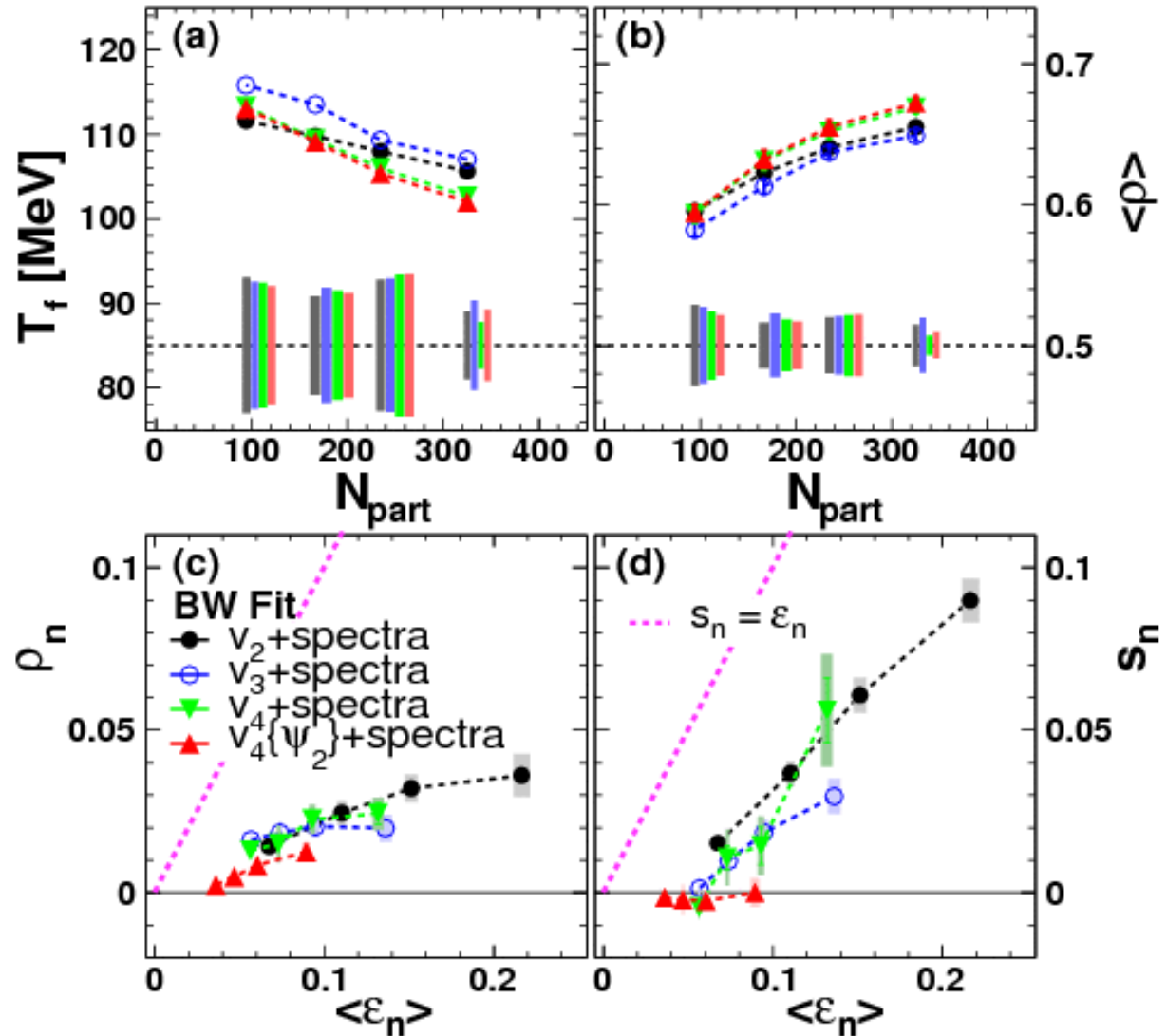


Both models predict weak oscillation will be seen in R_{side} and R_{out} .

Extracted parameters by different Blast wave model

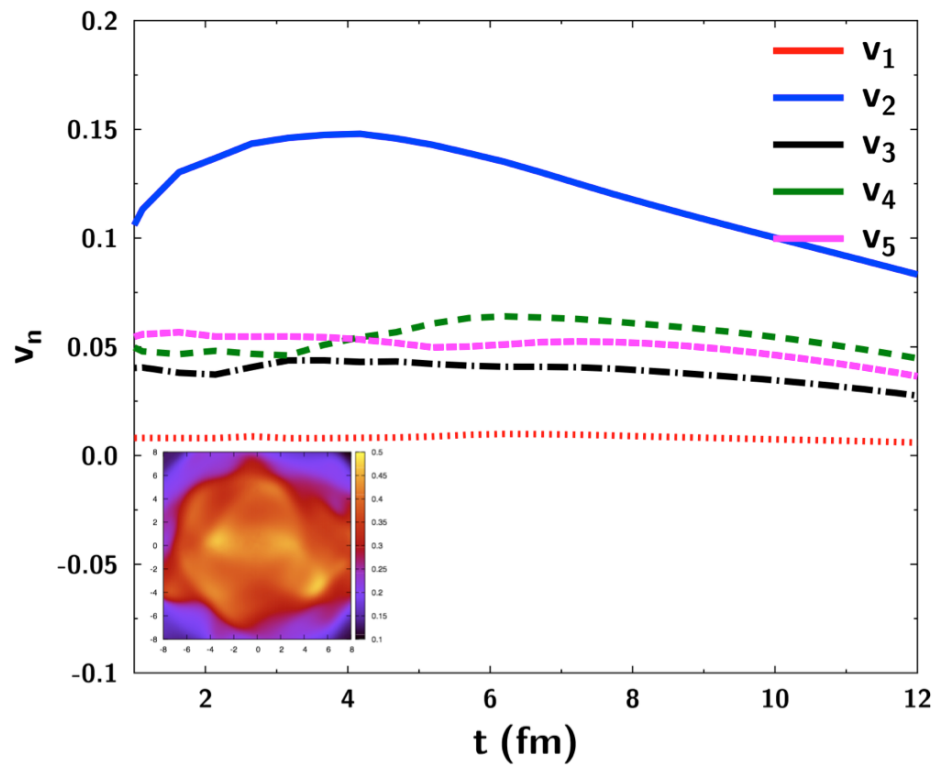
- Blast wave fit was performed for spectra and v_n

Flow anisotropy ρ_2 and ρ_3
have similar trend !



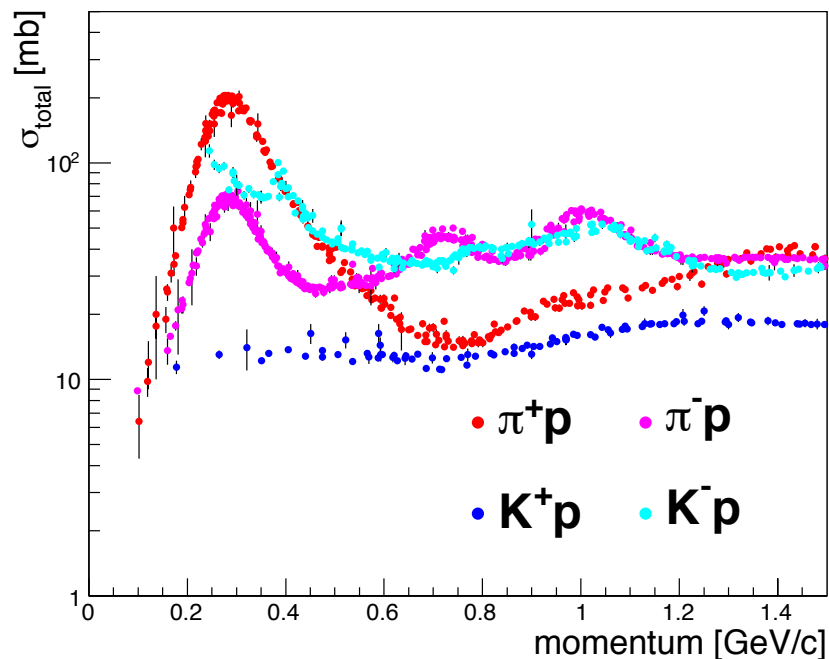
Time evolution of higher harmonic flow

Ideal hydrodynamic calculation
(Nonaka, correlation 2013)



Interaction cross section

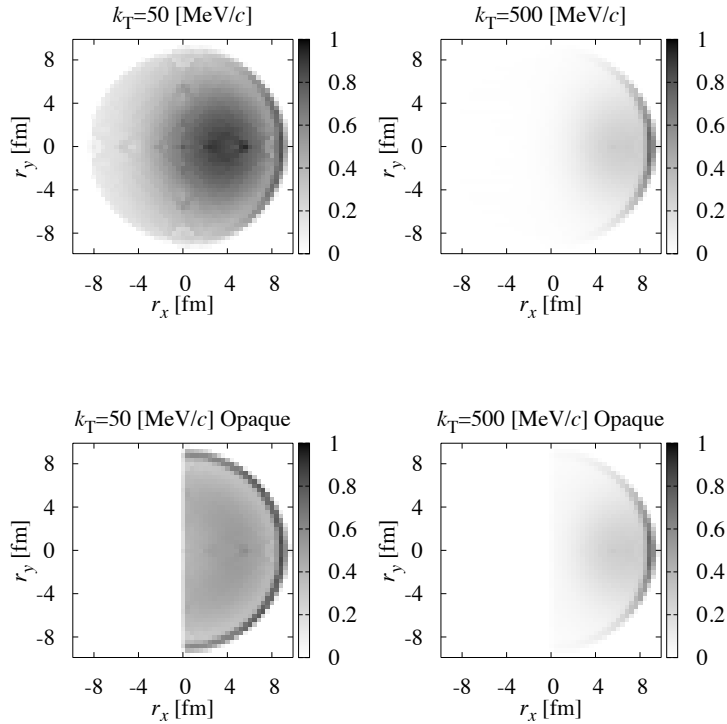
- Try looking at HBT radii for K^+K^+ and K^-K^- separately because they have different interaction cross section on nucleons.
- If π/K difference is due to the hadron rescattering with different cross sections, the difference may be also seen in positive and negative kaon pairs.



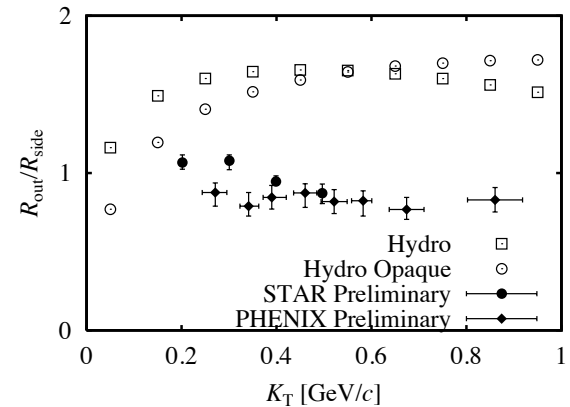
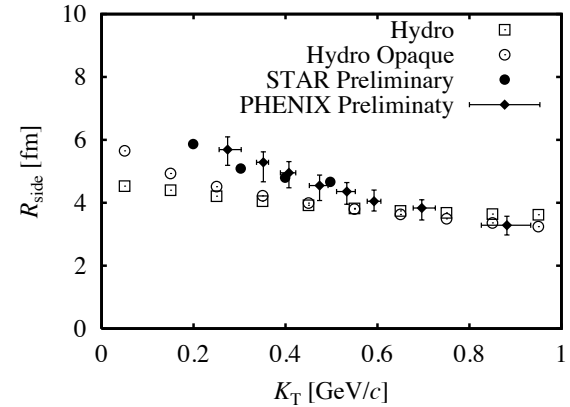
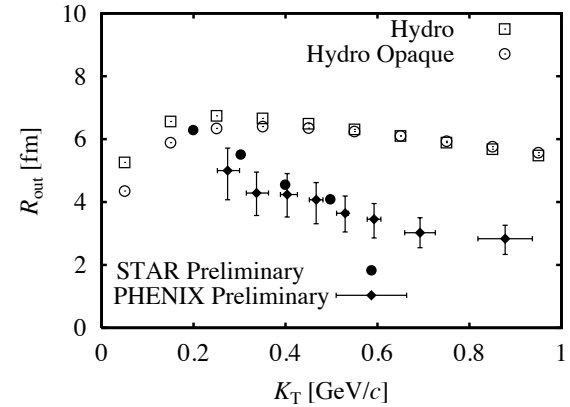
0.8 ~ 1.0 GeV/c
(π) = 0.69 ~ 1.1, $k_T(K) = 0.5 \sim 1.0$
($p_{T1} + p_{T2}$)/2

Effect of Opacity

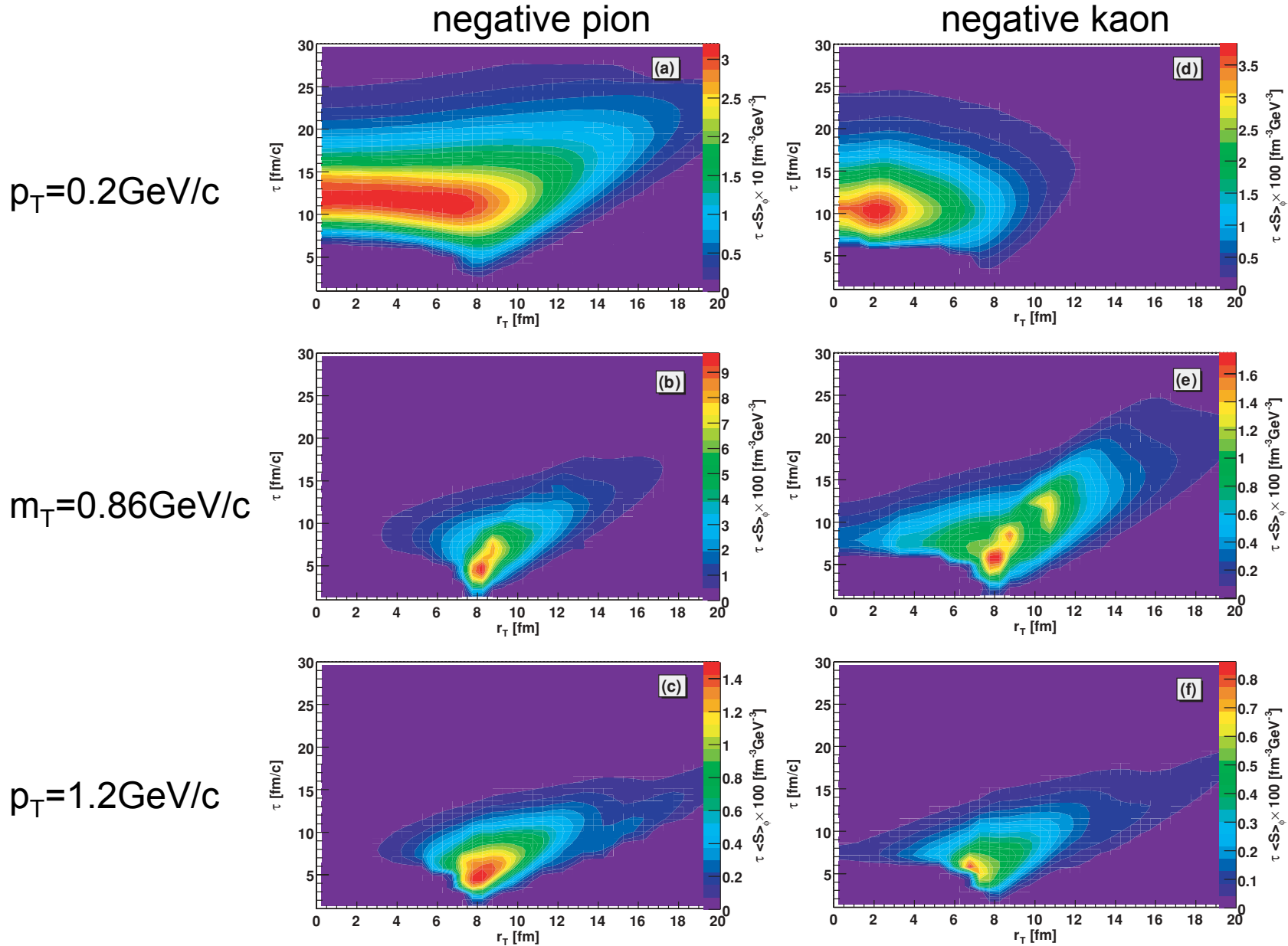
K. Morita, S. Muroya, Prog. Theor. Phys. 111, 93 (2004)



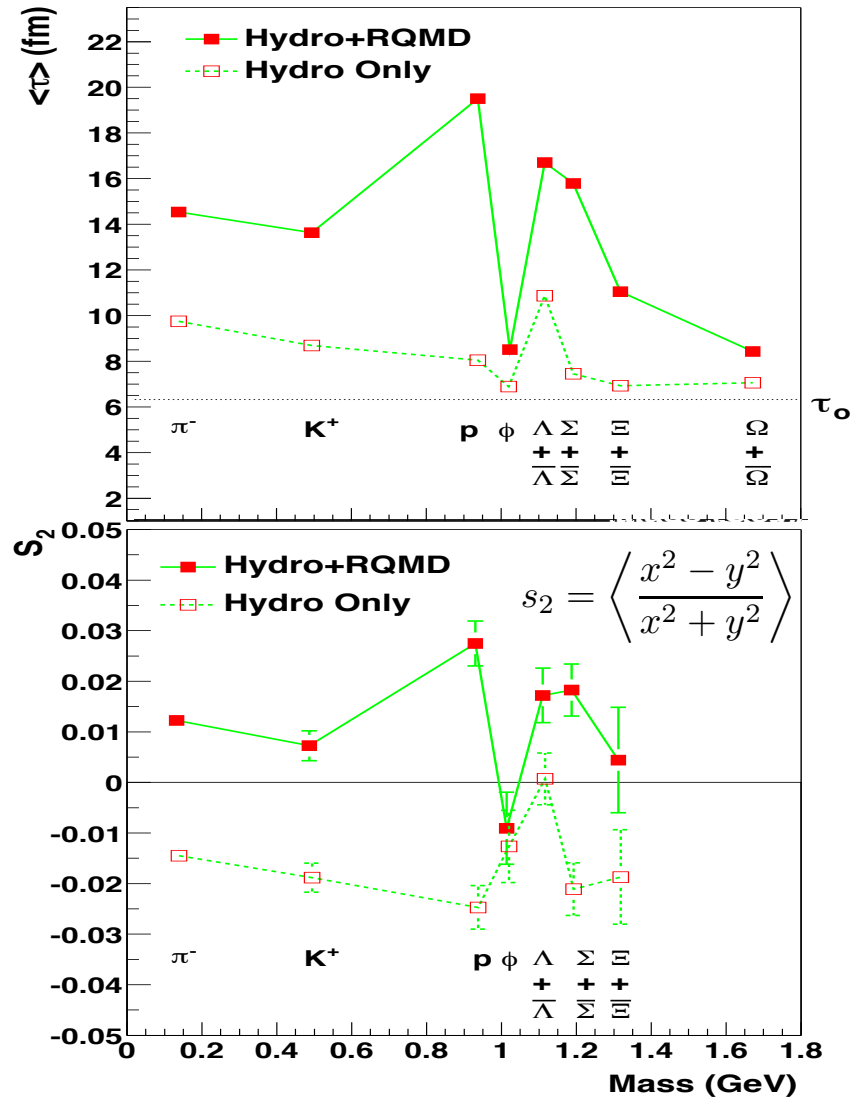
Opacity leads to a small decrease of R_o at low k_T and a small increase of R_s at low k_T .



Emission function from HKM model



Hydro+RQMD



Hydro only (including resonance decay)
vs
Hydro + hadronic rescattering