

Exploring the QCD phase diagram measured by cumulants of net-charge distributions in Au+Au collisions at the STAR experiment

STAR実験金+金衝突における
net-charge揺らぎを用いたQCD相図の探索

Tetsuro Sugiura
Defense
Feb. 8, 2019



My activity

♦ 2014 - 2015 Master course

♦ 2016 - 2019 Ph.D course

Conference and workshop : 10 talk and 2 poster presentations

• CiRfSE workshop, @University of Tsukuba	Mar. 2015 talk
• JPS @University of Miyazaki,	Jan. 2016 talk
★ QM2017 @Chicago,	Sep. 2016 talk
• JPS @Osaka University,	Feb. 2017 poster
• JPS @Utsunomiya University,	Mar. 2017 talk
• TCHoU workshop,	Sep. 2017 talk
★ QM2018 @Venice,	Mar. 2018 talk
• HIP/HIC @ Nagoya University	May. 2018 poster
★ APS/JPS joint meeting @Hawaii,	Jun. 2018 talk
★ QNP2018 @Tsukuba, - arXiv:1901.03639	Oct. 2018 talk
• TCHoU workshop,	Nov. 2018 talk
	Feb. 2019 talk

Blue: $\Delta\eta$ dependence analysis

Red: C_6 net-charge analysis

Green: Correction method

★ International conference

Master thesis:

- “ $\Delta\eta$ dependence of net-charge fluctuations in Au+Au collisions at RHIC-STAR experiment”

Teaching experience

Summer challenge TA @KEK
- (M2, D1, D2)

Paper: 1 paper contributed as 2nd author

T. Nonaka, T. Sugiura, S. Esumi, H. Masui, X. Luo

“Importance of separated efficiencies between positively and negatively charged particles for cumulant calculations” PRC. 94. 034909, Sep. 2016

STAR shift taking

: (D1, D2)

Outline

- Introduction
- Motivation
 - Sixth-order cumulants analysis (Experiment)
 - $\Delta\eta$ dependence of net-charge (Experiment)
 - Volume fluctuation study (Simulation)
- Analysis method (Experiment)
- Results (Experiment)

- Analysis method (Simulation)
- Results (Simulation)
- Summary and Outlook

Quark-Gluon Plasma (QGP)

- Quarks and gluons are confined inside of nucleons. (**Color confinement**)
- Interactions are very weak inside of nucleons. (**Asymptotic freedom**)

→

- Quarks are de-confined under high density or high temperature.
(Quark-Gluon Plasma ; QGP)

The diagram shows a chemical reaction-like equation: Pressure + Heat (creates pions) → Quark-Gluon Plasma. On the left, a blue arrow points to a red-bordered box containing the text: "Quarks are de-confined under high density or high temperature. (Quark-Gluon Plasma ; QGP)".

History of the Universe

A timeline diagram showing the evolution of the universe from the Big Bang to the present. It includes a legend for particle symbols and a key for physical phenomena.

Key:

- W, Z bosons
- Photon
- q quark
- g gluon
- e electron
- μ muon
- τ tau
- v neutrino
- q̄ meson
- baryon
- ion
- atom
- Galaxy
- Star
- Black hole

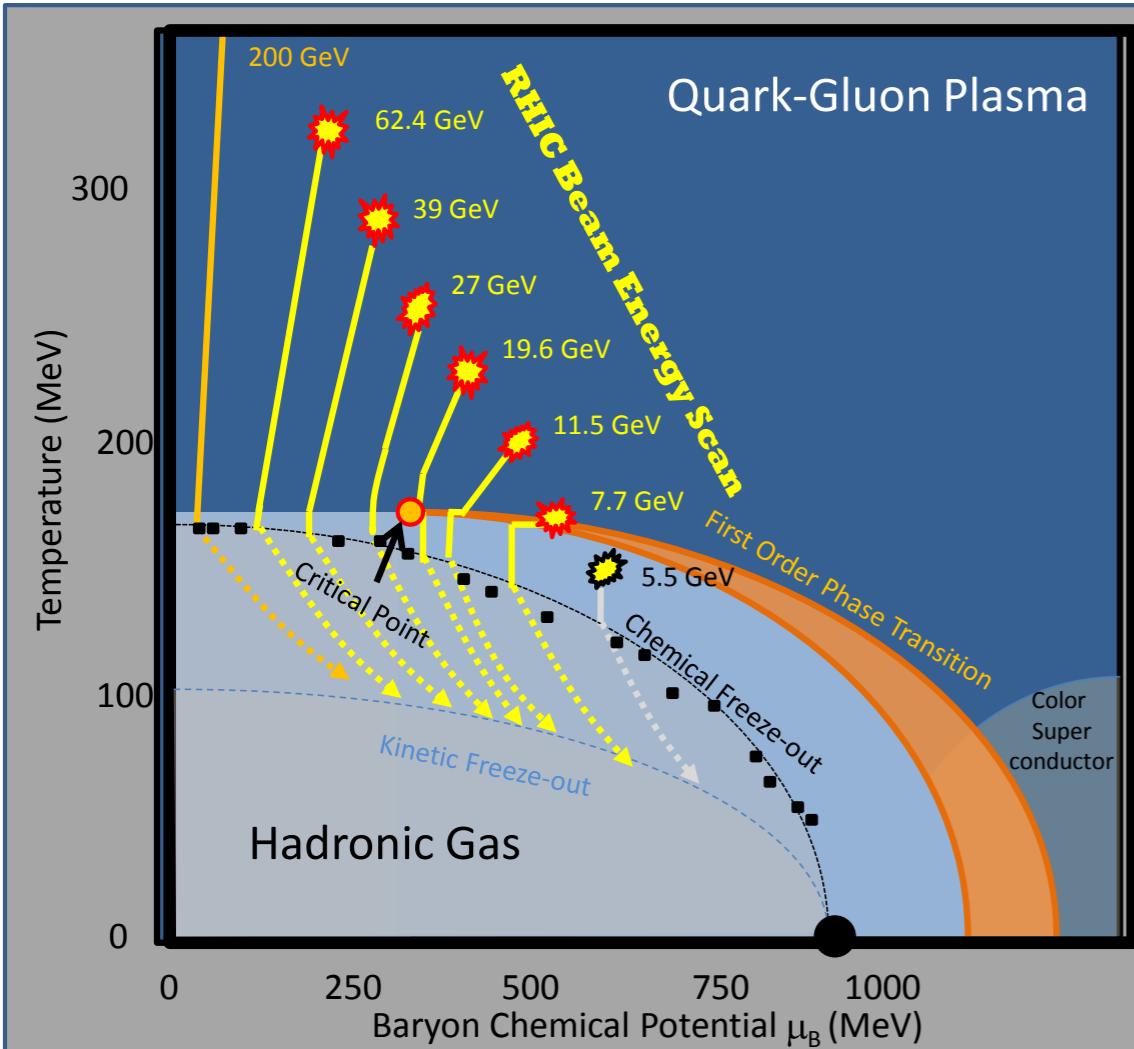
Particle Data Group, LBNL, © 2008. Supported by DOE and NSF

Pressure + Heat (creates pions) → Quark-Gluon Plasma

$T_c \sim 160 \text{ MeV}$

- QGP is considered to be exist in early universe, μs after big bang.

Phase transition and Beam Energy Scan



- Cross over transition around $\mu_B=0$ from Lattice QCD but there is no experimental evidence.
- Location of critical point is unknown.



**Detail structure of QCD diagram
should be discovered experimentally**

Beam Energy Scan I (BES-I) 2010-2014

- Varying the center of mass energy, $\sqrt{s_{NN}} = 7.7, 11.5, 14.5, 19.6, 27, 39, 62.4$ and 200 GeV in Au+Au collisions.
- $\sqrt{s_{NN}} = 54$ GeV in Au+Au collision data was taken in 2017.
- BES-II will start from 2019.



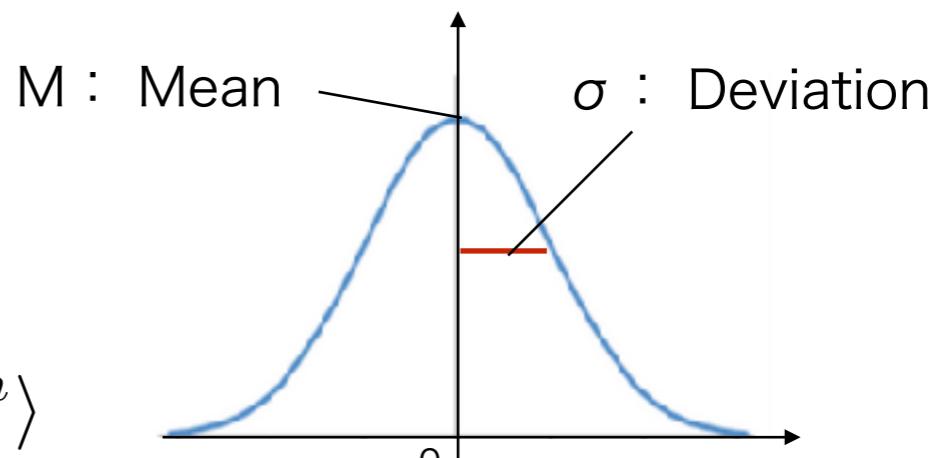
Event by event fluctuations are considered as one of the powerful tool to discover the phase diagram.

Event-by-event fluctuations

Net-charge : $N_+ - N_-$

- nth-order moment is defined by

$$\langle m^n \rangle = \sum_m m^n P(m), \quad \langle \delta m^n \rangle = \langle (m - \langle m \rangle)^n \rangle$$



- Cumulants can be written by moments

Cumulants \leftrightarrow Moments

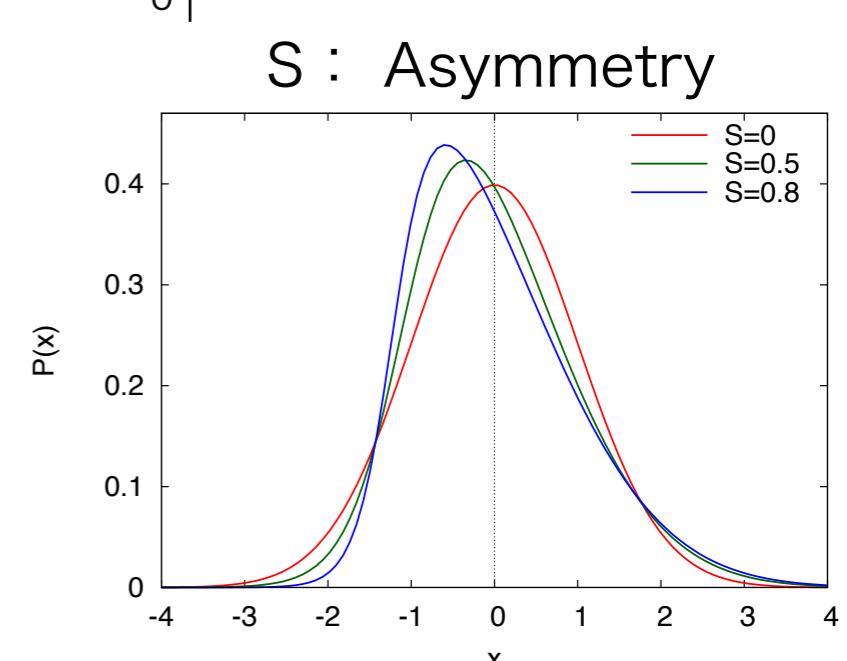
C_1	$= \langle m \rangle$
C_2	$= \langle \delta m^2 \rangle$
C_3	$= \langle \delta m^3 \rangle$
C_4	$= \langle \delta m^4 \rangle - \langle \delta m^3 \rangle$

$$M = C_1$$

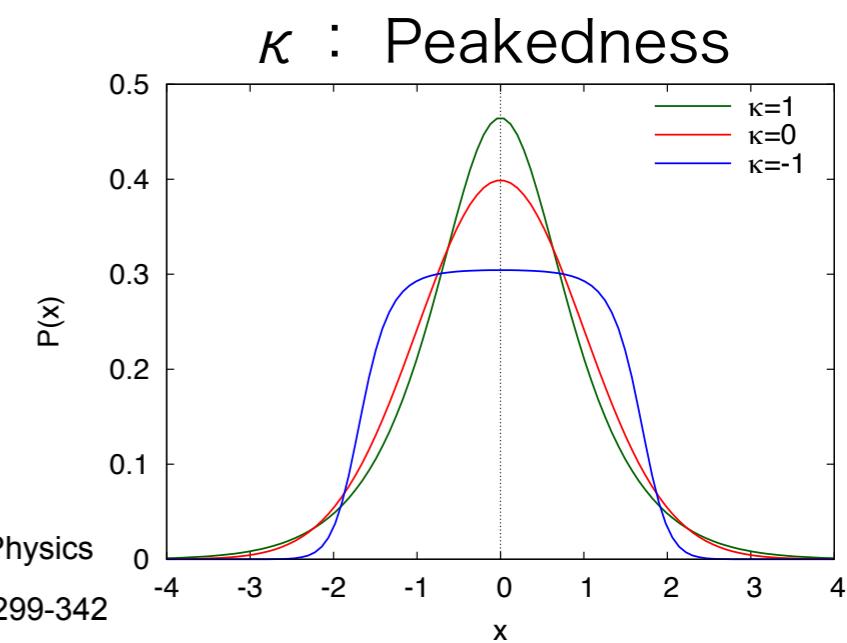
$$\sigma^2 = C_2$$

$$S = \frac{C_3}{(C_2)^{3/2}}$$

$$\kappa = \frac{C_4}{(C_2)^2}$$



- Cumulants are proportional to the volume
(Additivity)



Cumulant ratios and Critical behaviour

Cumulant
(can be measured experimentally)

$$C_n = (-T)^{n-1} \frac{\partial^n \Omega}{\partial \mu^n} \equiv \chi_n V$$

Susceptibility
(can NOT be measured experimentally)

- Cumulant ratios **can be directly compared to the ratio of susceptibilities.**
- By measuring the **cumulant ratios, effect from the volume can be canceled.**

$$S_\sigma = \frac{C_3}{C_2} = \frac{\chi_3}{\chi_2}$$

$$\kappa\sigma^2 = \frac{C_4}{C_2} = \frac{\chi_4}{\chi_2}$$

- Higher order cumulants are more sensitive to the signal from the critical point.
- These relations are only valid for cumulants of **conserved quantities.**

$$C_2 \approx \xi^2 \quad C_5 \approx \xi^{9.5}$$

$$C_3 \approx \xi^{4.5} \quad C_6 \approx \xi^{12}$$

$$C_4 \approx \xi^7$$

ξ : Correlation length

Statistical baseline (Skellam distributions)

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- If distribution of N_+ and N_- follow the Poisson distributions, net-charge distribution follow the **Skellam distribution**.

$$C_n = \lambda_1 + (-1)^n \lambda_2$$

λ_1 : Poisson parameter of N_+ distribution

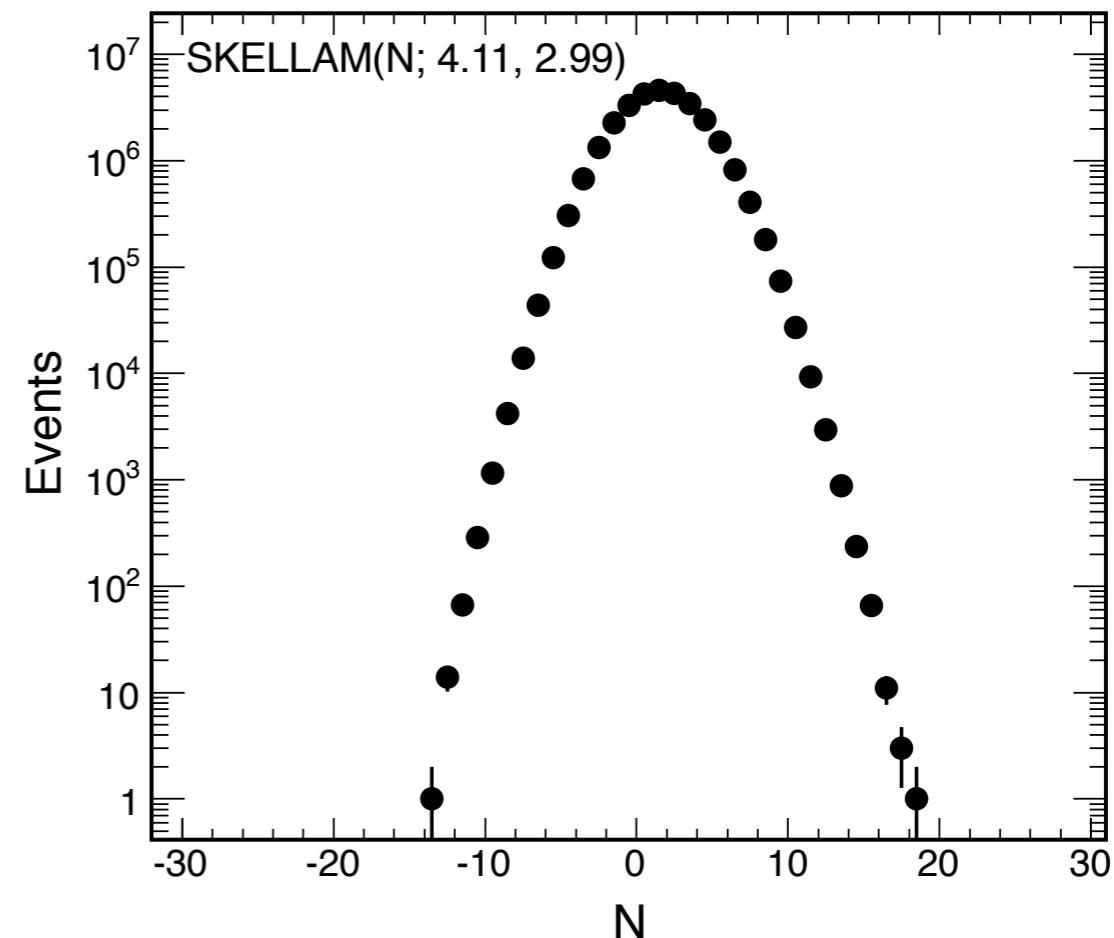
λ_2 : Poisson parameter of N_- distribution

$$\frac{C_{2n}}{C_{2m}} = \frac{C_{2n+1}}{C_{2m+1}} = 1$$

$$\frac{C_{2n+1}}{C_{2m}} = \frac{\lambda_1 - \lambda_2}{\lambda_1 + \lambda_2}$$

For example

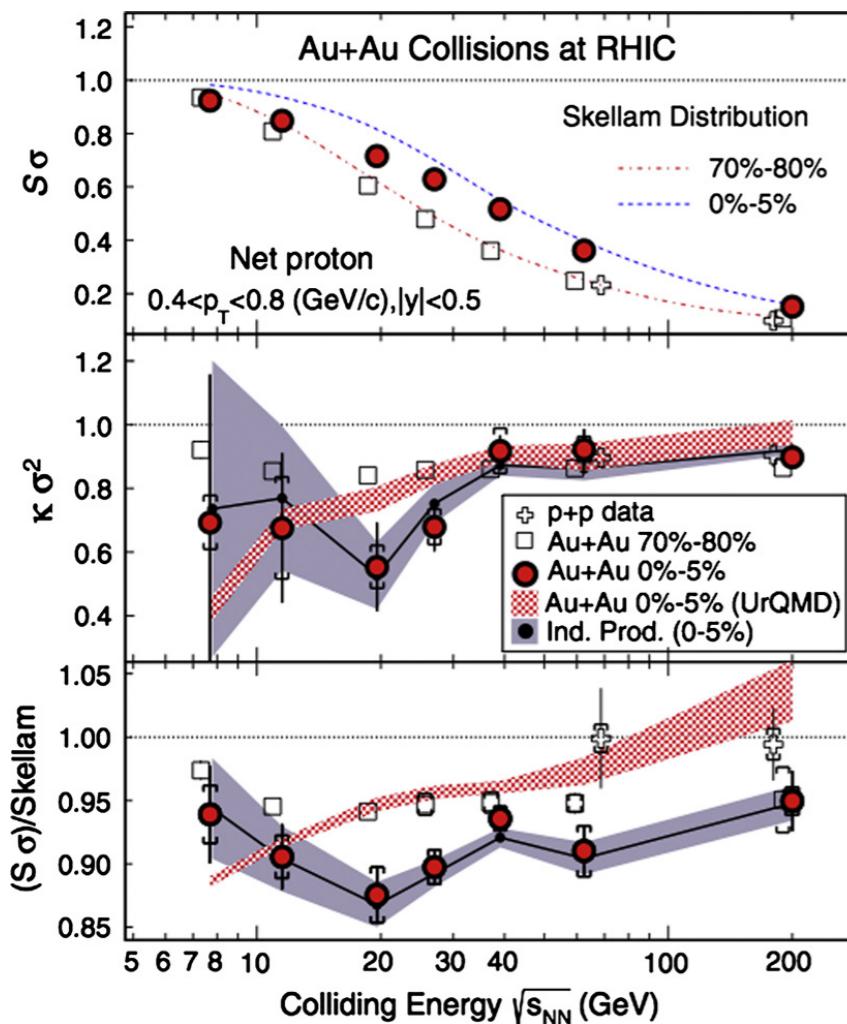
$$\frac{C_6}{C_2} = \frac{C_4}{C_2} = \boxed{1}$$



This values can be used as a statistical baseline.

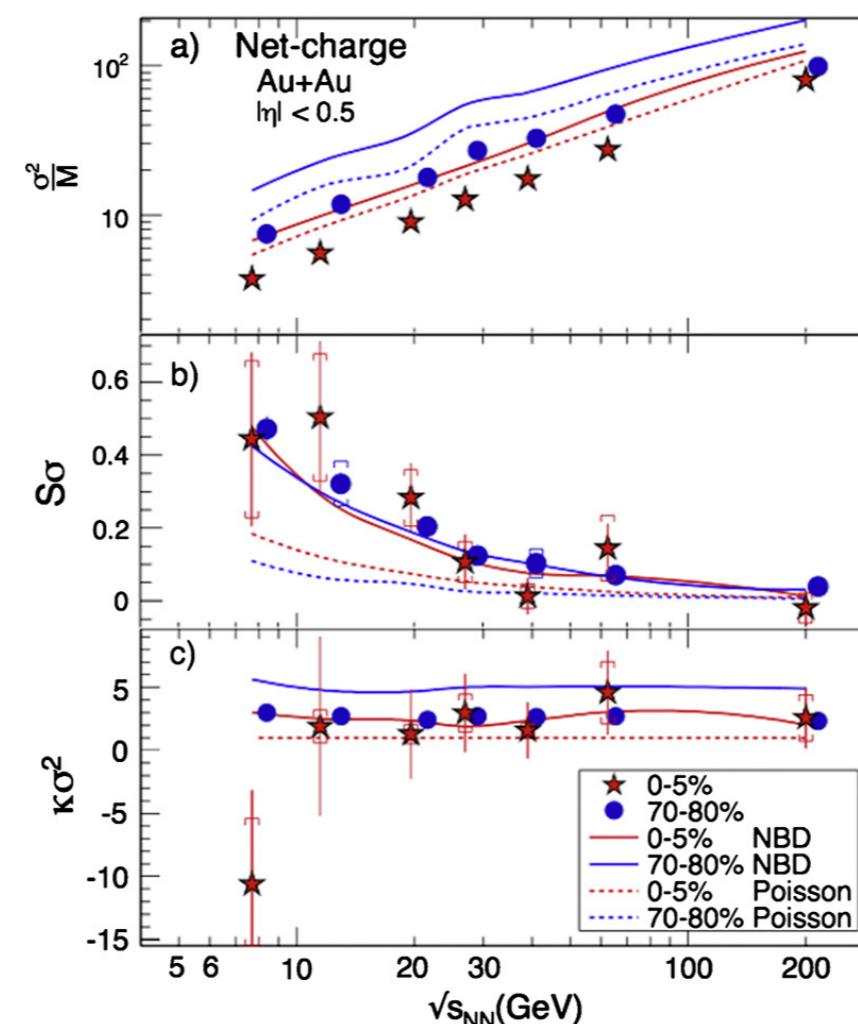
Published results

Net-proton



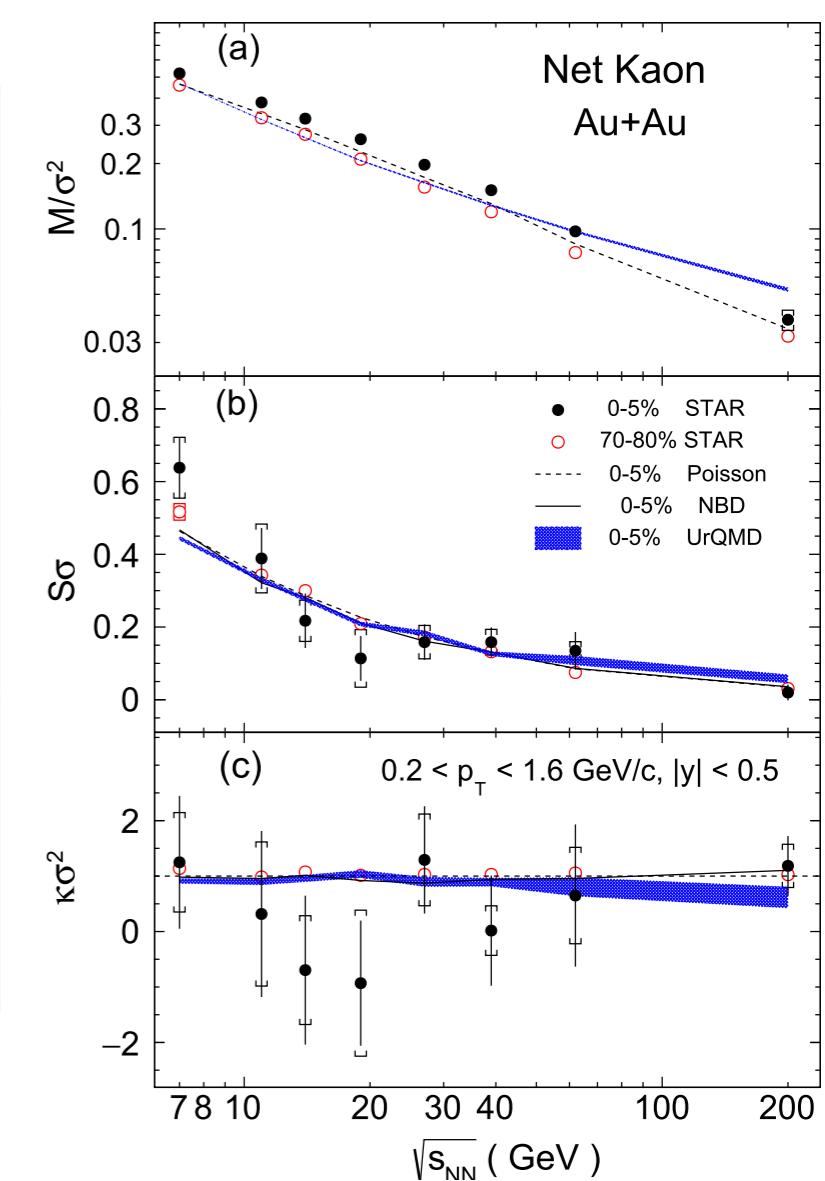
L. Adamczyk *et al.* (STAR Collaboration)
Phys. Rev. Lett. 112, 032302 (2014)

Net-charge



L. Adamczyk *et al.* (STAR Collaboration)
Phys. Rev. Lett. 113, 092301(2014)

Net-Kaon

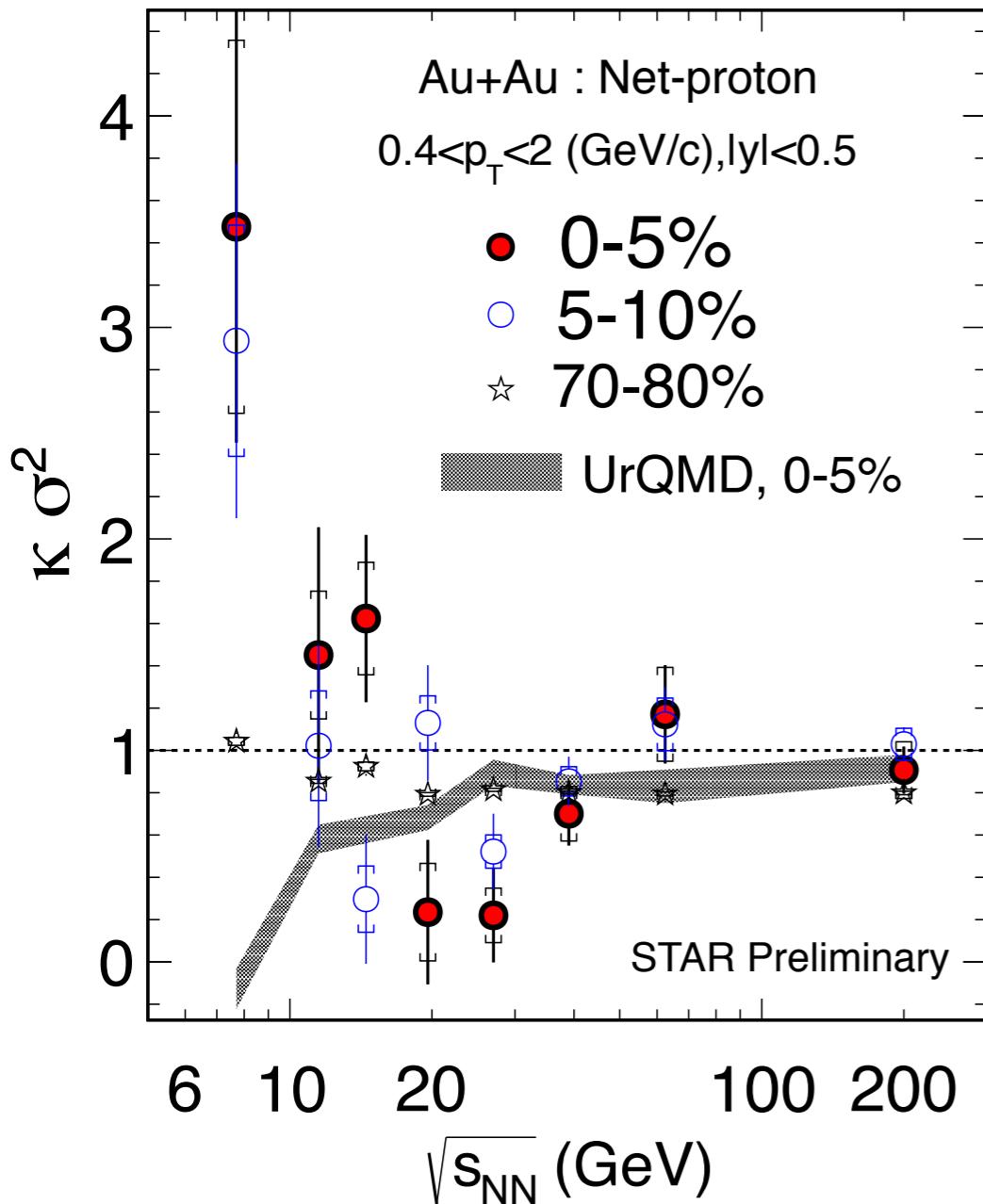


Physics Letters B [Volume 785](#),
10 October 2018, Pages 551-560

- **STAR published up to 4-th order fluctuations** of net-proton, net-charge and net-kaon distributions.
- **Critical phenomena have not been observed yet.**

Net-proton preliminary results

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- Interesting energy dependences were observed by **expanding p_T range by using TOF.**
- Minimum around $\sqrt{s_{NN}} = 20$ GeV
- Larger than unity at $\sqrt{s_{NN}} = 7.7$ GeV



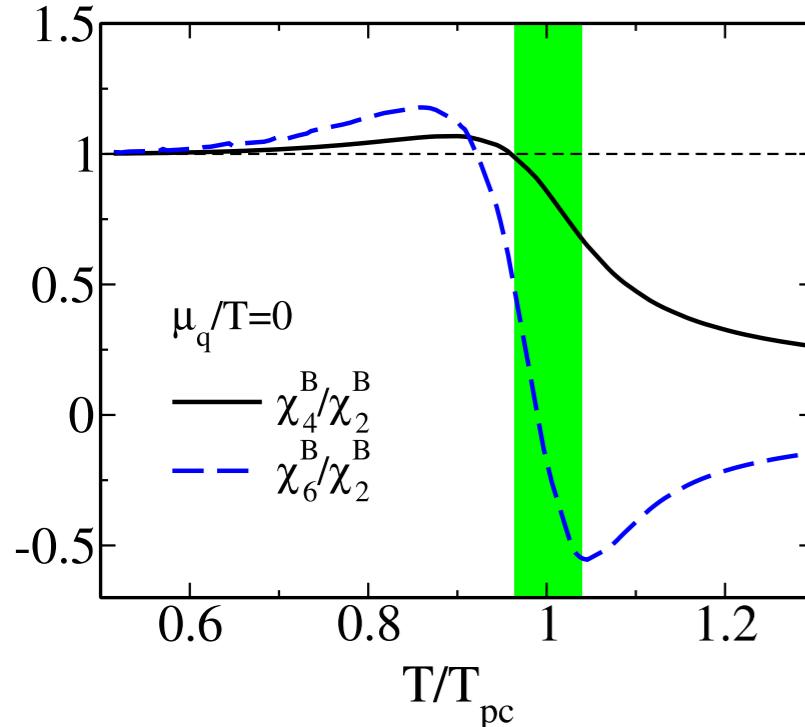
- Signal from phase transition or critical point?
- More statistics is necessary.

X. Luo, Nucl. Phys. A956 (2016) 75–82, [1512.09215].

Motivation 1 -C₆ analysis-

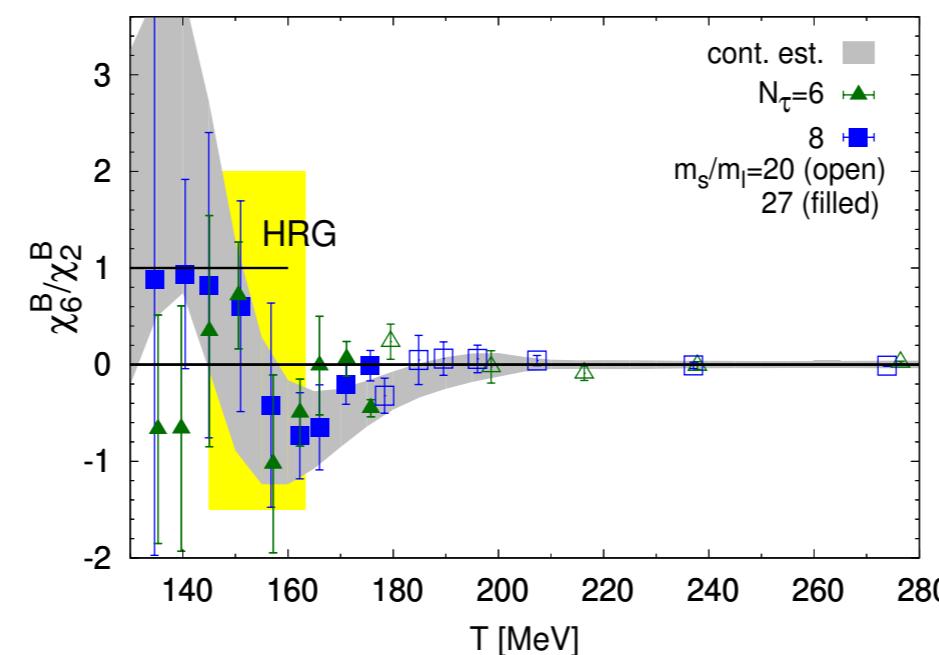
PQM model prediction

Friman et al, Eur. Phys. J. C(2011) 71:1694

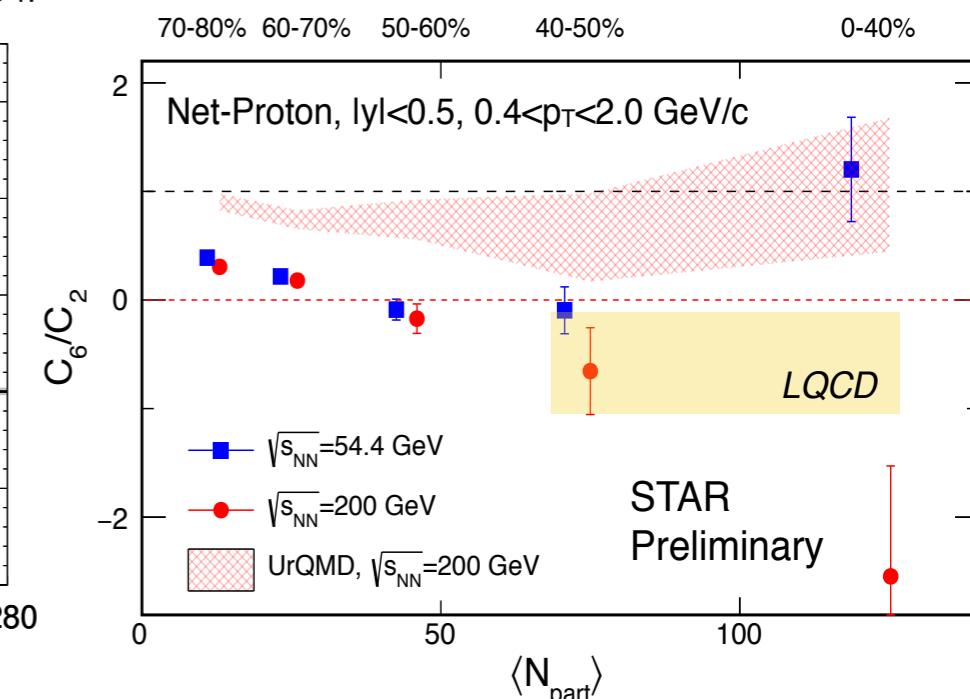


Lattice QCD prediction

A. Bazavov et al, Phys.Rev. D95 (2017) no.5, 054504.



T.Nonaka (STAR Collaboration),
QNP2018



freeze-out conditions	χ_4^B/χ_2^B	χ_6^B/χ_2^B	χ_4^Q/χ_2^Q	χ_6^Q/χ_2^Q
HRG	1	1	~ 2	~ 10
QCD: $T^{freeze}/T_{pc} \lesssim 0.9$	$\gtrsim 1$	$\gtrsim 1$	~ 2	~ 10
QCD: $T^{freeze}/T_{pc} \simeq 1$	~ 0.5	< 0	~ 1	< 0

- **Negative C₆** of conserved quantities are predicted to be the signal of **cross-over phase transition**.
- C₆ of net-proton is measured at $\sqrt{s_{NN}} = 200$ and 54 GeV in Au+Au collisions.
- **C₆ of net-charge** should be measured to compare net-proton results and theoretical predictions.
- 54 GeV net-charge results have not been reported yet.

Motivation2 - $\Delta\eta$ dependence-

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D-measure (D)

$$D = 4 \frac{C_2}{\langle N_{ch} \rangle} \quad N_{ch} = N^+ + N^-$$

Theoretical predictions

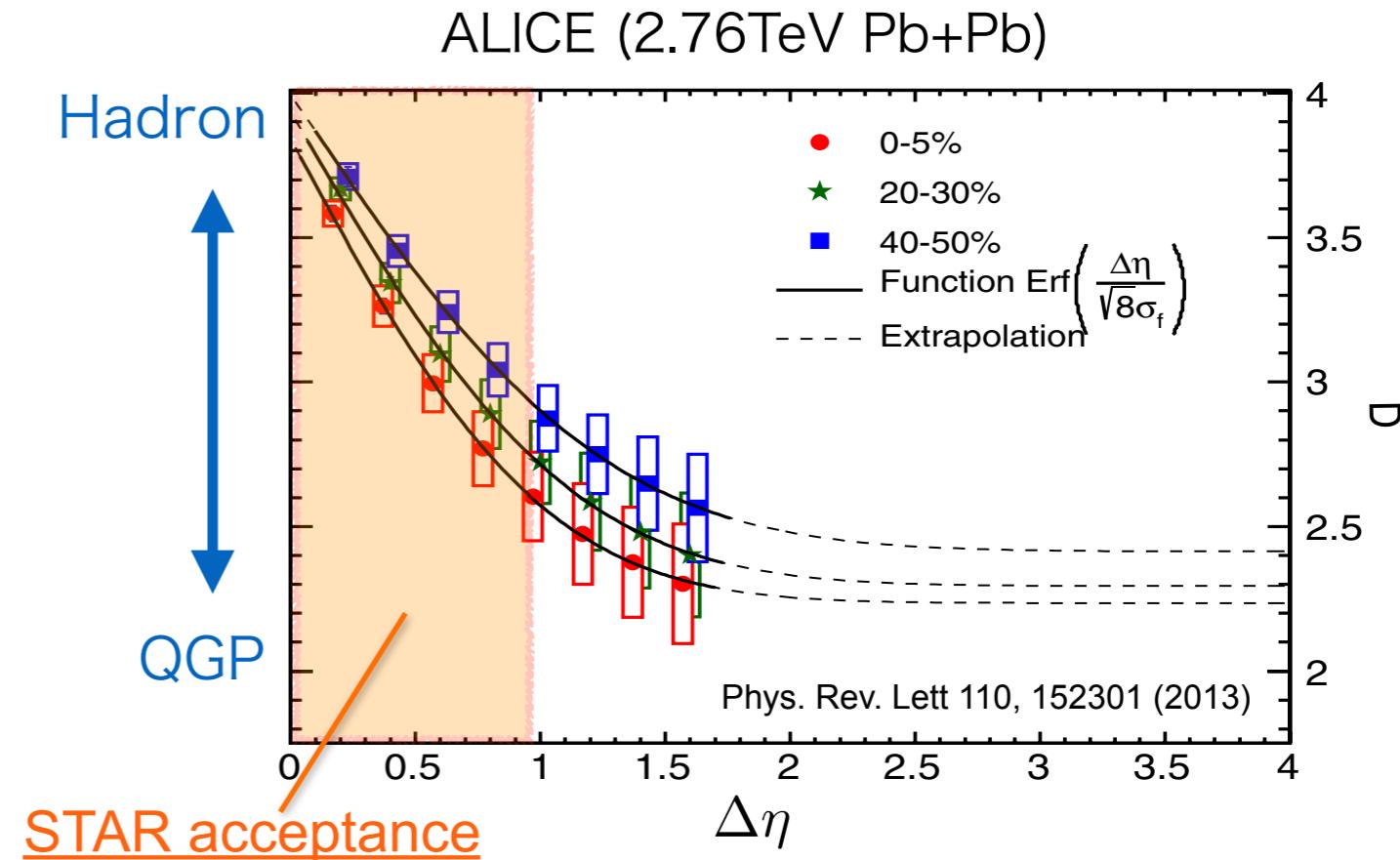
(HRG model, Lattice QCD, etc)

QGP fluctuation :

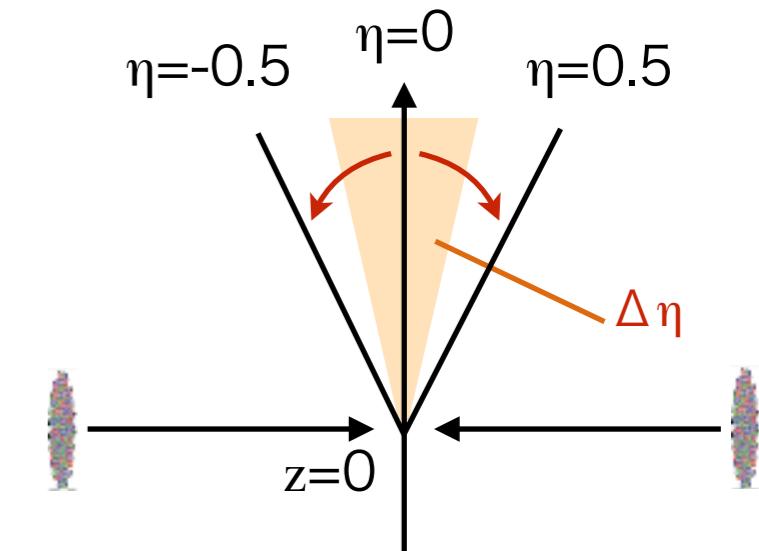
$$D = 1-1.5$$

Hadron fluctuation :

$$D = 3-4$$



- D-measure is observed to decrease with expanding $\Delta\eta$ in Pb-Pb collisions at 2.76 TeV at ALICE.**
- D-measure also decrease from peripheral to central collisions.

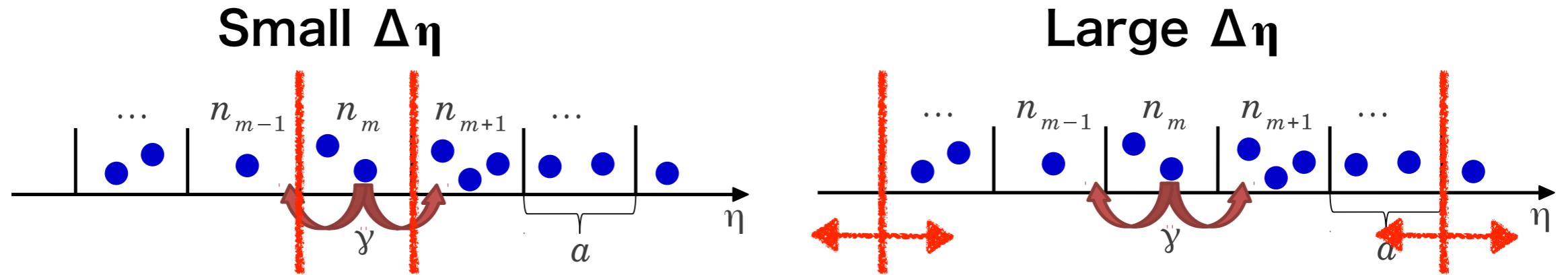


Expanding $\Delta\eta$: The signal of QCD transition might been observed



It is important to measure $\Delta\eta$ dependence of D-measure in lower energies.

Motivation2 - $\Delta\eta$ dependence-



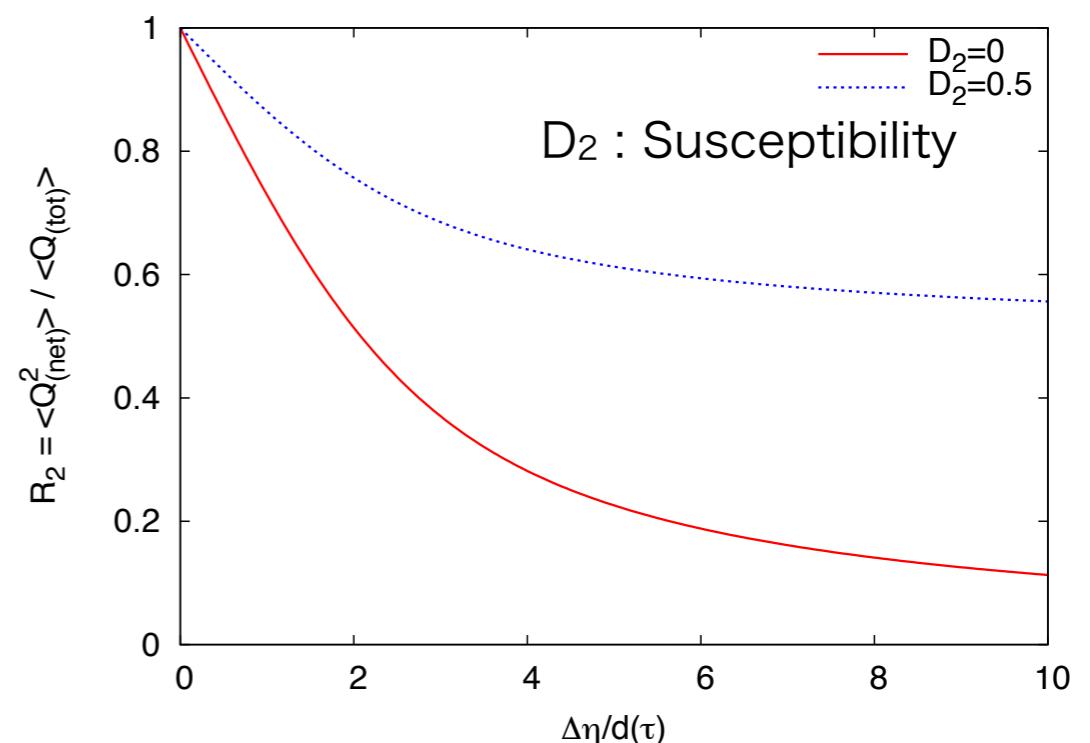
$$\partial_\tau P(\mathbf{n}, \tau) = \gamma(t) \sum_m [(n_m + 1) \{P(\mathbf{n} + \mathbf{e}_m - \mathbf{e}_{m+1}, \tau) + P(\mathbf{n} + \mathbf{e}_m - \mathbf{e}_{m-1}, \tau)\} - 2n_m P(\mathbf{n}, \tau)]$$

$$\mathbf{n} = (\dots, n_{m-1}, n_m, n_{m+1}, \dots)$$

$$a \rightarrow 0$$

- Decreasing D-measure with $\Delta\eta$ can be described by diffusion model.
- Third and fourth-order cumulants by this model are implemented.

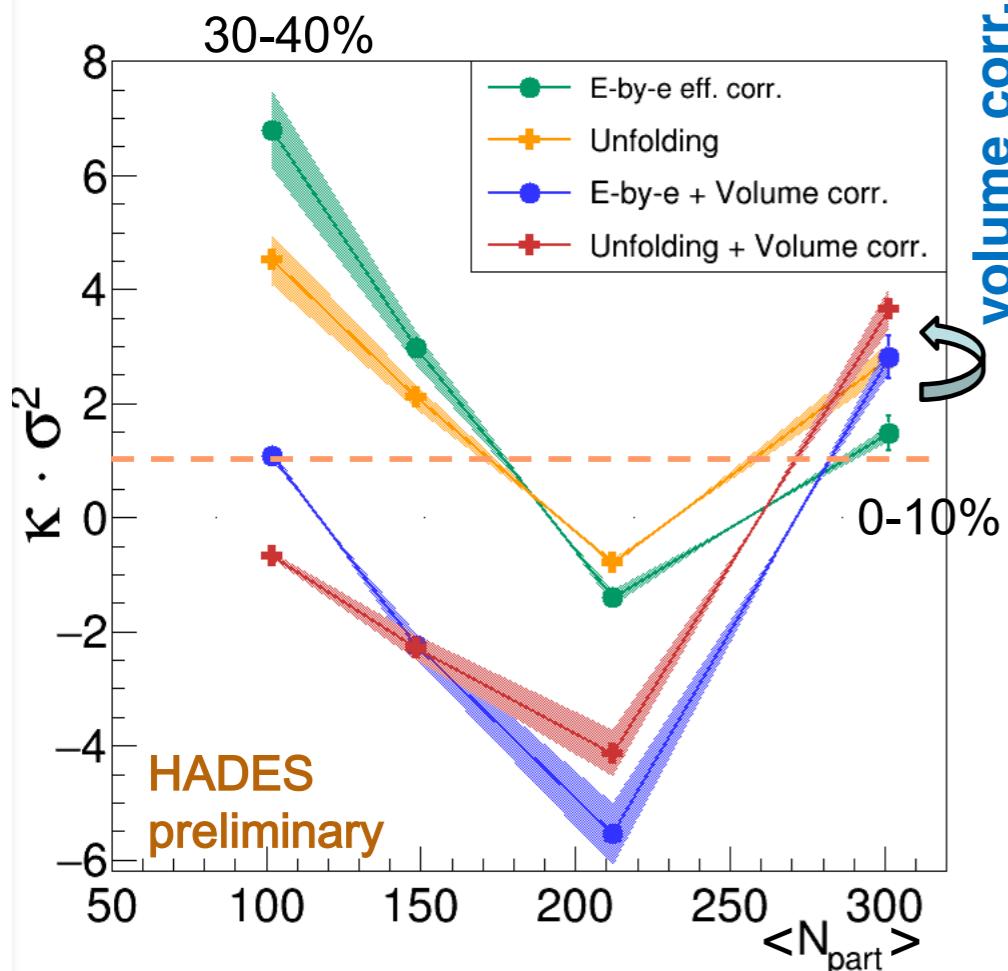
M.Asakawa, M.Kitazawa, Progress in Particle and Nuclear Physics
Volume 90, September 2016, Pages 299-342



Experimental measurement are necessary up to 4th-order

Motivation3 - Volume Fluctuation study-

R.Holzmann (HADES collaboration) QM 2017



- Initial volume fluctuation (VF) which correspond to the e-by-e fluctuation of number participant nucleons ($N_{\text{part}}=N_w$) artificially enhance cumulants and **should be subtracted from measured cumulants**
- STAR applied **Centrality Bin Width Correction (CBWC)** to eliminate VF.
- New correction **Volume fluctuation Correction (VFC)** was proposed and HADES experiment have already applied VFC to experimental data.

However...

- VFC is model dependent** correction whereas **CBWC is data driven** correction.

Importance of subtracting VF and validity of VFC are studied by using toy model and UrQMD in both net-charge and net-proton cases.

Why net-charge?

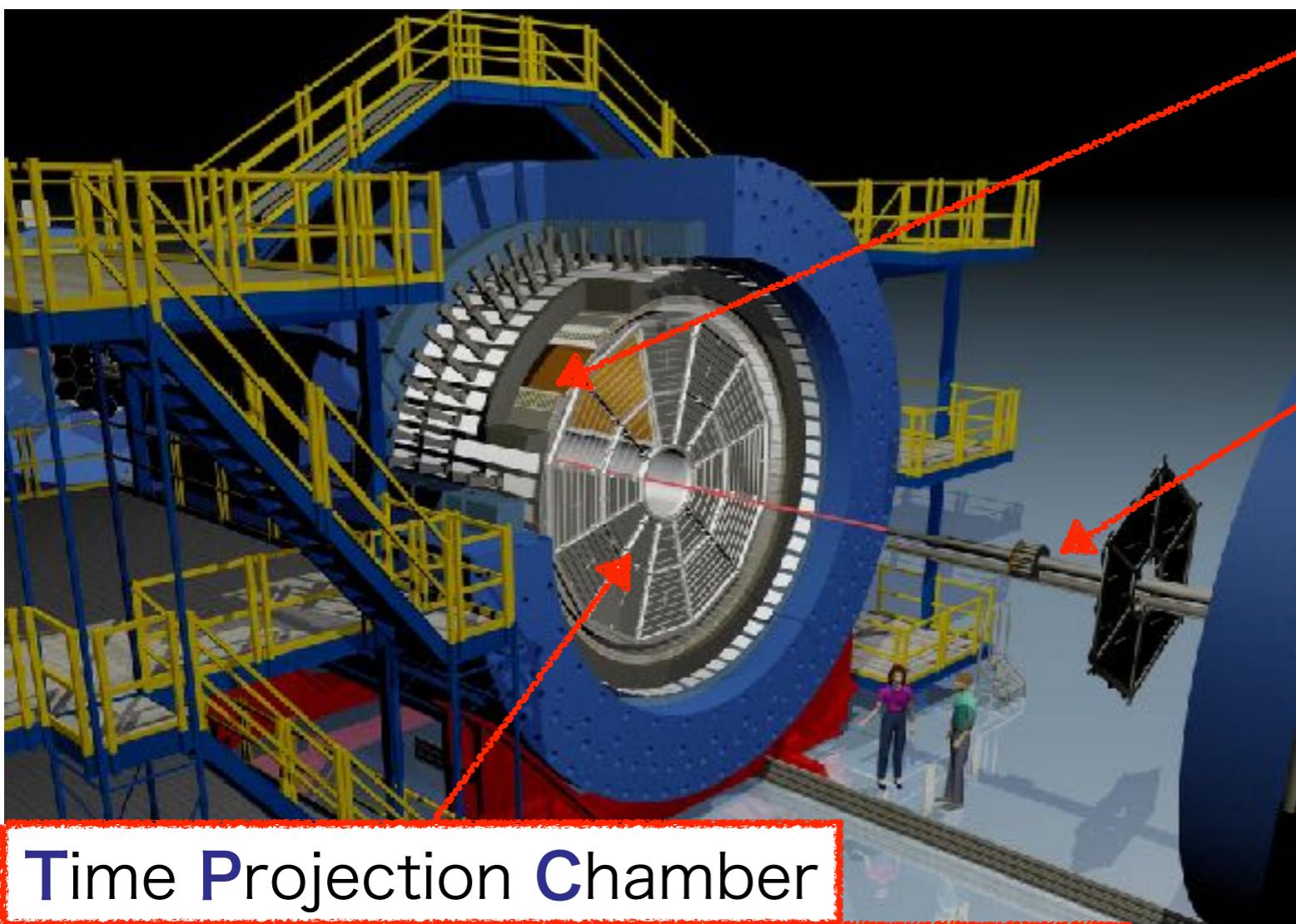
- It is important to measure **both net-proton and net-charge** cumulants and compare the results in order to understand the phase transition.
- Net-charge is conserved quantity** whereas net-proton and net-Kaon are measured as a proxy of net-baryon and net-strangeness.
 - My PhD. thesis, University of Tsukuba (2019)
 - T.Nonaka's PhD. thesis, University of Tsukuba (2018)

	net-charge	net-proton	net-Kaon
Conserved quantity up to C_4	net-charge	net-baryon	net-strangeness
up to C_6	published (7.7-200 GeV)	published (7.7-200 GeV)	published (7.7-200 GeV)
$\Delta\eta$ dependence	200GeV, 54GeV	200GeV, 54GeV	-
	7.7-200 GeV (STAR) 2.76 TeV (ALICE)	-	-

- In **C_6 analysis**, analysis and correction **methods are improved**.
- $\Delta\eta$ dependence** of D-measure (**2nd order** fluctuation) is reported **by ALICE** whereas **up to 4th-order** cumulants are measured **in this thesis**.
- VFC is studied by T.Nonaka's Doctoral thesis but more studies are necessary.

Analysis mothod (Experiment)

The STAR Detector



Time Projection Chamber

- Used for tracking, momentum determination and PID.
- Used to determine centrality and to calculate net-charge cumulants.
- using different kinematic window to avoid auto-correlation

Time-Of-Flight

- PID in a broad p_T region.
- Used to remove pile-up events.

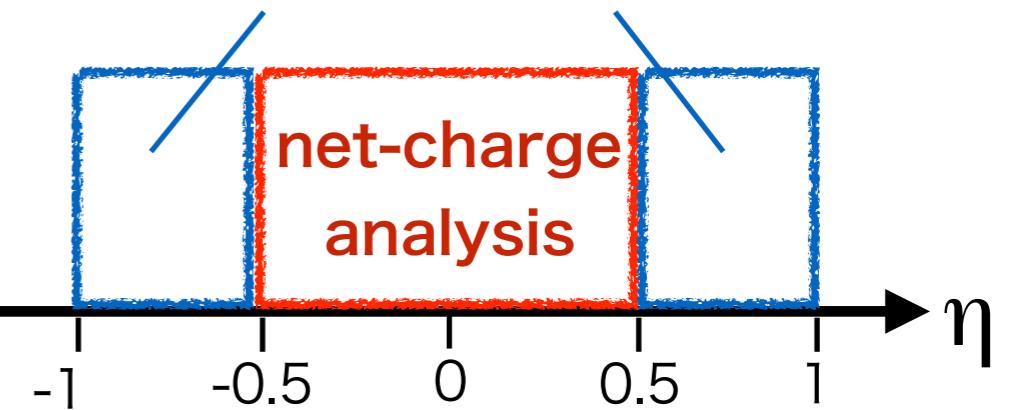
Vertex Position Detector

Zero Degree Calorimeter

- Minimum bias trigger
- Used for start time (VPD)
- Collision vertex (VPD)
- Luminosity monitor (ZDC)

ZDC : 18m from vertex

Determine centrality (Refmult2)



Data set

$\sqrt{s_{NN}} =$ 200 GeV	Run11(0-80%)	Run10(10-80%)	Run10(0-10%)	UrQMD
NEvent	485M	211M	196M	45M

$\sqrt{s_{NN}} =$ 54 GeV	Run17
NEvent	543M

$\Delta\eta$ analysis

$\sqrt{s_{NN}}$	NEvent (Million)
7.7	1.55
11.5	2.57
14.5	12
19.6	15.5
27	27.5
39	85.3
62.4	50.4
200	97.8

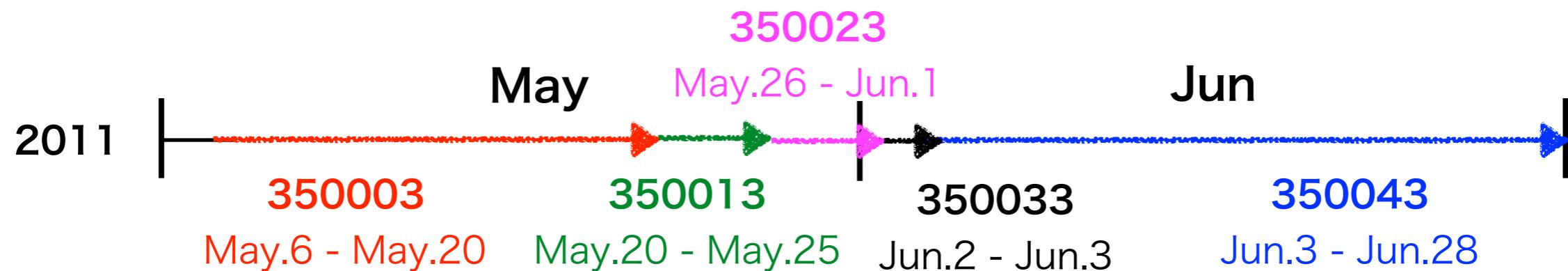
C₆ analysis

- C₆ analysis needs large statistics compared to up to 4th order fluctuation analysis
 - 200 GeV data was taken from 2010 to 2011 and 54 GeV data was taken in 2017.
 - Statistics of both of them are larger than 100M.
 - Some analysis and correction method are improved compared to published results.
- All BES-I data were used for $\Delta\eta$.
 - analysis method and cuts are almost same as published results (up to 4th order).

Trigger ID

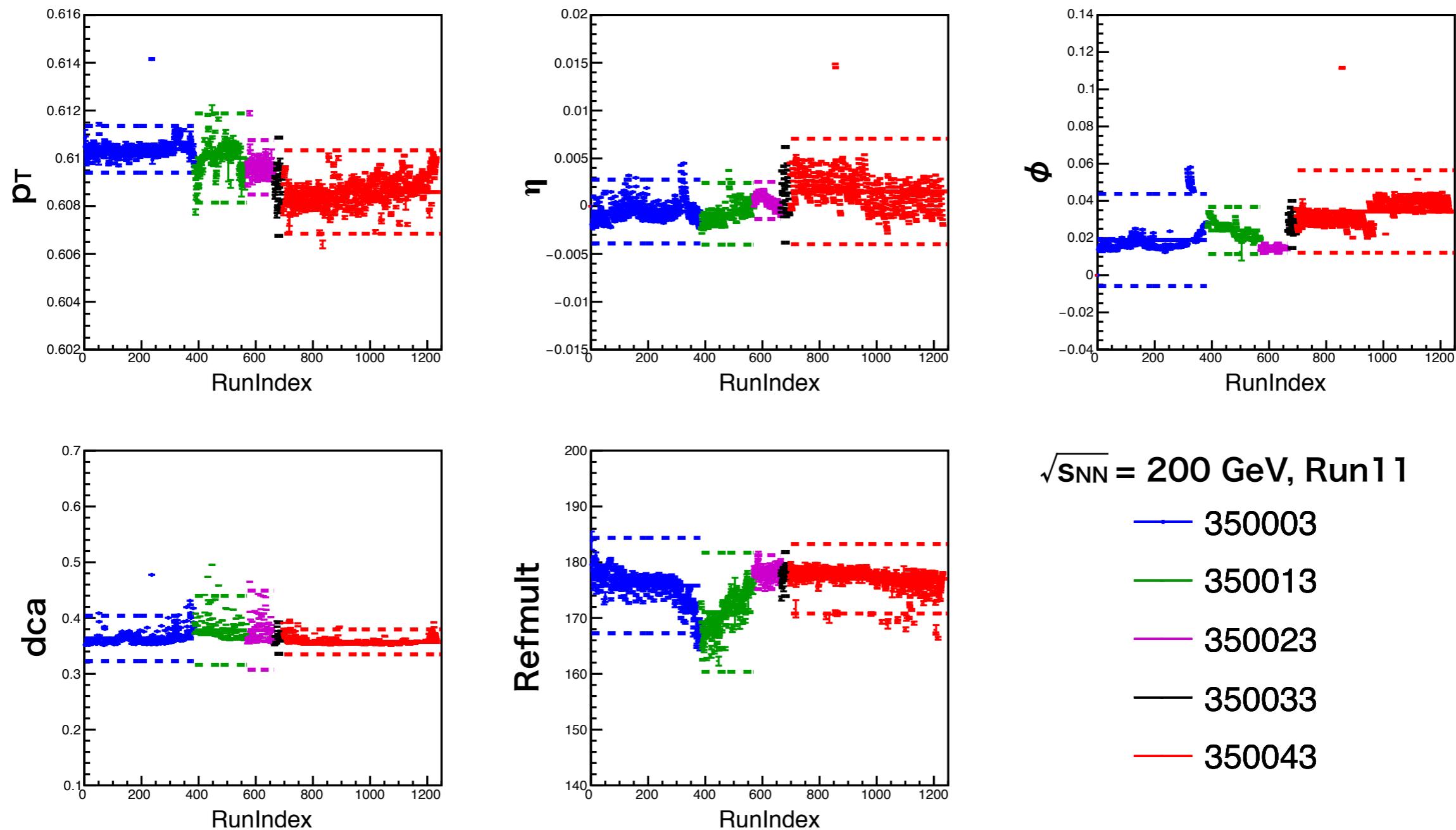
- Run11 data correspond to the data taken in 2011.
- There are several trigger ID. (Minimul bias trigger)
- Trigger ID changed when detector conditions were changed.

$\sqrt{s_{NN}} =$ 200 GeV	350003	350013	350023	350033	350043
Run11	200M	74M	15M	12M	187M



- For consistency check, cumulants are calculated for each trigger ID once and then merged at 200GeV.

Run by run QA

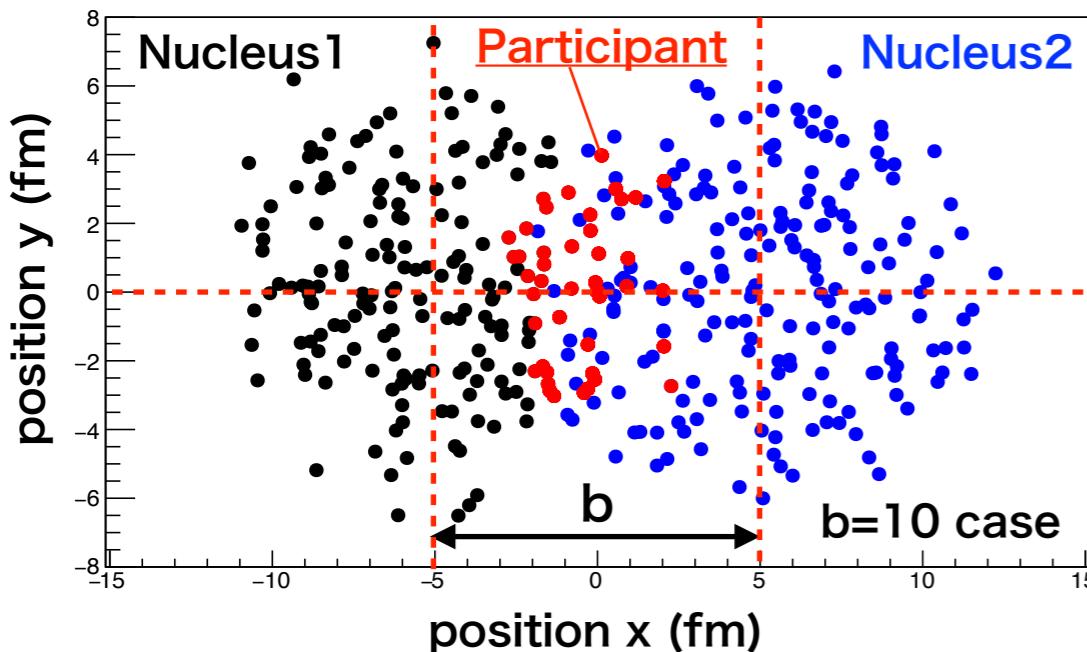


- Run by run $\langle p_T \rangle$, $\langle \eta \rangle$, $\langle \phi \rangle$, $\langle dca \rangle$, and $\langle \text{Refmult} \rangle$ are measured
- The outlier runs of 3σ were rejected as bad runs for each trigger ID

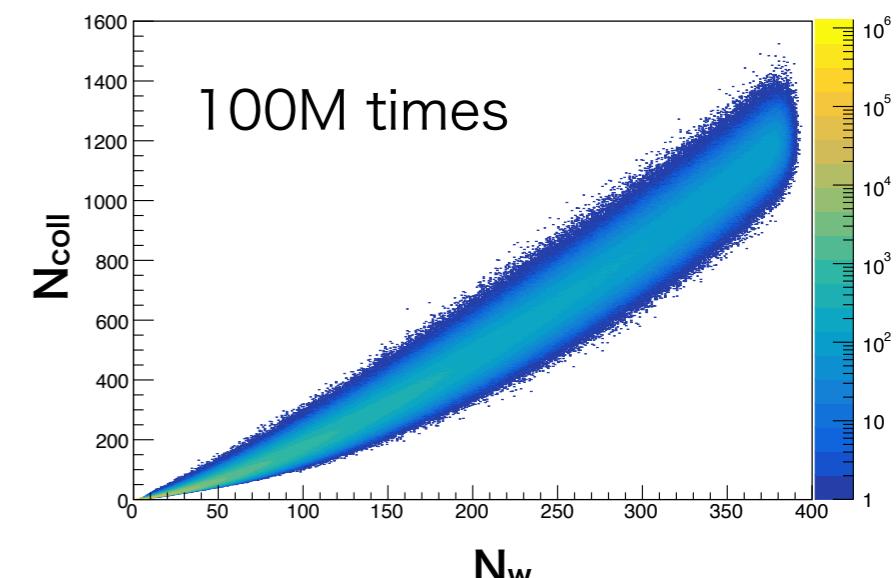
Glauber model

- N_w can't be measured directly by experiment so estimated by Glauber model.

1. Calculate N_w and N_{coll} by Monte Carlo simulation.



- Width : **0.535**
 - Radius : **6.4**
 - σ : **42 (mb)**
- Repeat



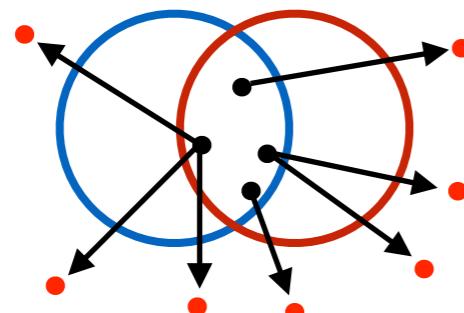
2. Calculate multiplicity by two component model and NBD.

Two component model

Number of source :

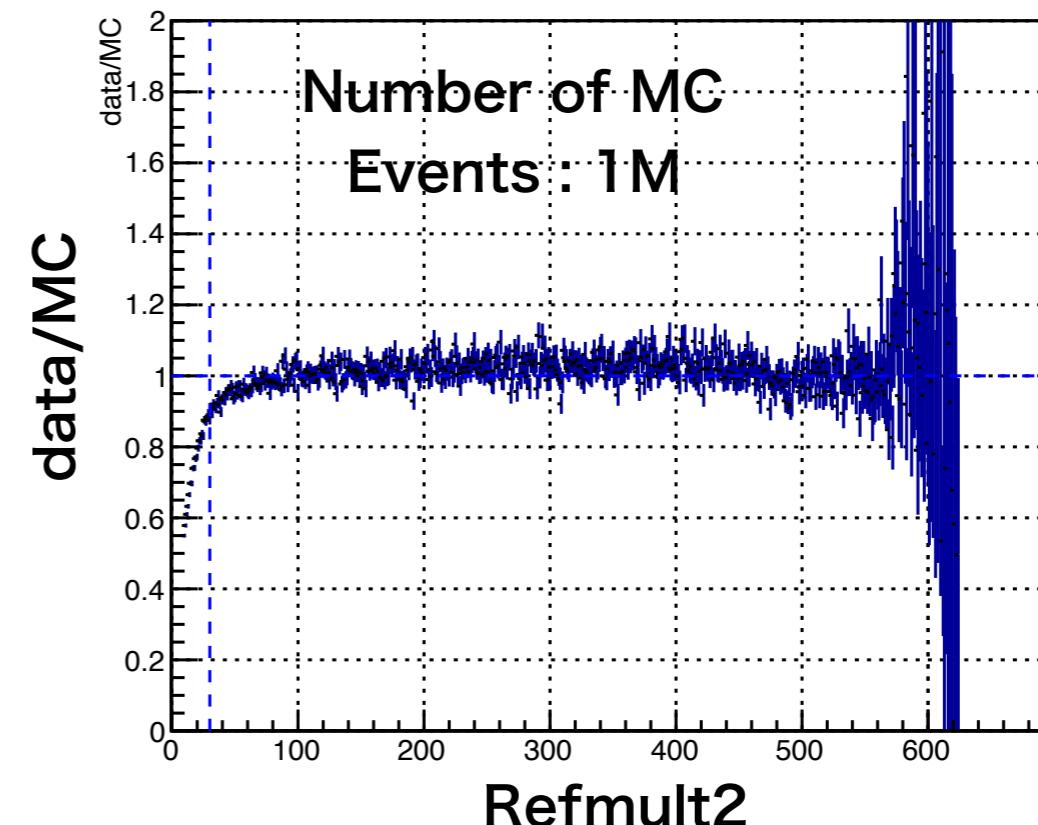
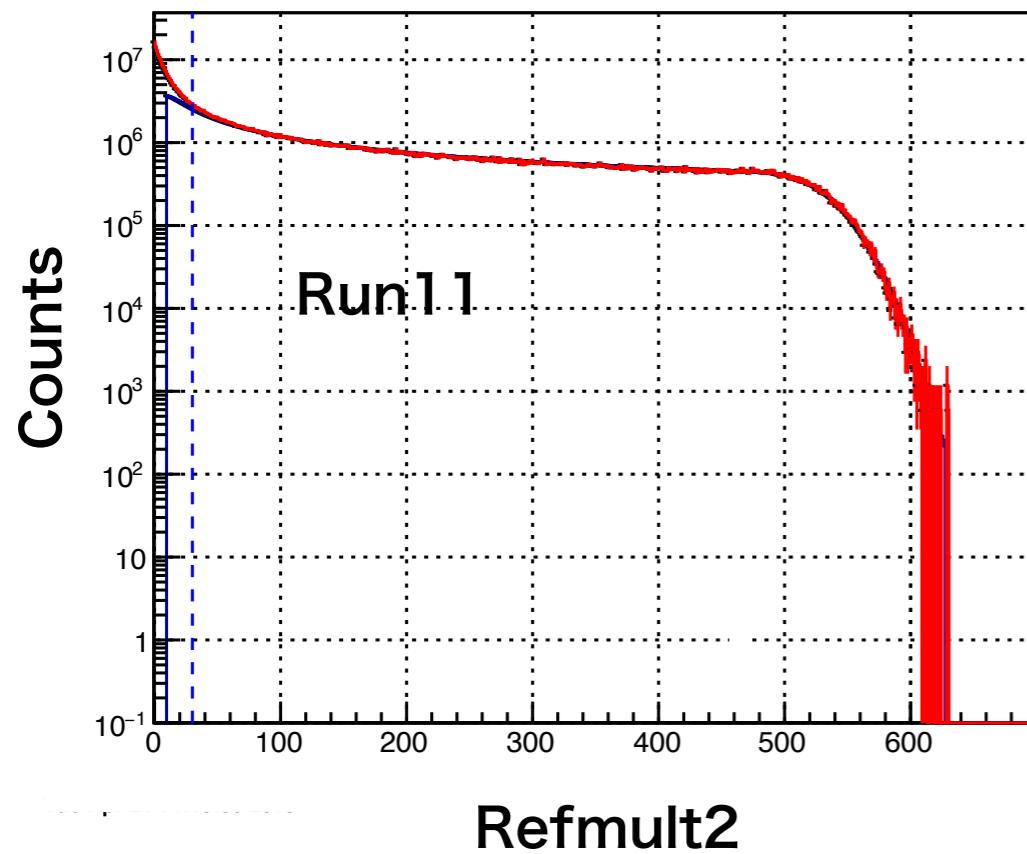
$$(1-x) * (N_w / 2.) + x * N_{coll}$$

- : initial source $x=0.13$
- : final state multiplicity



- Final state particles are generated from each source independently. (IPP)
- Number of particles from each source is fluctuating under the NBD.

Glauber fit results

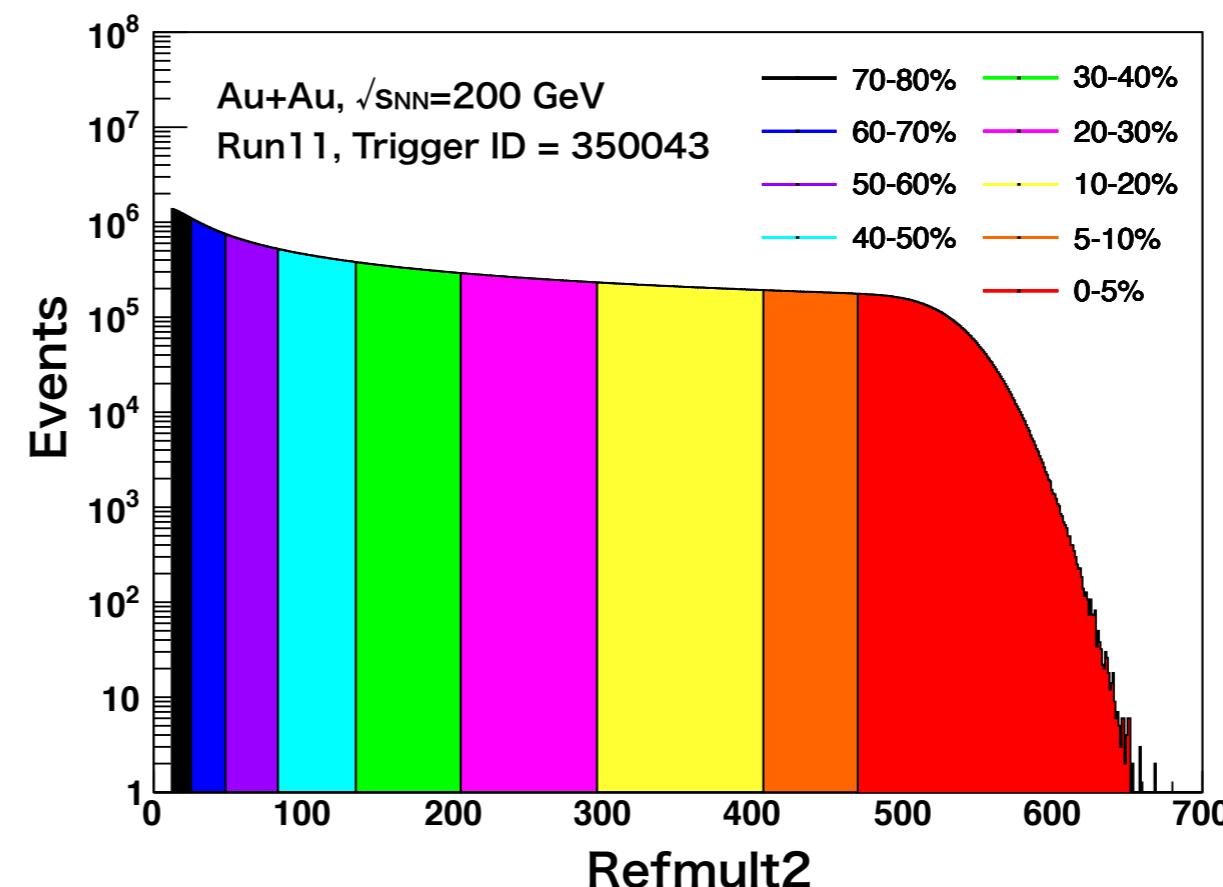


Fitting parameters

n_{pp} : Mean number of generated particles from each source.

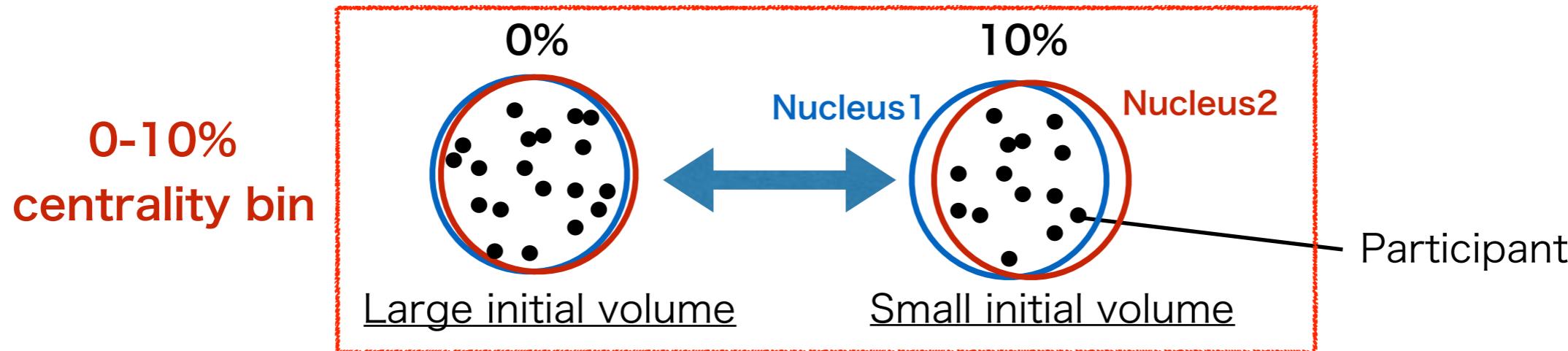
k : Parameters of NBD.

eff : tracking efficiency



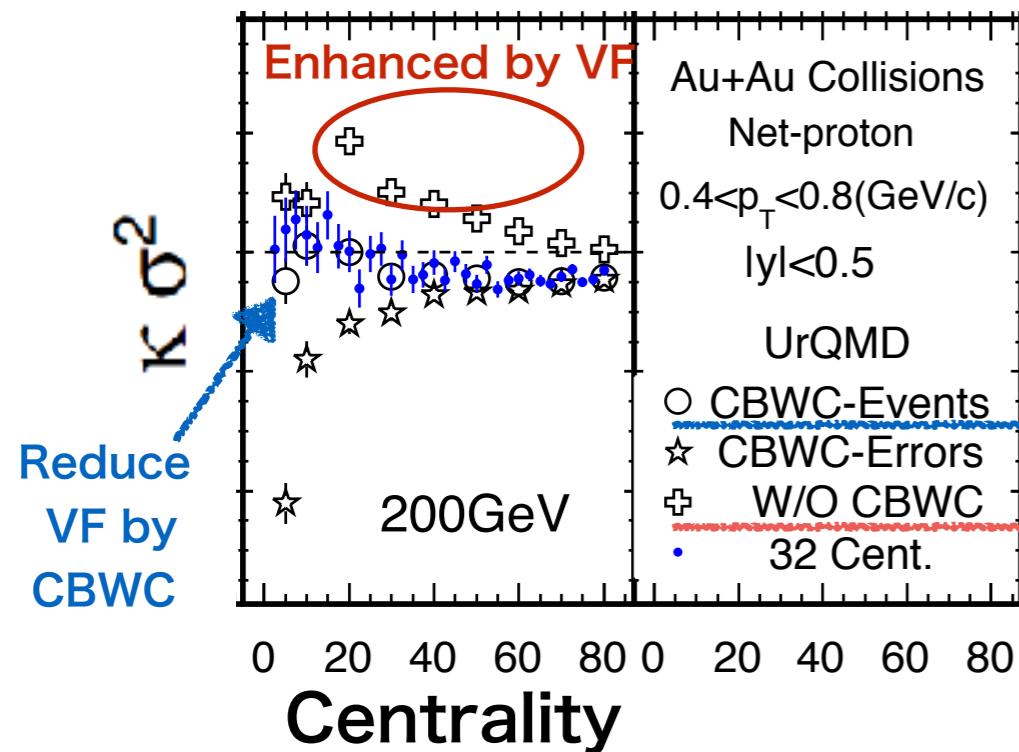
Centrality Bin Width Correction

- Initial volume which correspond to number participant nucleons (N_w) are different even in the same centrality bins.



- Because of VF, cumulants are artificially enhanced which would be the background which should be eliminated.

Xiaofeng Luo, J. Phys. G: Nucl. Part. Phys. 40 105104 (2013)



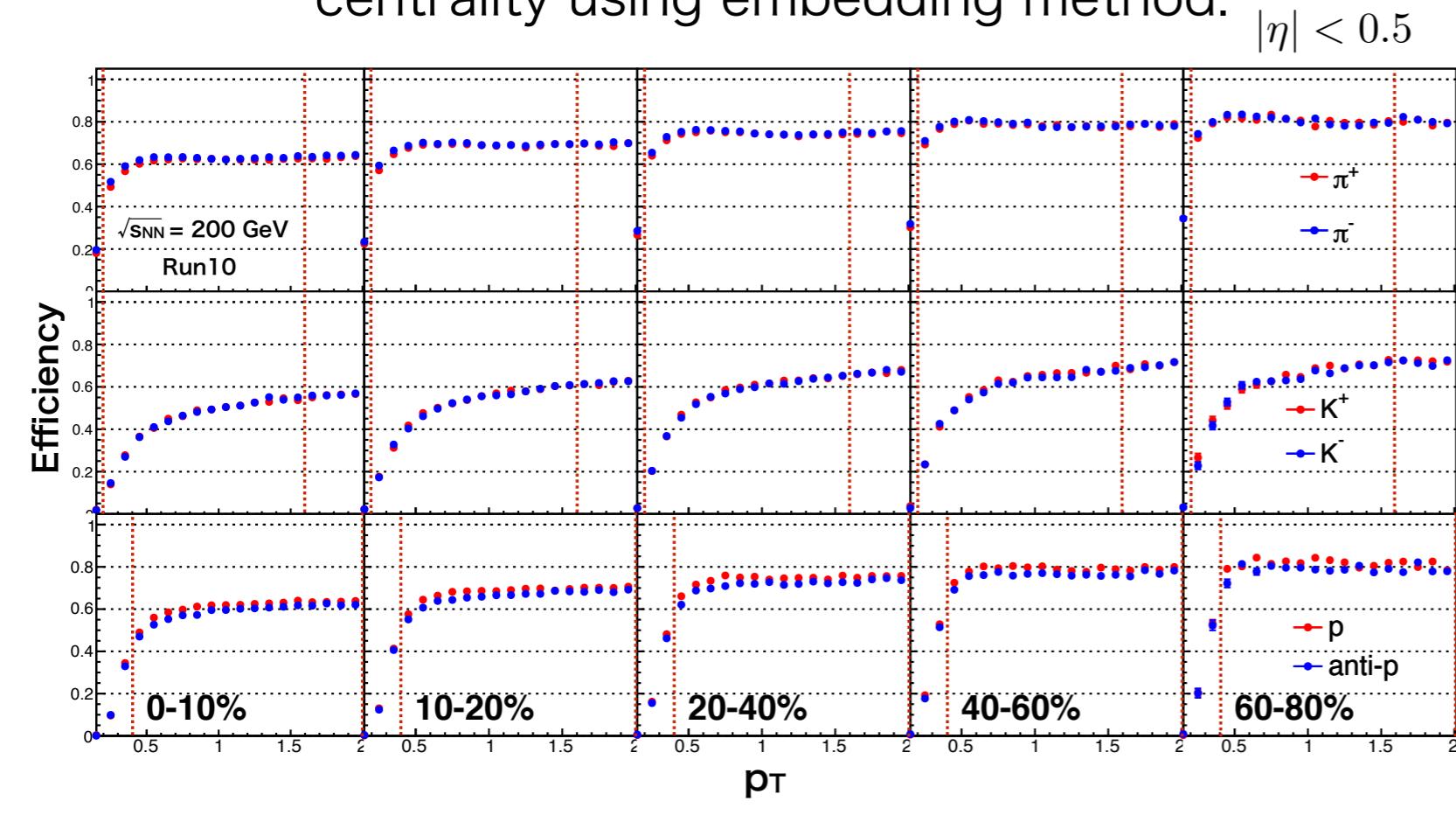
- STAR applied **Centrality Bin Width Correction (CBWC)**.
- In this correction, cumulants are **calculated for each multiplicity bins**.

$$C_n = \sum_r w_r C_{(n,r)} \quad w_r = \frac{N_r}{\sum_r N_r}$$

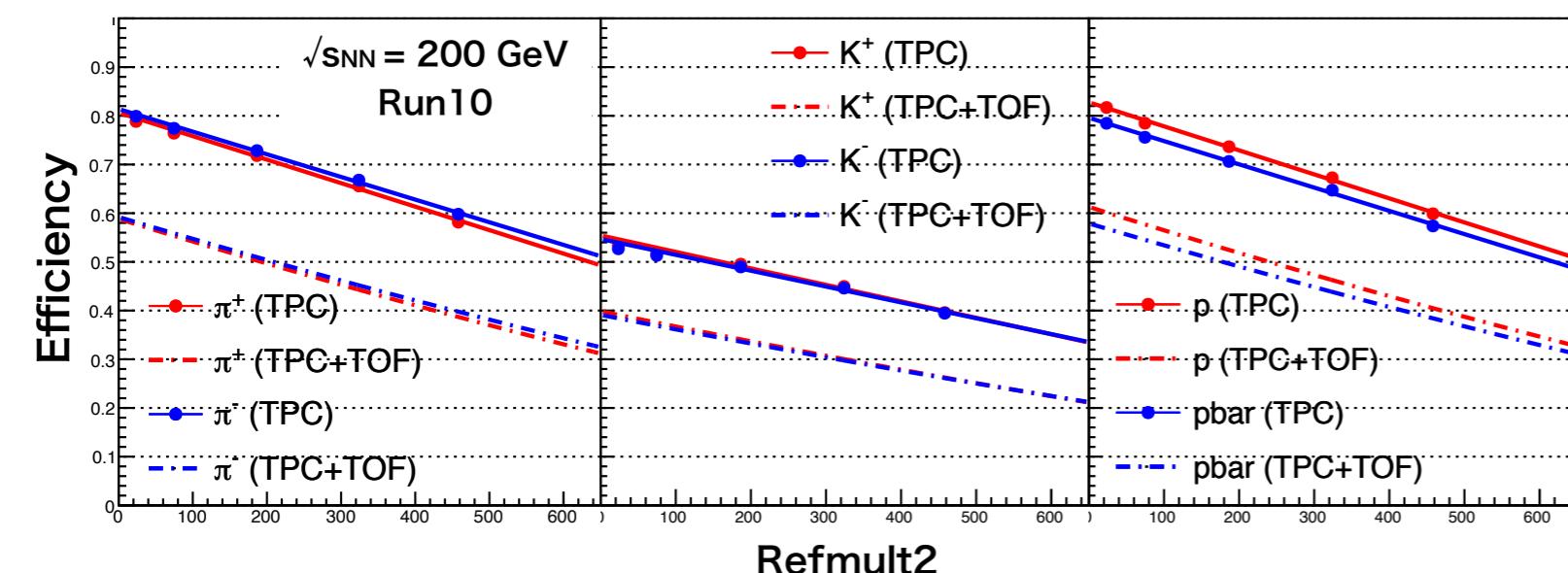
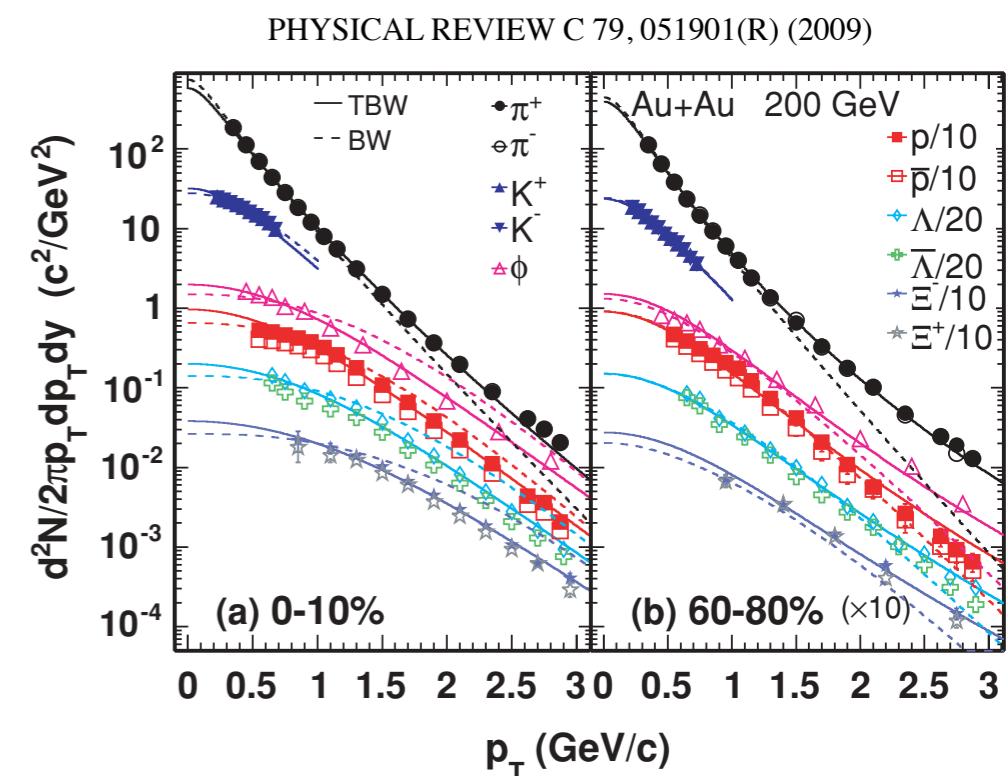
Cumulants for each centrality bins Weight Cumulants in r-th multiplicity bins

Efficiency estimation

① p_T dependent efficiencies were estimated for each centrality using embedding method.



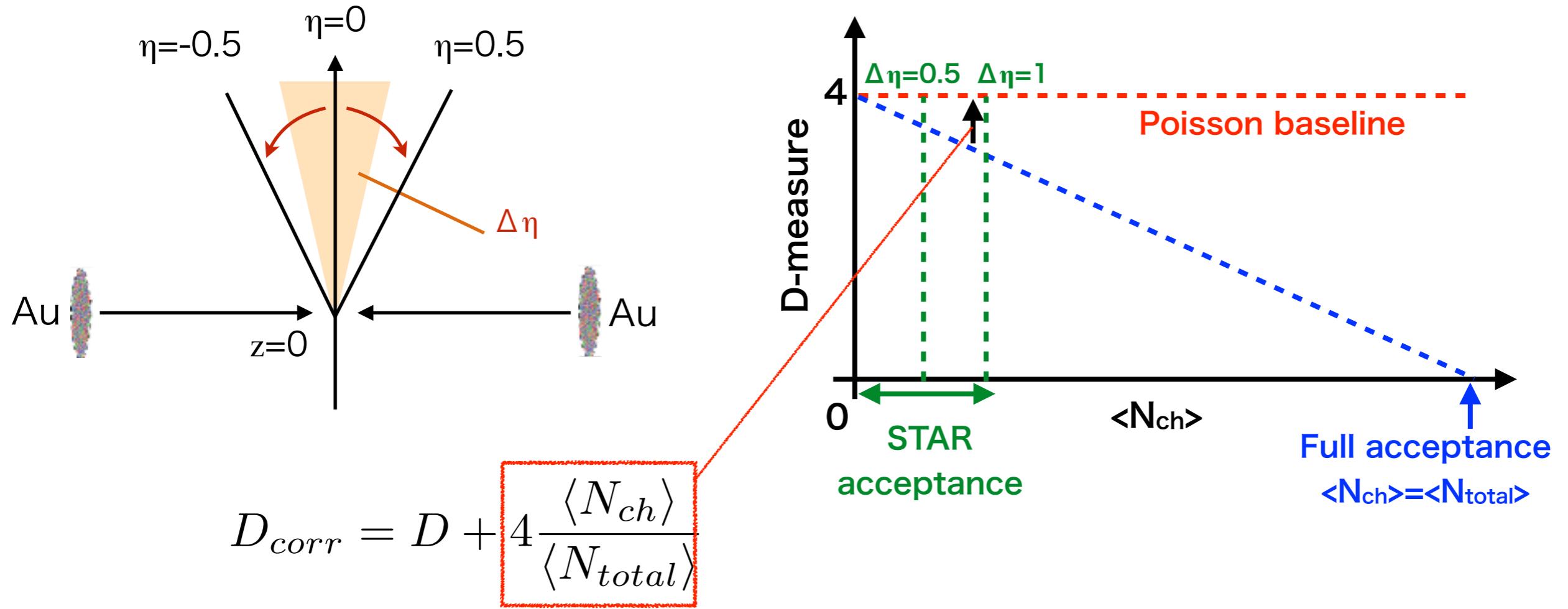
② p_T integrated efficiencies were estimated using results of ① and spectra data.



$$\epsilon_i = \frac{\int \epsilon_i(p_T) p_T f(p_T) dp_T}{\int p_T f(p_T) dp_T}$$

③ poll fit and estimate tracking efficiencies for each multiplicity bin.

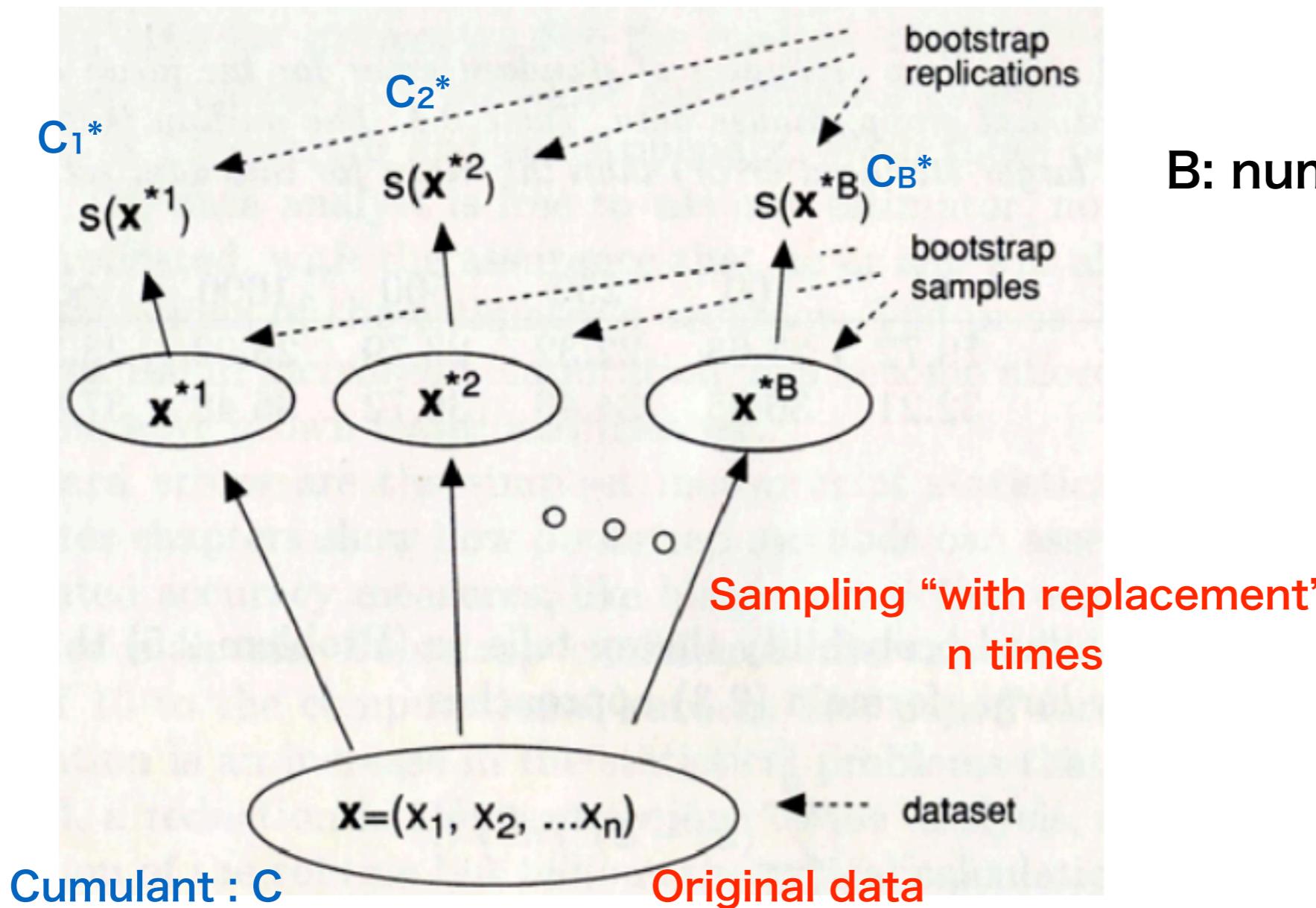
Charge conservation correction on D



- D-measure decrease with expanding $\Delta\eta$ due to charge conservation.
- This effect is critical to measure $\Delta\eta$ dependence of D-measure
- Corrected by adding correction term.
- N_{total} is total multiplicity in full acceptance and estimated from PHOBOS experiment results.

Statistical error estimation

- Statistical errors are estimated by **Bootstrap** method.



-
- Error of C_n is estimated by calculating **standard error of C_n^* distribution**.
 - Number of bootstrap is more than 100 times in this analysis.

How to estimate Systematic uncertainty?

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$$RMS = \sqrt{\frac{1}{n} \sum_i \left(\frac{Y_i - Y_{st.cut}}{Y_{st.cut}} \right)^2}$$
$$Sys.Err = Y_{st.cut} \sqrt{\sum_j (RMS)^2}$$

Net-charge (published)

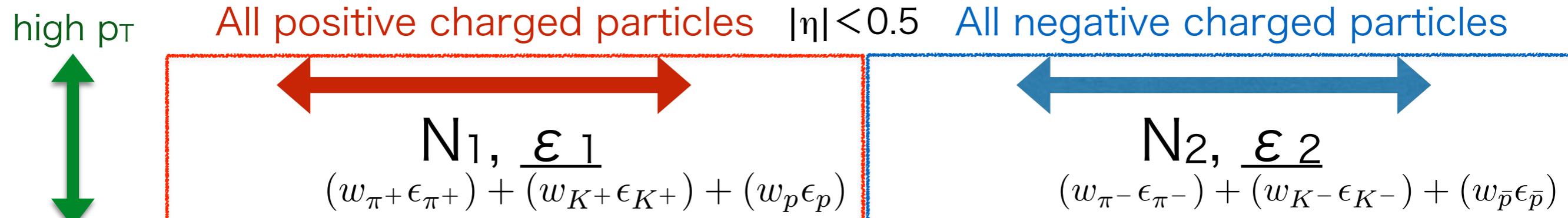
dca	0.8, 1.0(default), 1.2
nFitPoints	-18, 20(default), 22
nhitsdedx	8, 10(default), 12
efficiency	$\pm 5\%$ + α?

- Efficiency of positively charged particle and negatively charged particle are changed $\pm 5\%$ at published results.
(pos,neg) = (+5%, +5%) and (-5%, -5%)
- In this analysis, efficiencies are changed 0.3% separately in addition to conventional method.
(pos,neg) =(+5%, +5%), (-5%, -5%), (+0.3%, -0.3%), (+0.3%, -0.3%)

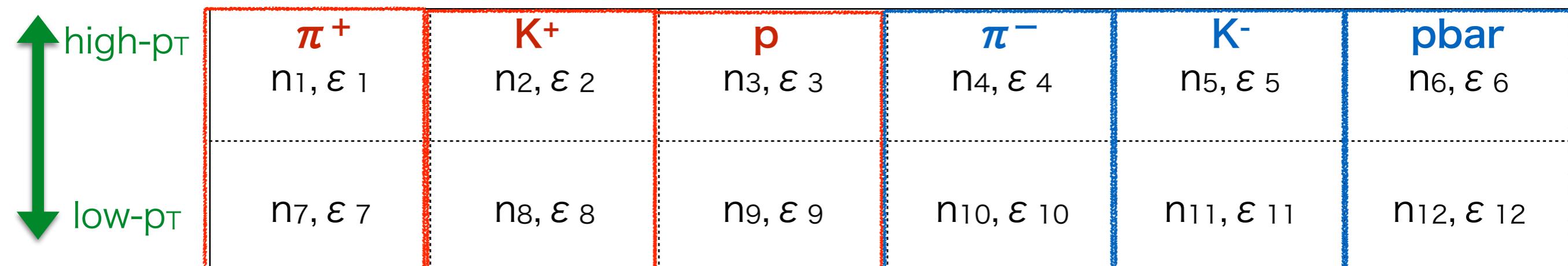
Analysis and correction method improvement (C₆ analysis)

Difference between published and current method²⁹

- Earlier method



- Current method



Net-charge : $(n_1 + n_2 + n_3 + n_7 + n_8 + n_9) - (n_4 + n_5 + n_6 + n_{10} + n_{11} + n_{12})$

\uparrow \uparrow \uparrow ...

Efficiency : $\epsilon_1 \quad \epsilon_2 \quad \epsilon_3 \quad \dots$ 12 efficiency bins were used

- **Earlier method** : Count all charged particle with averaged efficiency.
- **Current method** : Count π , K , p with separated efficiencies.

Efficiency correction for different phase space³⁰

It was suggested that efficiency correction using the average efficiency for different region of the phase space (different p_T bins for different particle species) would give artificial results when distributions do not follow Poisson.

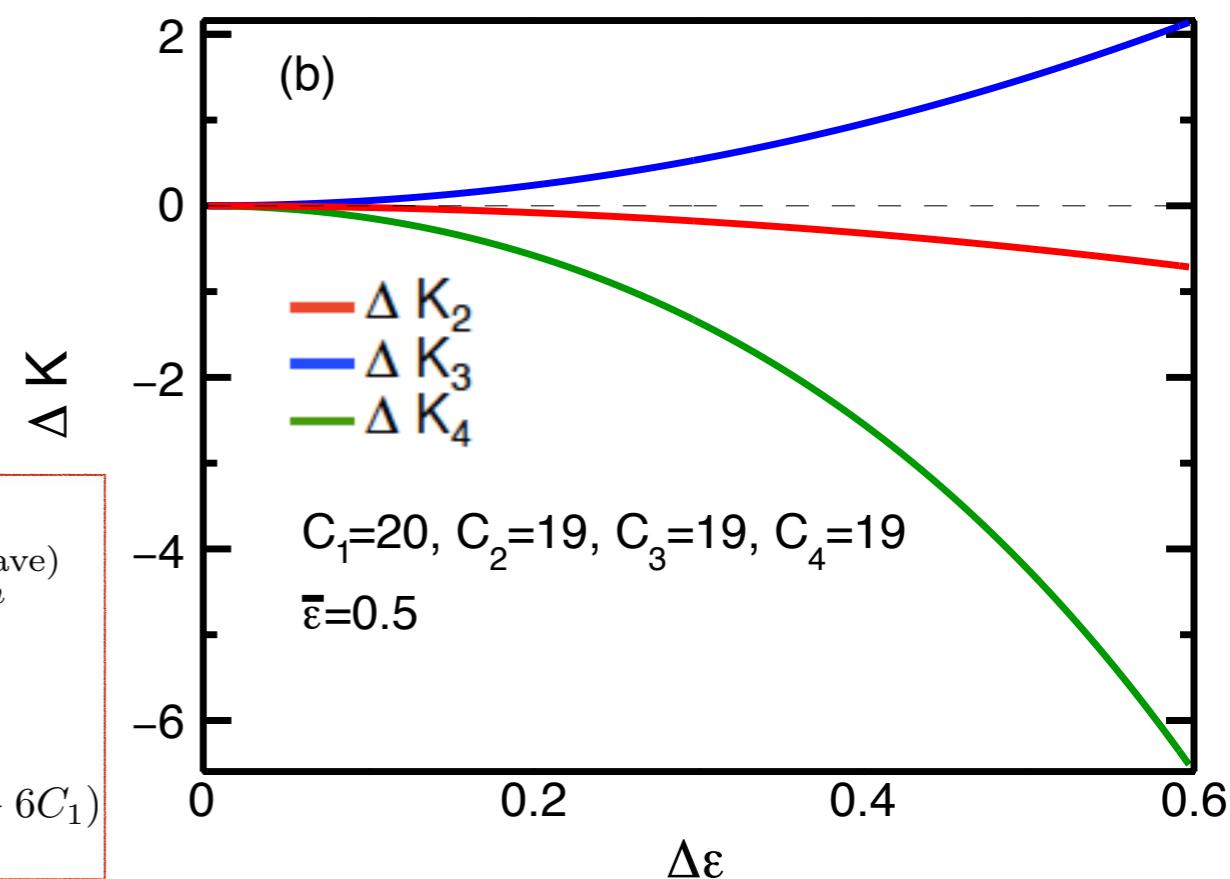
2 distribution model

n_A, ϵ_A n_B, ϵ_B
 $\Delta\epsilon = \epsilon_A - \epsilon_B$

5% smaller cumulants than Poisson

$$\Delta K_2 = \frac{1}{2} \left(\frac{\Delta\epsilon}{\bar{\epsilon}} \right)^2 (C_2 - C_1), \quad \Delta K_m = K_m - K_m^{(\text{ave})} = 2C_m - K_m^{(\text{ave})}$$
$$\Delta K_3 = \frac{3}{2} \left(\frac{\Delta\epsilon}{\bar{\epsilon}} \right)^2 (C_3 - 2C_2 + C_1),$$
$$\Delta K_4 = \frac{1}{2} \left(\frac{\Delta\epsilon}{\bar{\epsilon}} \right)^2 (6C_4 - 18C_3 + 19C_2 - 7C_1) + \frac{1}{8} \left(\frac{\Delta\epsilon}{\bar{\epsilon}} \right)^4 (C_4 - 6C_3 + 11C_2 - 6C_1)$$

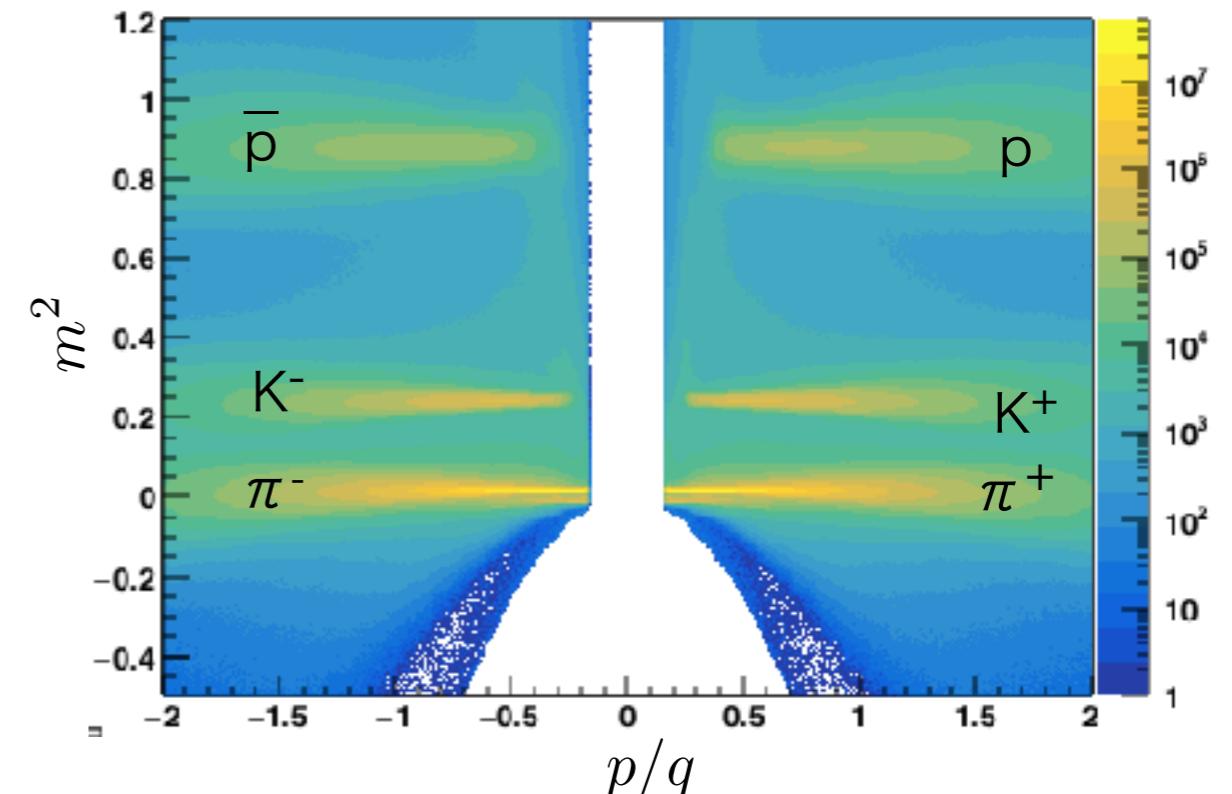
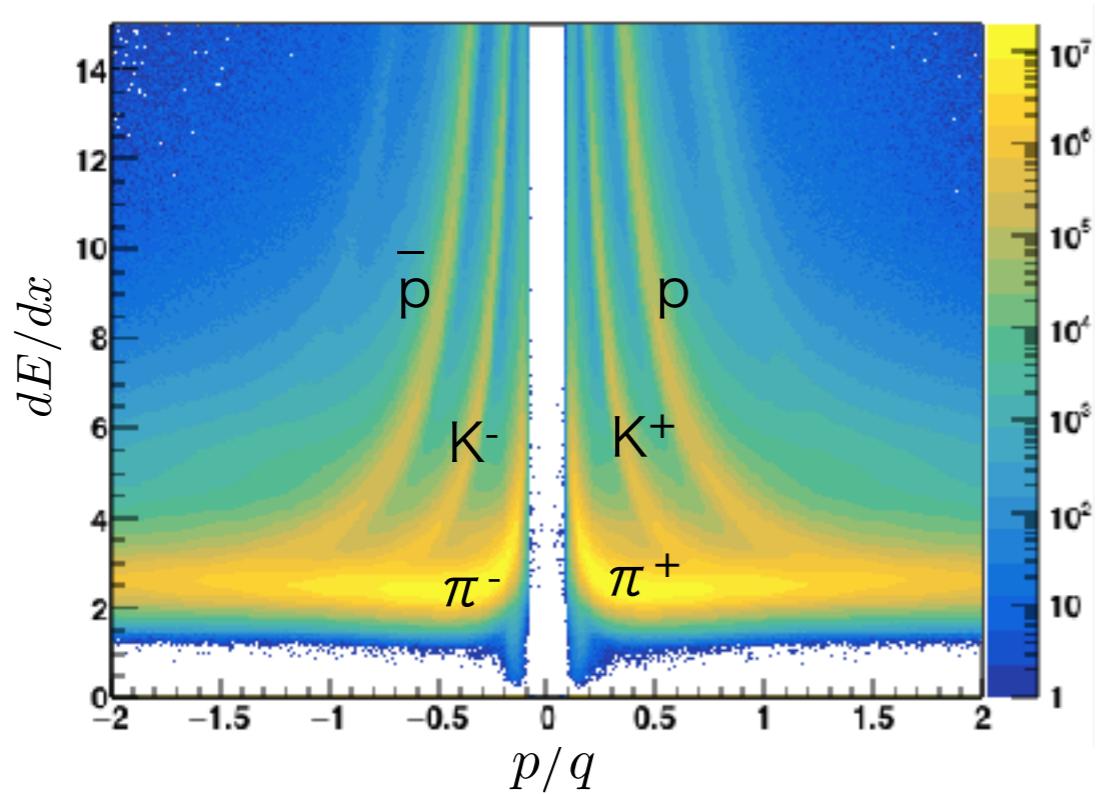
Phys. Rev. C 95, 064912 T.Nonaka, M.Kitazawa and S.Esumi



We can't estimate true corrected cumulant when
uses average efficiencies if $\Delta\epsilon$ is large.

Particle Identification

Au+Au, $\sqrt{s_{NN}} = 200\text{GeV}$
STAR Preliminary



- PID have been done by using dE/dx of TPC
- TOF is also used for PID at high p_T region

Track cut	π	K	p
p_T (TPC only)	0.2 - 0.5 GeV/c	0.2 - 0.4 GeV/c	0.4 - 0.8 GeV/c
p_T (TPC+TOF)	0.5 - 1.6 GeV/c	0.4 - 1.6 GeV/c	0.8 - 2.0 GeV/c

Factorial cumulant method

Cumulant (measured)

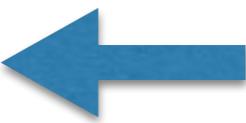


Factorial Cumulant (measured)



Efficiency correction

Cumulant (true)



Factorial Cumulant (true)

1-4th-order cumulants (factorial cumulant formula)

$$\langle Q \rangle_c = \langle q_{(1,1)} \rangle_c,$$

$$\langle Q^2 \rangle_c = \langle q_{(1,1)}^2 \rangle_c + \langle q_{(2,1)} \rangle_c - \langle q_{(2,2)} \rangle_c,$$

$$\langle Q^3 \rangle_c = \langle q_{(1,1)}^3 \rangle_c + 3\langle q_{(1,1)}q_{(2,1)} \rangle_c - 3\langle q_{(1,1)}q_{(2,2)} \rangle_c + \langle q_{(3,1)} \rangle_c - 3\langle q_{(3,2)} \rangle_c + 2\langle q_{(3,3)} \rangle_c,$$

$$\begin{aligned} \langle Q^4 \rangle_c = & \langle q_{(1,1)}^4 \rangle_c + 6\langle q_{(1,1)}^2 q_{(2,1)} \rangle_c - 6\langle q_{(1,1)}^2 q_{(2,2)} \rangle_c + 4\langle q_{(1,1)} q_{(3,1)} \rangle_c + 3\langle q_{(2,1)}^2 \rangle_c \\ & + 3\langle q_{(2,2)}^2 \rangle_c - 12\langle q_{(1,1)} q_{(3,2)} \rangle_c + 8\langle q_{(1,1)} q_{(3,3)} \rangle_c - 6\langle q_{(2,1)} q_{(2,2)} \rangle_c \\ & + \langle q_{(4,1)} \rangle_c - 7\langle q_{(4,2)} \rangle_c + 12\langle q_{(4,3)} \rangle_c - 6\langle q_{(4,4)} \rangle_c, \end{aligned}$$

Phys. Rev. C 95, 064912

T.Nonaka, M.Kitazawa and S.Esumi

Cumulant (true)

Cumulant (measured)

$$q_{(r,s)} = q_{(a^r / p^s)} = \sum_{i=1}^M (a_i^r / p_i^s) n_i.$$

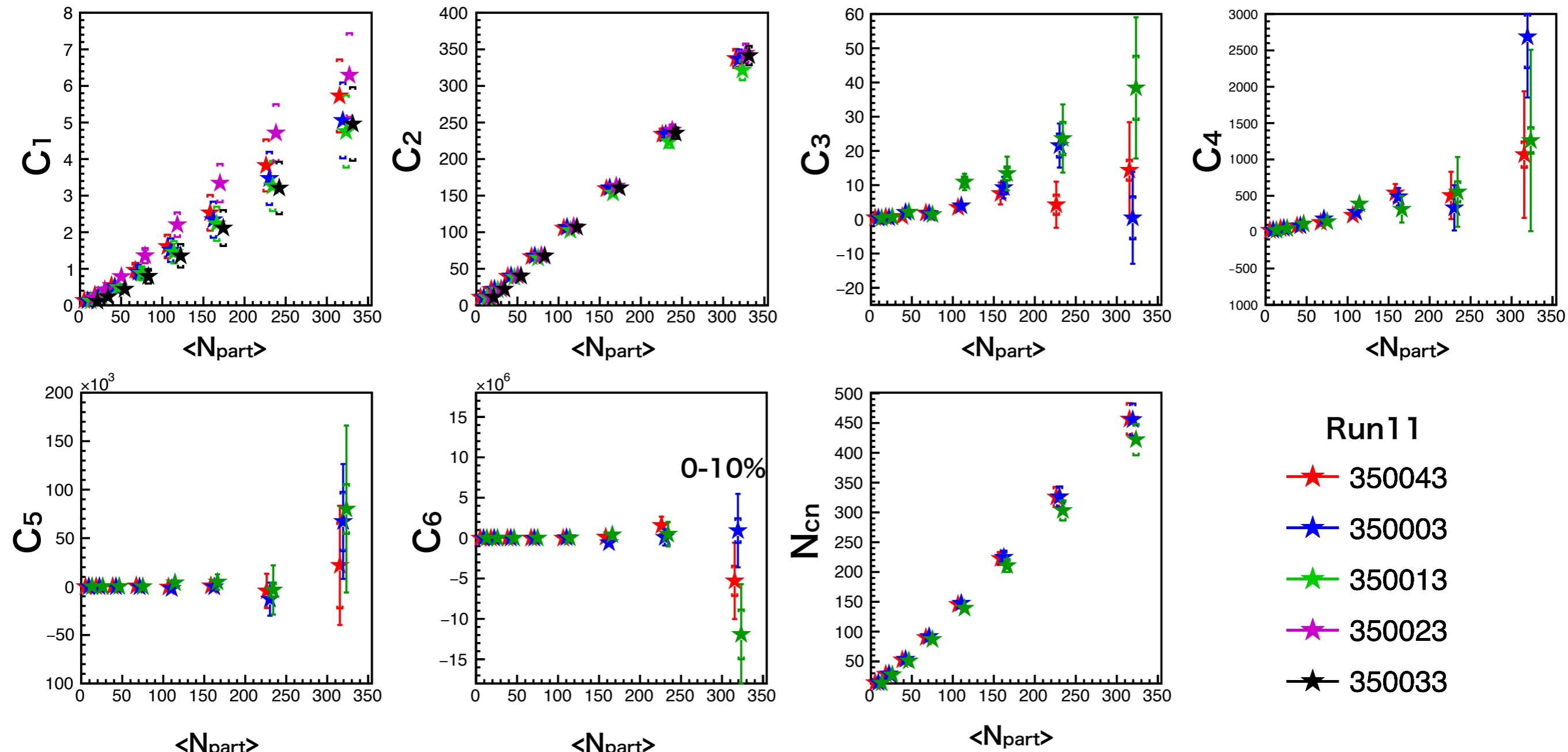
+1 (positive charge) or -1 (negative charge)

M Number of phase space
i=1 Number of particle
Efficiency

Results (C₆ analysis)

Cumulants (200GeV, Run11)

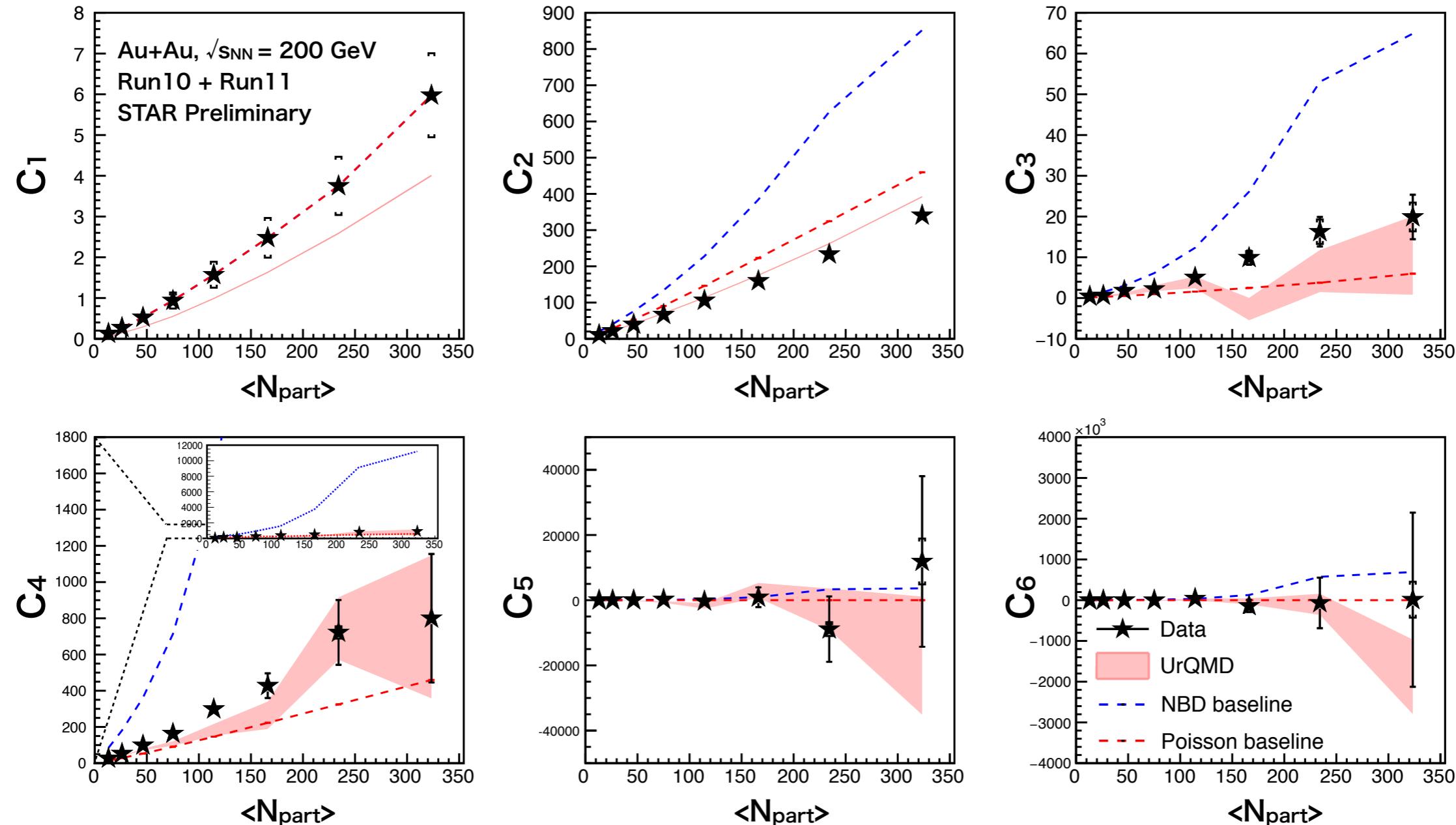
34



- Cumulants of all trigger ID are consistent within statistical or systematic errors.

Cumulants (200GeV, Run10+Run11)

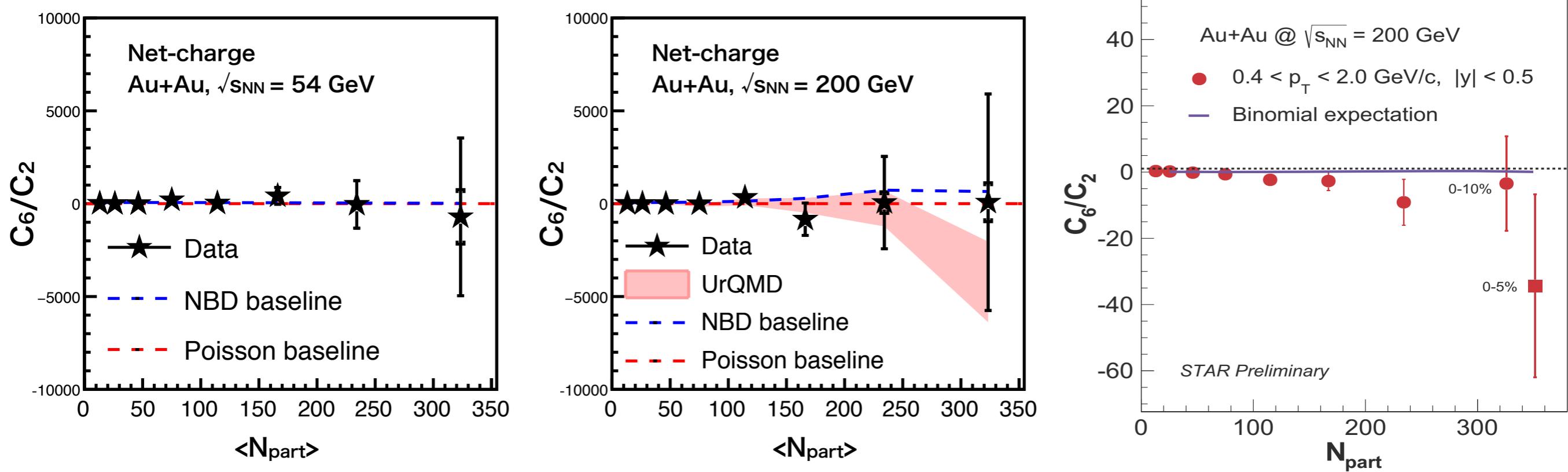
35



- From 1st- to 6th-order cumulants were measured as a function of $\langle N_{\text{part}} \rangle$
- Linear trend can be seen.
- For higher-order cumulants, statistical uncertainties are larger than the systematic ones.
- NBD baseline becomes larger than experimental values especially at C_4 .

C_6/C_2 at 200 GeV and 54GeV

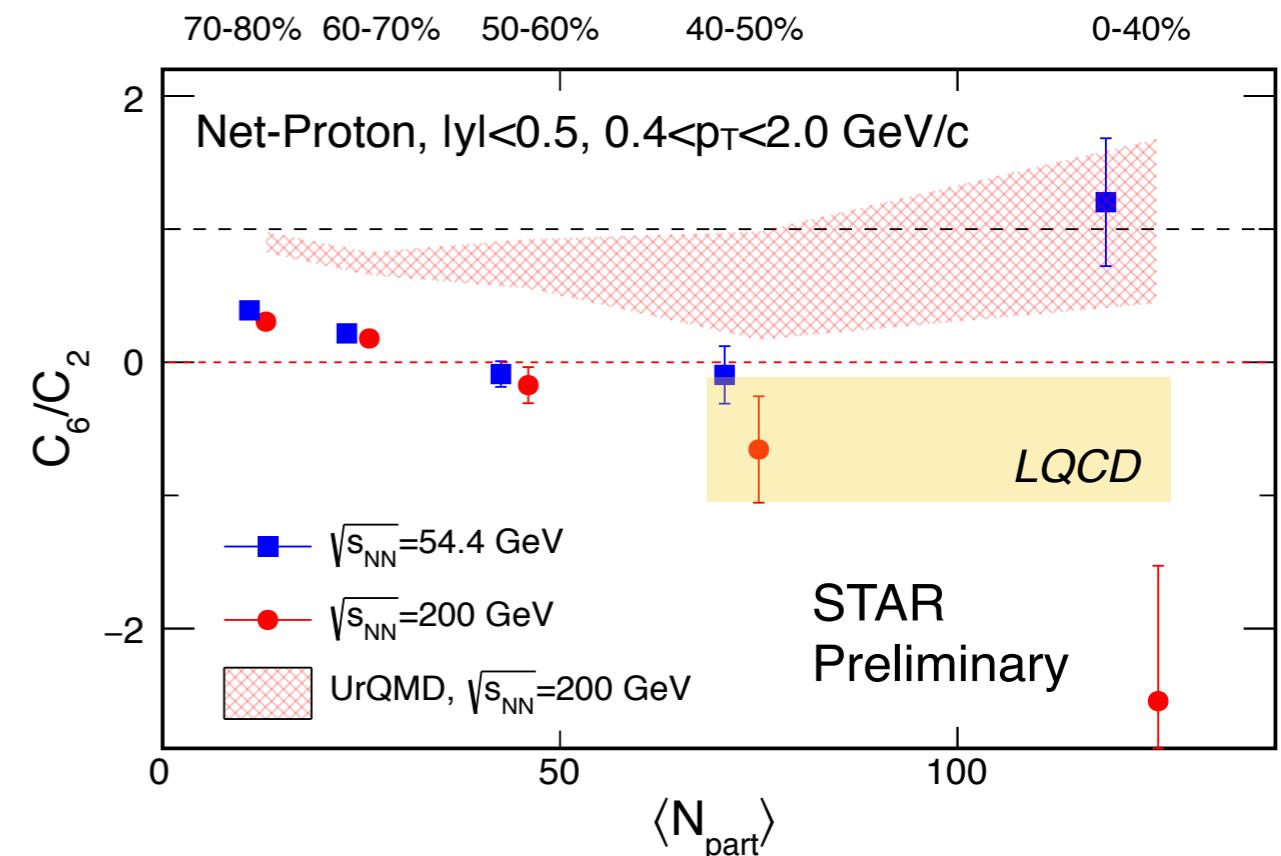
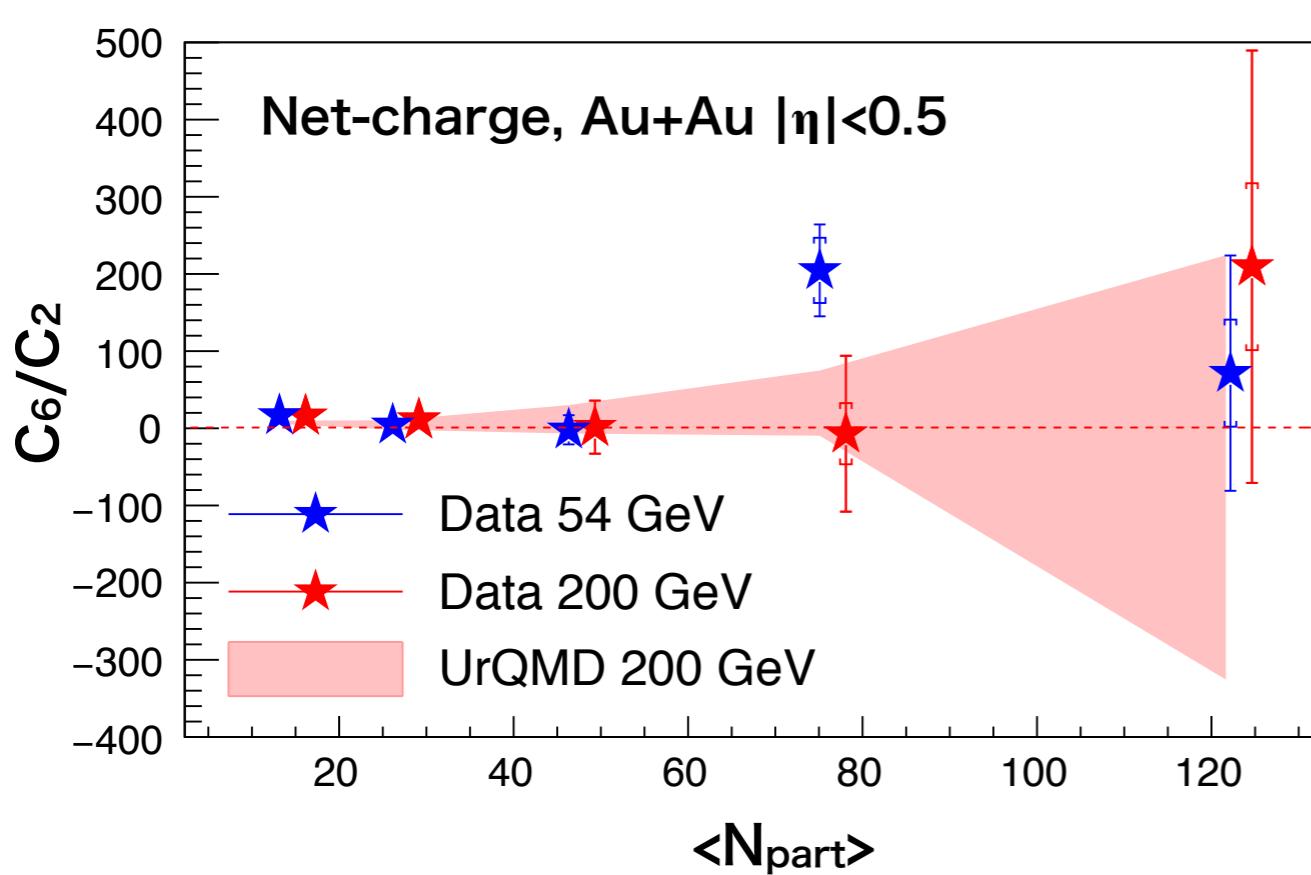
36



- Centrality dependence of C_6/C_2 of net-charge at 200 GeV and 54 GeV are compared to that of net-proton at 200 GeV.
- Compared to Poisson, NBD, and UrQMD.
- All statistical baseline and UrQMD results of C_6/C_2 are consistent within statistical uncertainties and large deviations are not observed.

C₆/C₂ (0-40% merged)

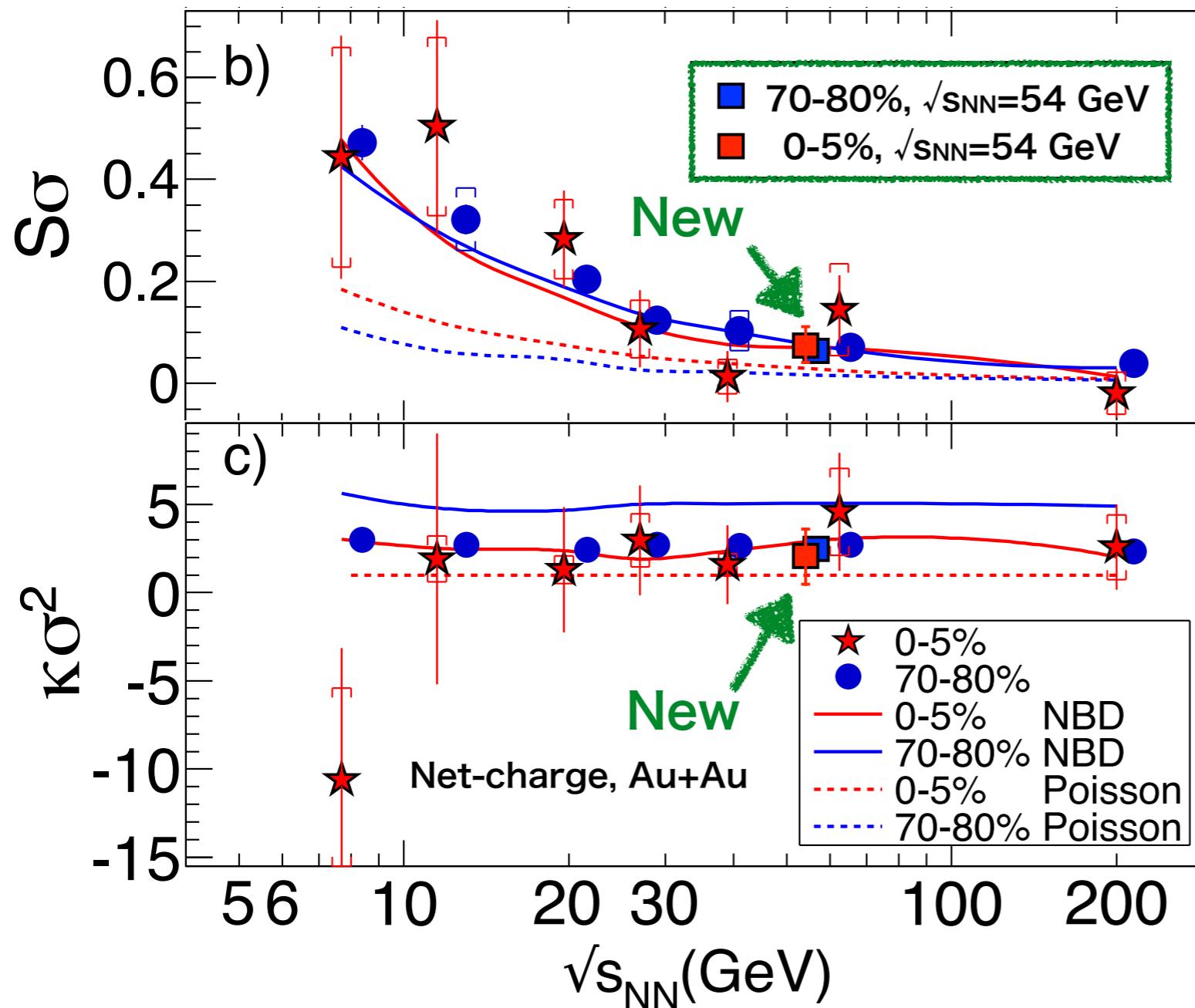
37



- 0-40% centralities are merged.
- Deviation between 200 GeV and 54 GeV is observed in 0-40% centrality at net-proton and 40-50% at net-charge.
→ Signal from cross-over??
- Statistical errors are large and which results should be compare to lower energies in future.

C₃/C₂ and C₄/C₂ at 54 GeV

38



$$S\sigma = \frac{C_3}{C_2} = \frac{\chi_3}{\chi_2}$$

$$\kappa\sigma^2 = \frac{C_4}{C_2} = \frac{\chi_4}{\chi_2}$$

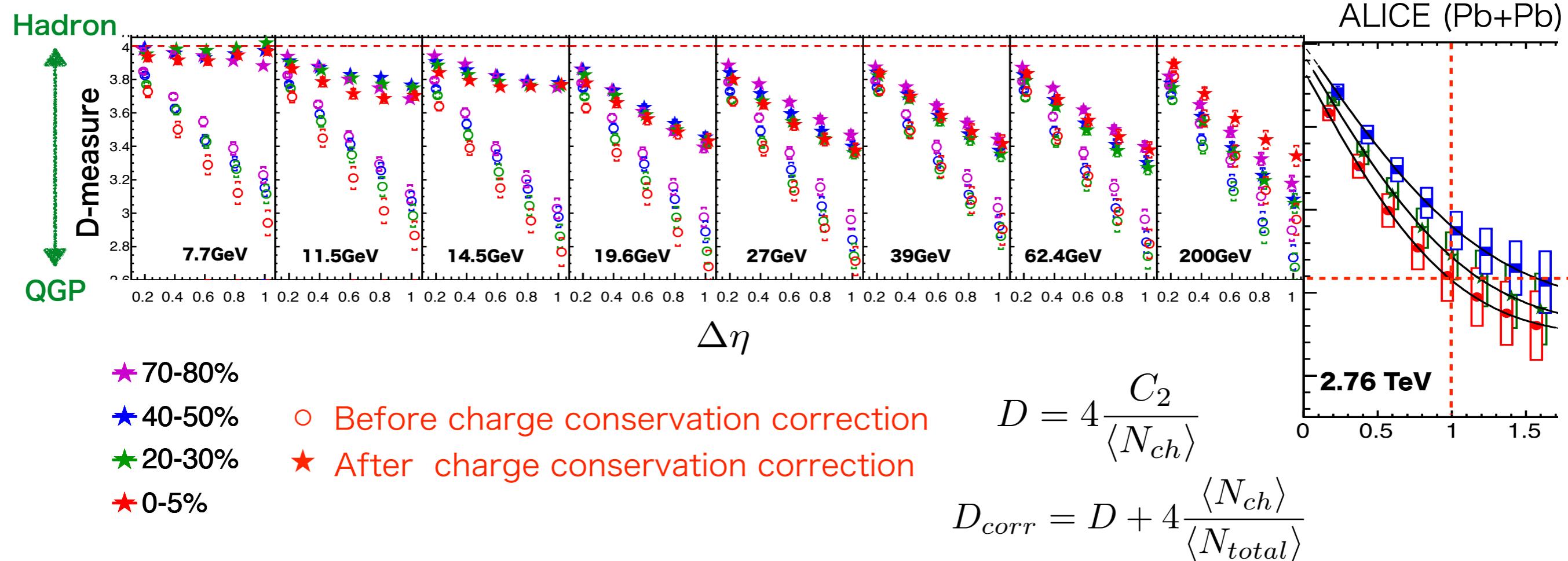
L. Adamczyk *et al.* (STAR Collaboration)
Phys. Rev. Lett. 113, 092301(2014)

- C₃/C₂ and C₄/C₂ of net-charge at 54 GeV are newly measured in addition to BES energies.
- Close to 39 GeV and 62.4 GeV results.
→ Results at 54 GeV are not conflict with published BES-I results.

Results ($\Delta\eta$ dependence)

$\Delta\eta$ dependence of D-measure

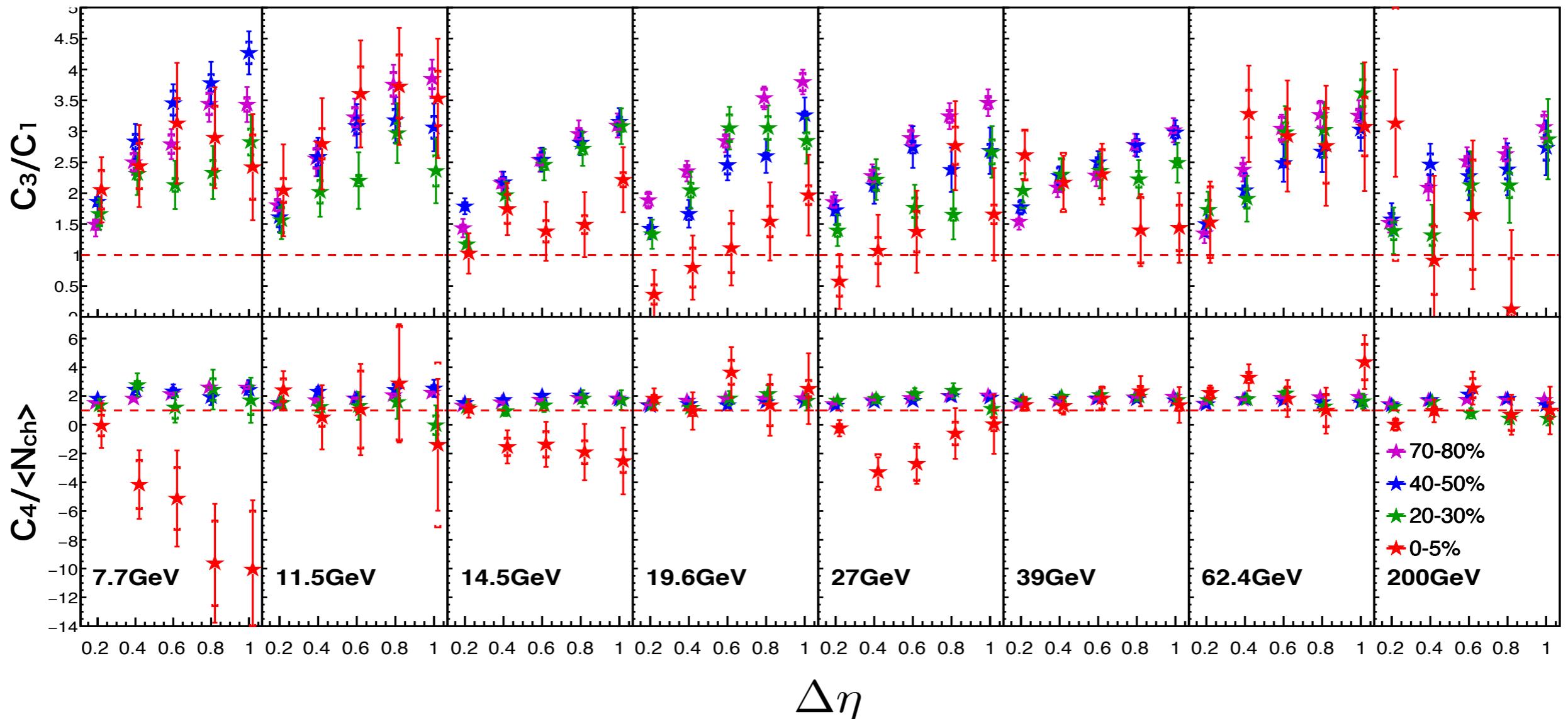
40



- D-measure are observed to decrease with $\Delta\eta$.
- Even though charge conservations are applied, D-measure is decrease with $\Delta\eta$ and this trend is stronger with collision energies.
- These results do not conflict with ALICE and the trends of $\Delta\eta$ dependence might represent the time evolution of phase transition.

Experimental results of C_3/C_1 and $C_4/\langle N_{ch} \rangle$

41



- At C_3/C_1 and $C_4/\langle N_{ch} \rangle$, most of the results are observed to increase with $\Delta\eta$ without most central collision at $C_4/\langle N_{ch} \rangle$.
- These trends can not be described by diffusion process whereas the trends of D-measure can be explained by this model.

Theoretical predictions of C_3/C_1 and $C_4/\langle N_{ch} \rangle$

42

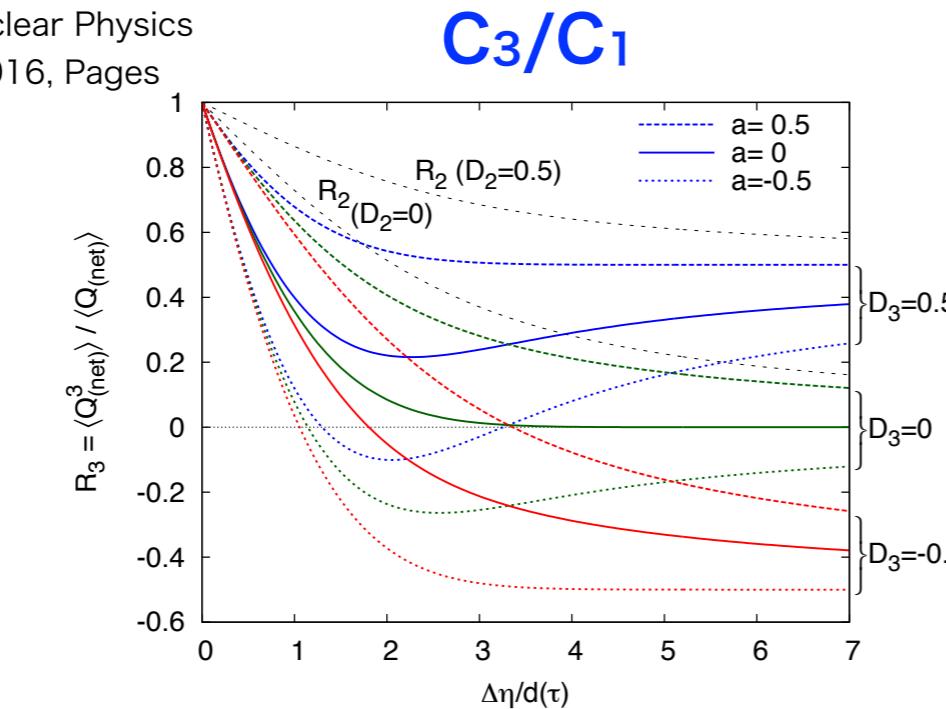
M.Asakawa, M.Kitazaw,

Progress in Particle and Nuclear Physics

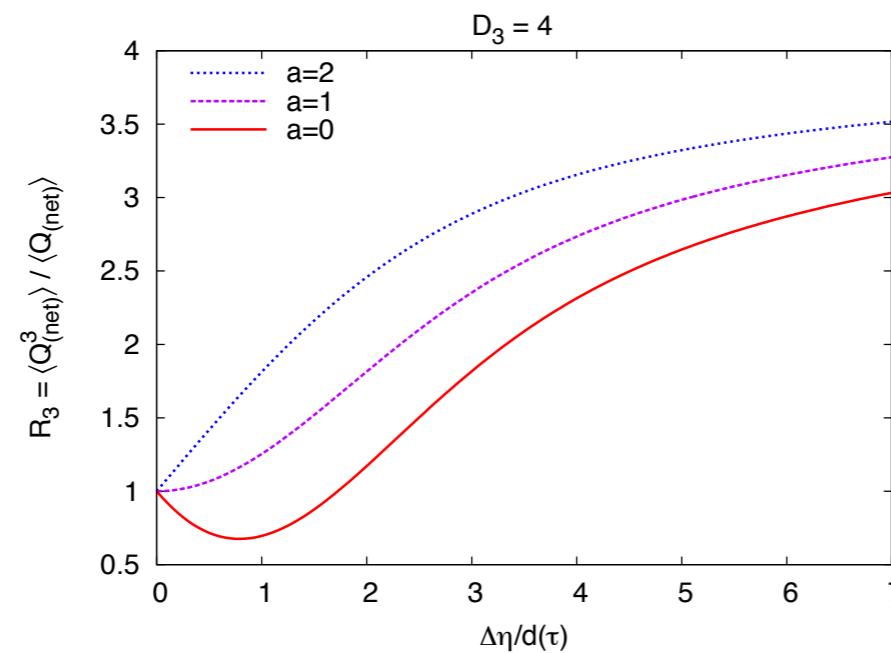
Volume 90, September 2016, Pages

299-342

**Small
susceptibilities**



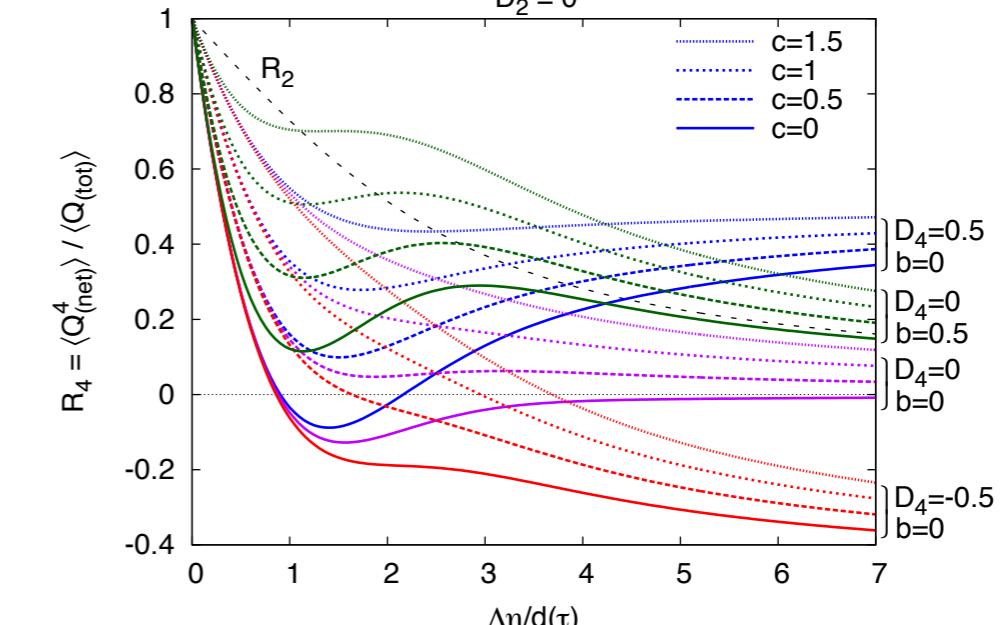
**Large
susceptibilities**



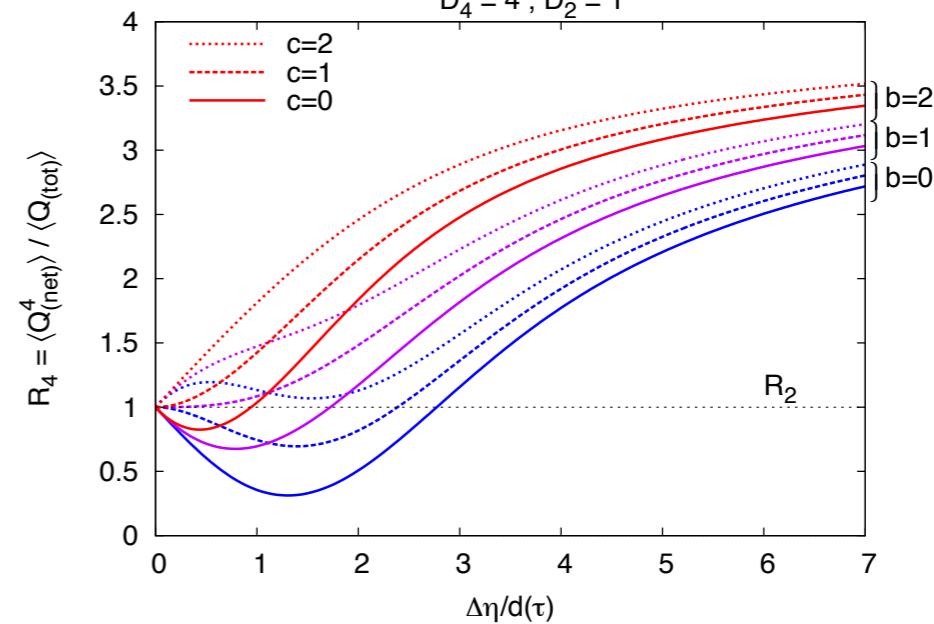
C_3/C_1

$C_4/\langle N_{ch} \rangle$

$D_2 = 0$



$D_4 = 4, D_2 = 1$



- $\Delta\eta$ dependence of **third and fourth-order cumulant ratios are also predicted** by the model but there are a lot of parameters and depend on initial condition.
- Experimental results are close to theoretical expectation with large susceptibility results.

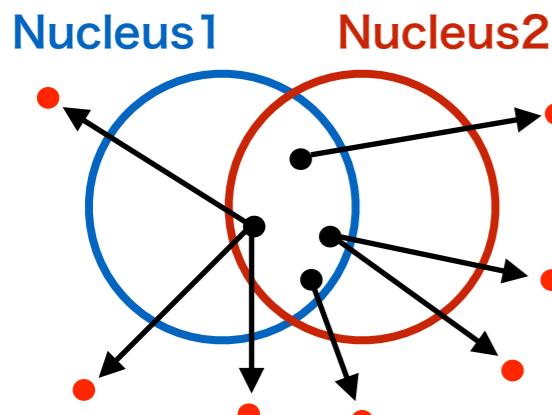
Volume fluctuation correction

Volume fluctuation correction (VFC)

44

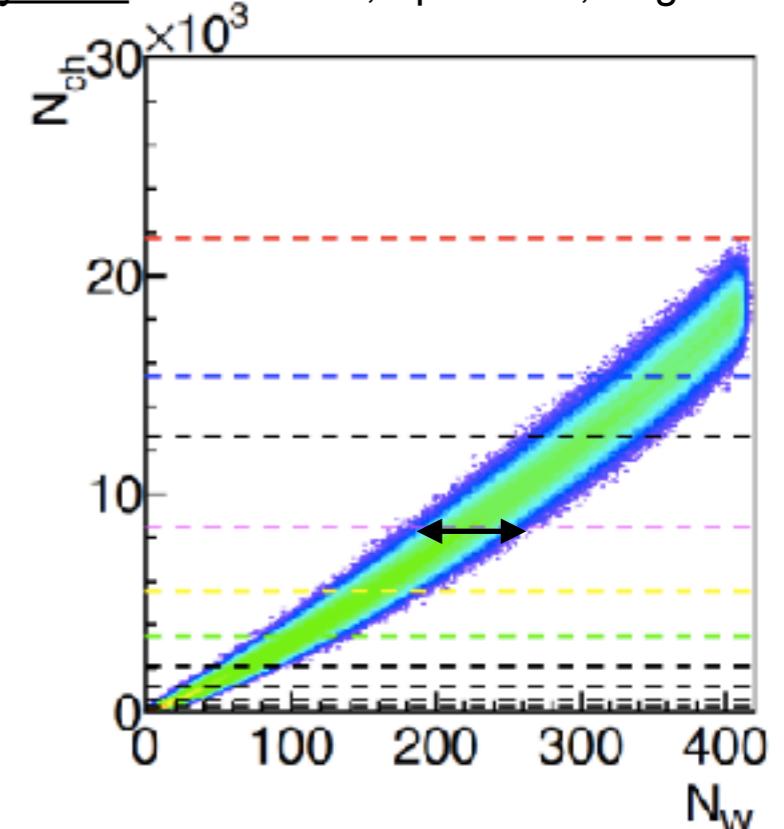
Nuclear Physics A Volume 960, April 2017, Pages 114-130

- N_w is fluctuating even if final state multiplicity is fixed.
→ CBWC may not be enough to eliminate VF



Independent particle production
(IPP) model

- : initial source
- : final state multiplicity



- Under the IPP model, measured cumulants include additional term from VF which can be written by N_w cumulants.
- We can know **true cumulants by subtracting additional terms (VFC)**

Measured cumulant True cumulant

$$\begin{aligned} \kappa_1(\Delta N) &= \langle N_w \rangle \kappa_1(\Delta n) \\ \kappa_2(\Delta N) &= \langle N_w \rangle \kappa_2(\Delta n) + \langle (\Delta n)^2 \rangle \kappa_2(N_w) \\ \kappa_3(\Delta N) &= \langle N_w \rangle \kappa_3(\Delta n) + 3\langle (\Delta n) \rangle \kappa_2(\Delta n) \kappa_2(N_w) + \langle (\Delta n)^3 \rangle \kappa_3(N_w) \\ \kappa_4(\Delta N) &= \langle N_w \rangle \kappa_4(\Delta n) + 4\langle (\Delta n) \rangle \kappa_3(\Delta n) \kappa_2(N_w) \\ &\quad + 3\kappa_2^2(\Delta n) \kappa_2(N_w) + 6\langle (\Delta n)^2 \rangle \kappa_2(\Delta n) \kappa_3(N_w) + \langle (\Delta n)^4 \rangle \kappa_4(N_w) \end{aligned}$$

Δn : number of net-particle

Additional term caused from N_w fluctuation

Toy model method

- Generate 2 independent Poisson distribution
- Parameter : λ_+, λ_-

Glauber : 500M events

Poisson : 100M events

$$\begin{array}{l}
 \lambda_+ * \langle N_w \rangle = N_+ \\
 \lambda_- * \langle N_w \rangle = N_- \\
 (\lambda_+ - \lambda_-) * \langle N_w \rangle = (N_+ - N_-)
 \end{array}$$

Number of particles

\leftarrow

N_w Calculated from Glauber

- λ_+ and λ_- are determined to N_+ and N_- describe the real experiment.

Nw fixed : Generate Poisson $\langle N_w \rangle$ times.

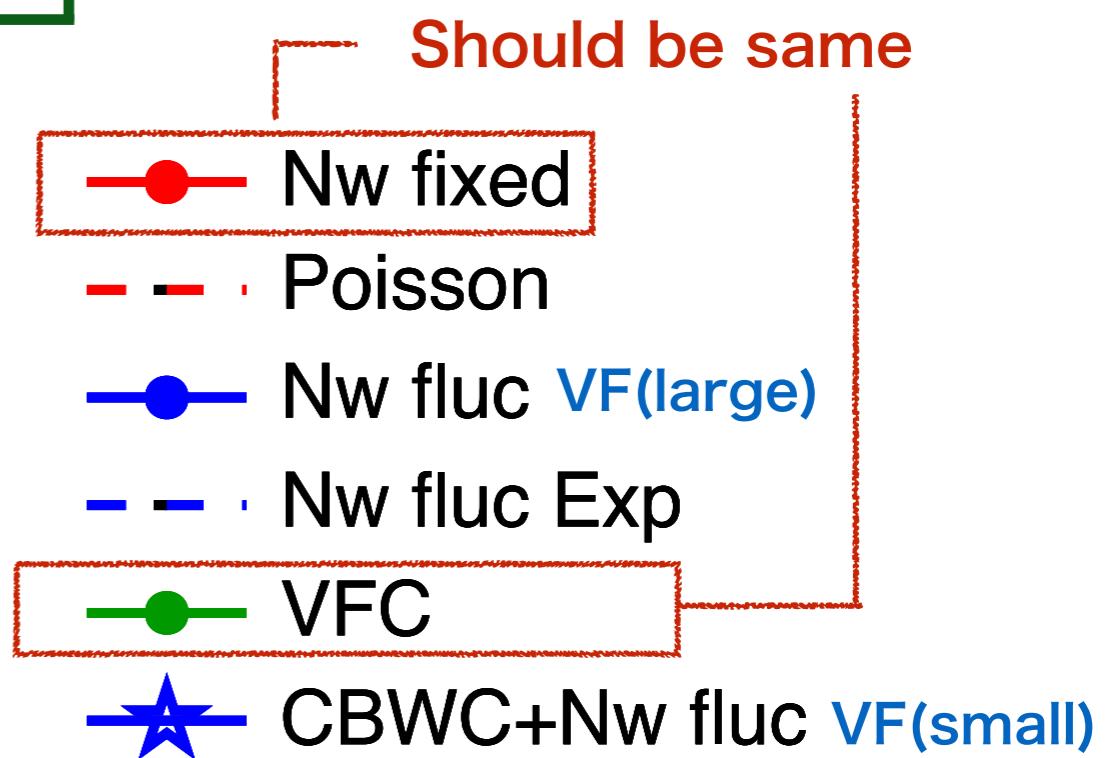
…No N_w fluctuation

Nw fluc : Generate Poisson N_w times.

… N_w fluctuation

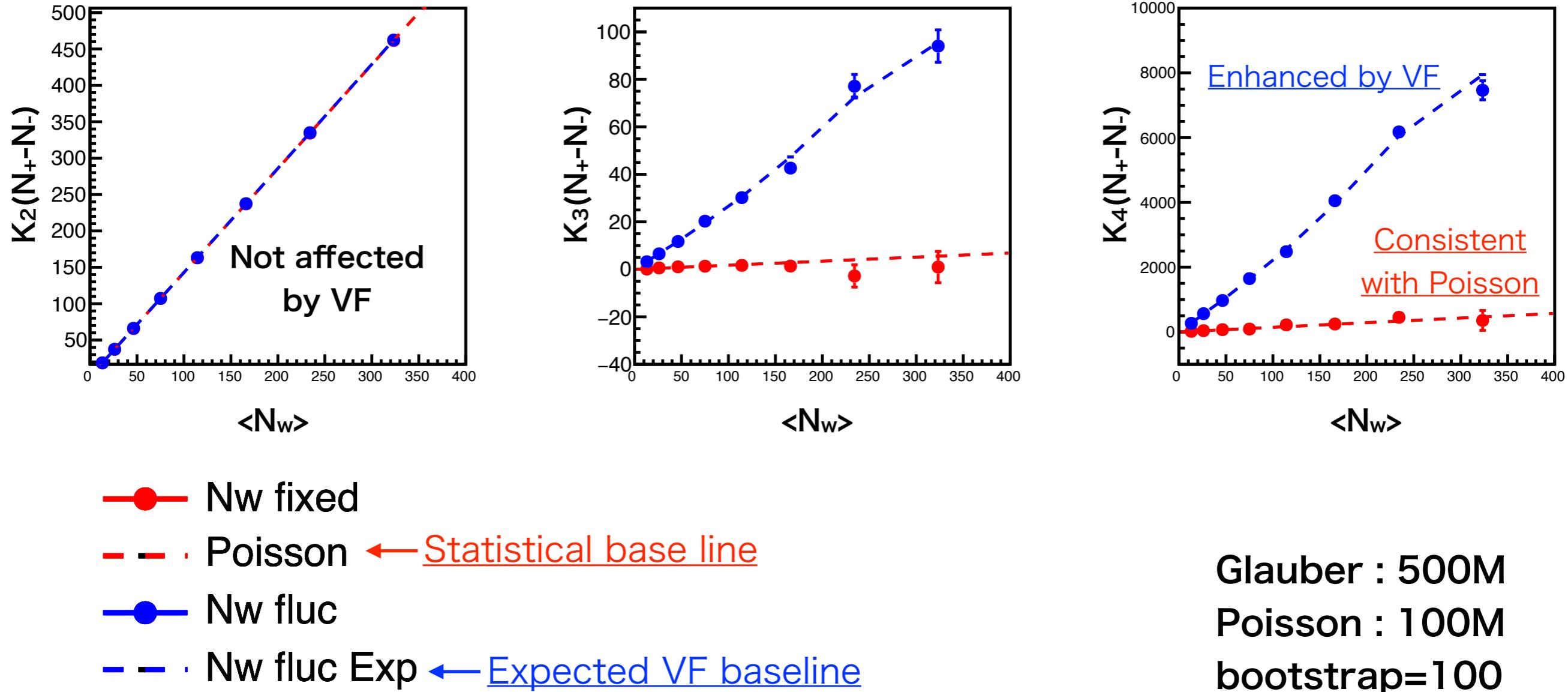
- **Nw fixed** and **VFC** results should be same.
- Results from CBWC are also compared.

VFC : Subtract N_w fluctuation



Toy model VFC check (10%step)

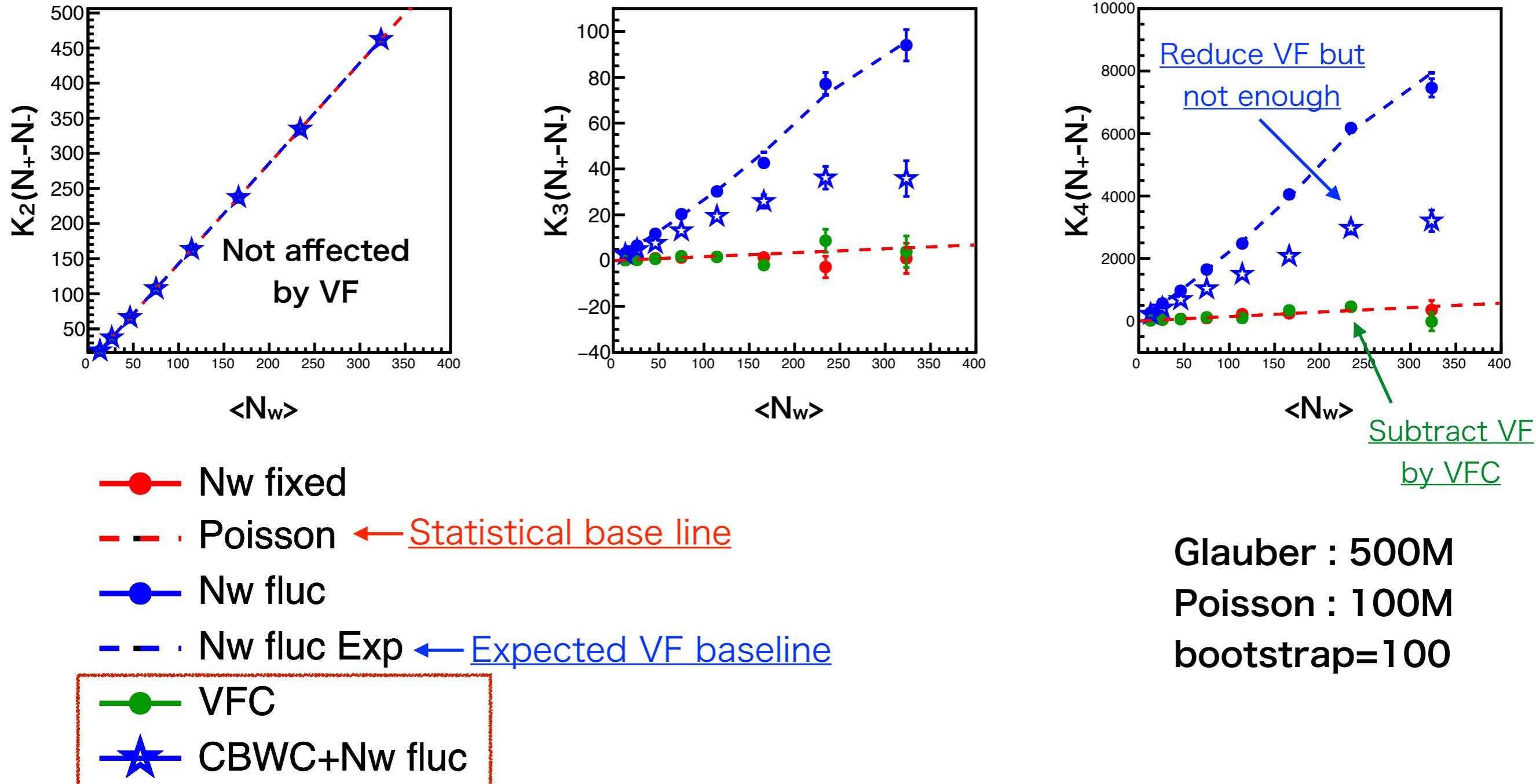
46



- K_2 does not largely affected by participant fluctuation.
- N_w fixed cumulants are consistent with Poisson baseline.
- N_w fluc cumulants are enhanced by VF.

Toy model VFC check (10%step)

47



- VFC works well up to 4-th order cumulants.
- CBWC can reduce VF but not enough.

UrQMD model

- VFC will be done by using UrQMD model in which IPP model might be broken.
- In this model, we can directly know N_w without Glauber fit.

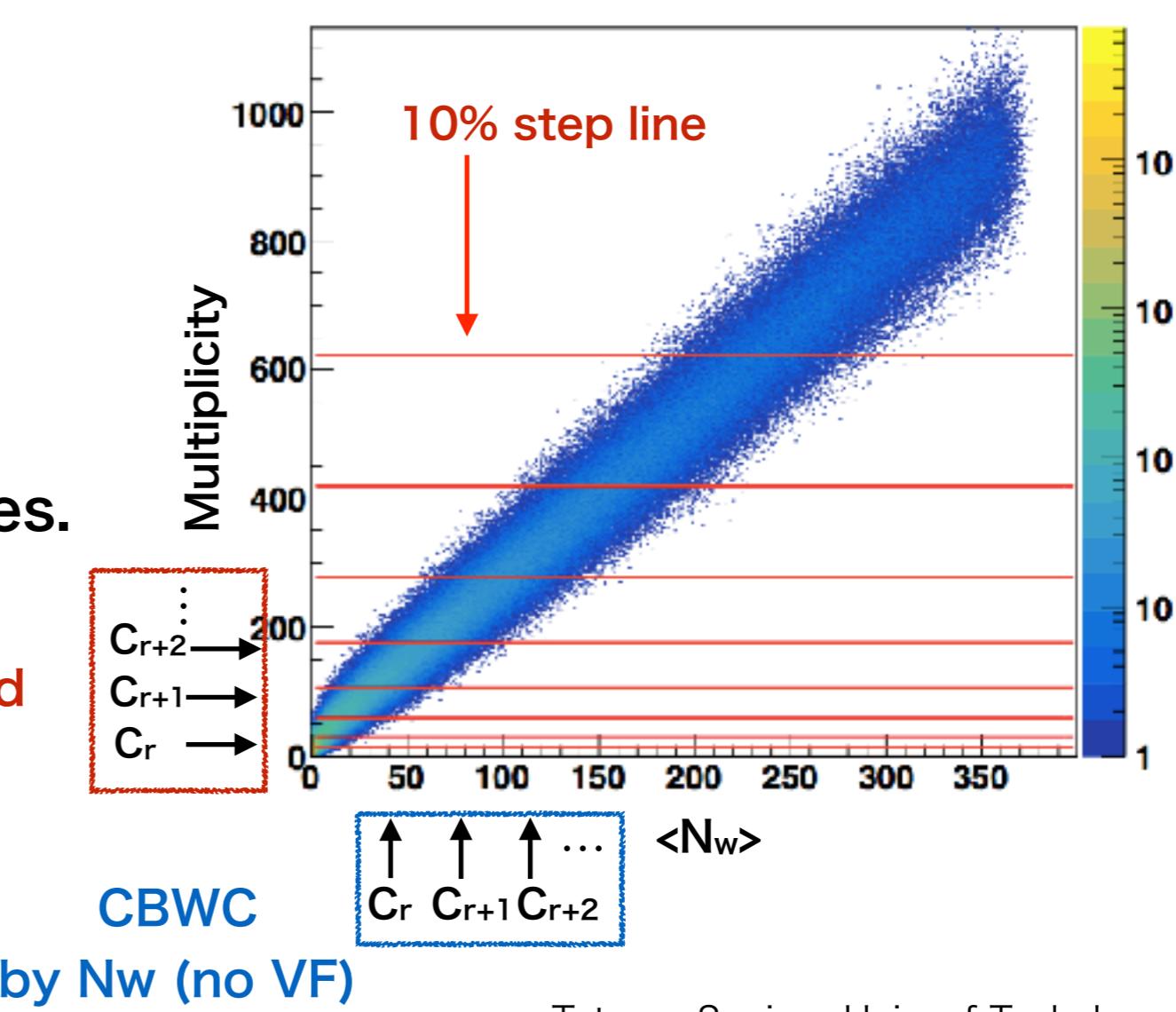
How to know “true” cumulants?

- In order to compare to CBWC and VFC results, we want to know “true” (no VF) cumulants.
- “True” cumulants can be calculated by CBWC using N_w instead of multiplicities.

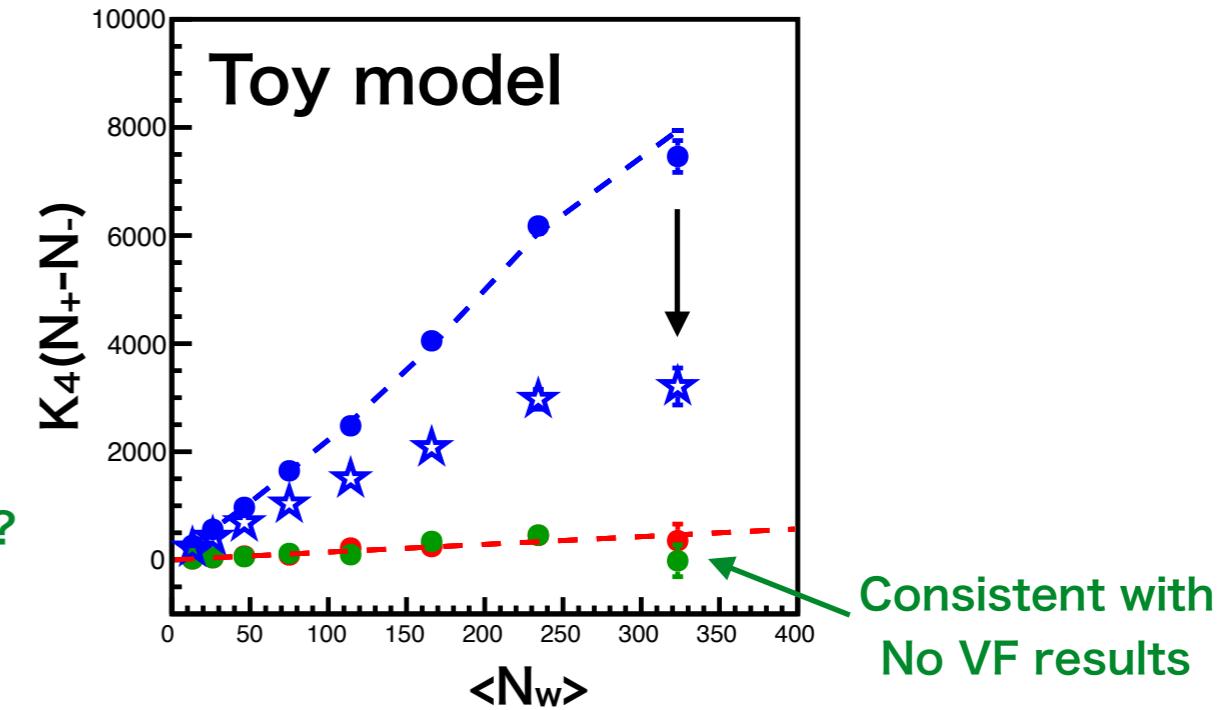
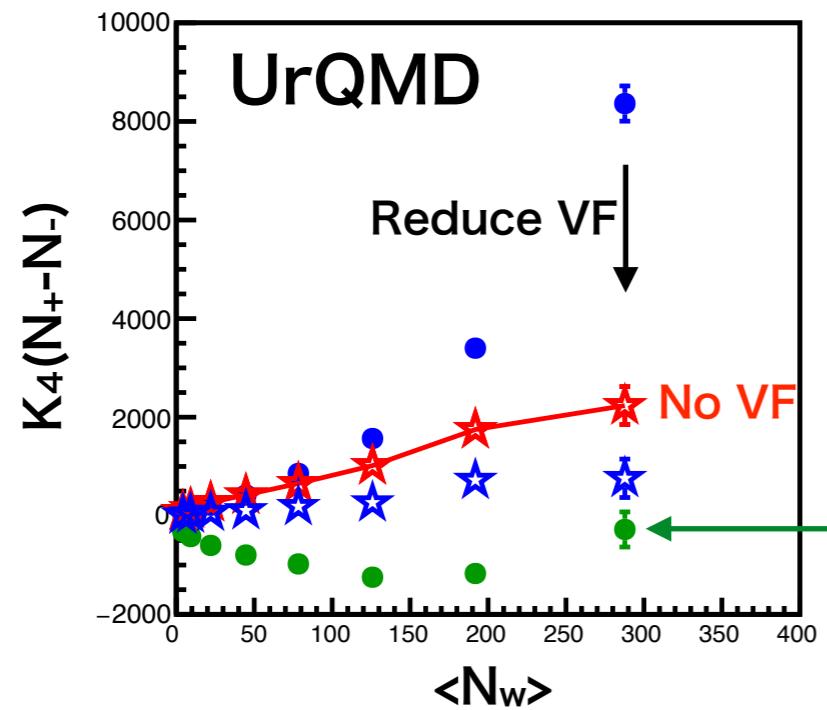
$$C_n = \sum_r w_r C_{(n,r)}$$

Cumulants in
r-th N_w bins

Standatd
CBWC



UrQMD results (Net-charge)



—●— Raw

: No correction

—●— VFC

: VFC was applied

—★— CBWC-Nw (no VF)

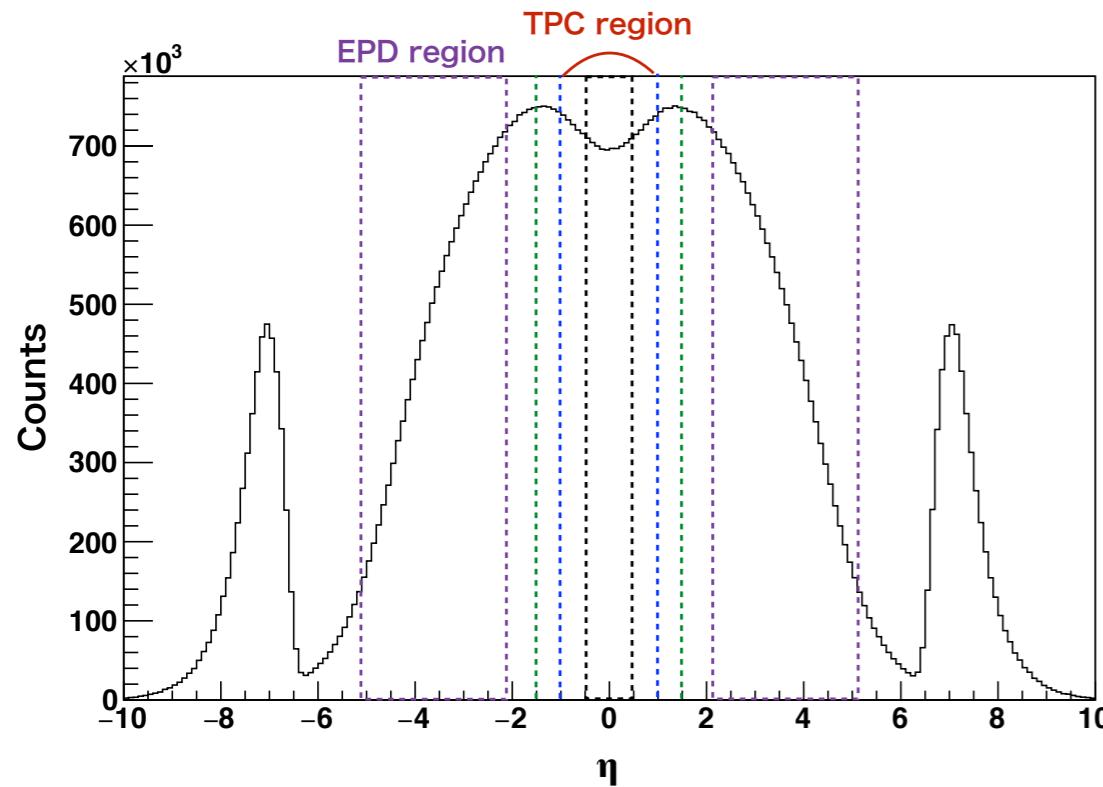
: CBWC by N_w (true)

—★— CBWC

: CBWC by multiplicity (experimental CBWC)

- CBWC-Nw (no VF, true) cumulants are larger than CBWC results even though CBWC results include VF.
 - IPP assumption is expected to be broken in UrQMD.
- CBWC results are larger than VFC results and close to CBWC-Nw cumulants.
 - CBWC is better than VFC?

Correlation in UrQMD



① $|\eta| < 0.5$ (Refmult)

② $0.5 < |\eta| < 1$ (Refmult2)

- experimentally used for net-c analysis

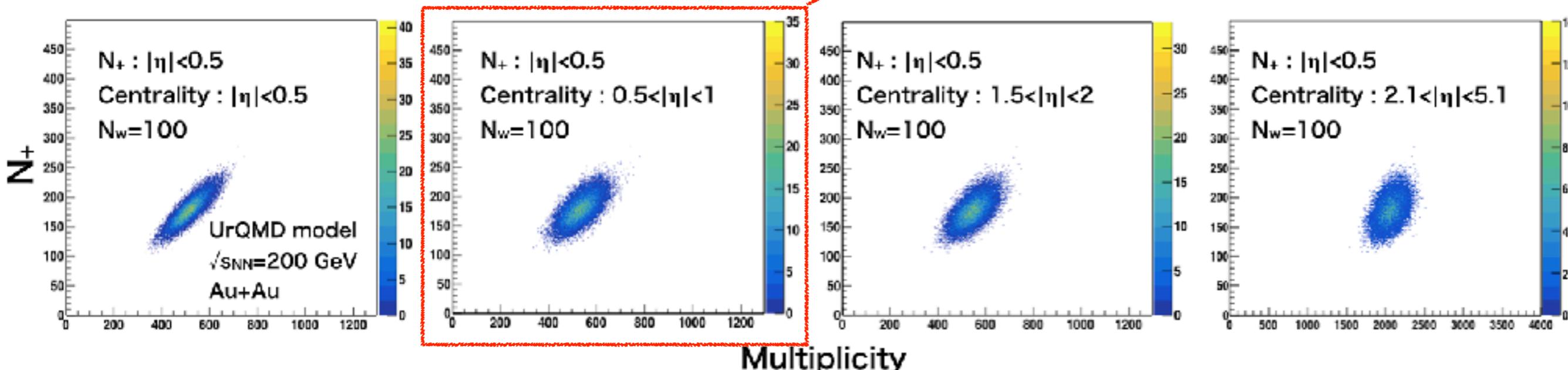
③ $1 < |\eta| < 1.5$

④ $2.1 < |\eta| < 5.1$ (EPD region)

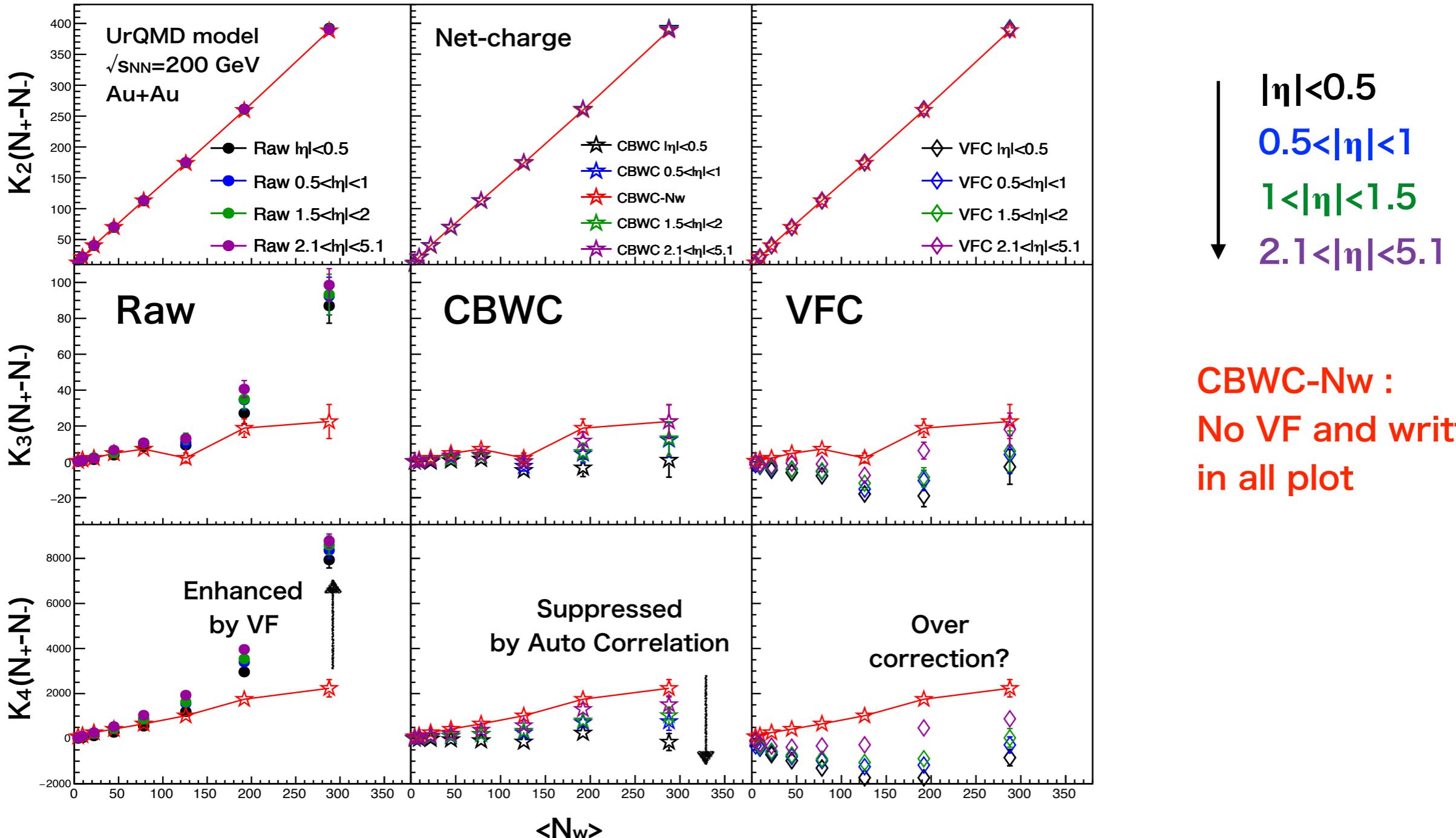
- centrality determination from BES II

- In Toy model, particles used for centrality determination and particles used for the net-charge calculation are produced independently.
- On the other hand, cumulants in UrQMD may include auto-correlation effect.

Correlation remain even though different eta regions were used for centrality determination and particles net-charge calculations

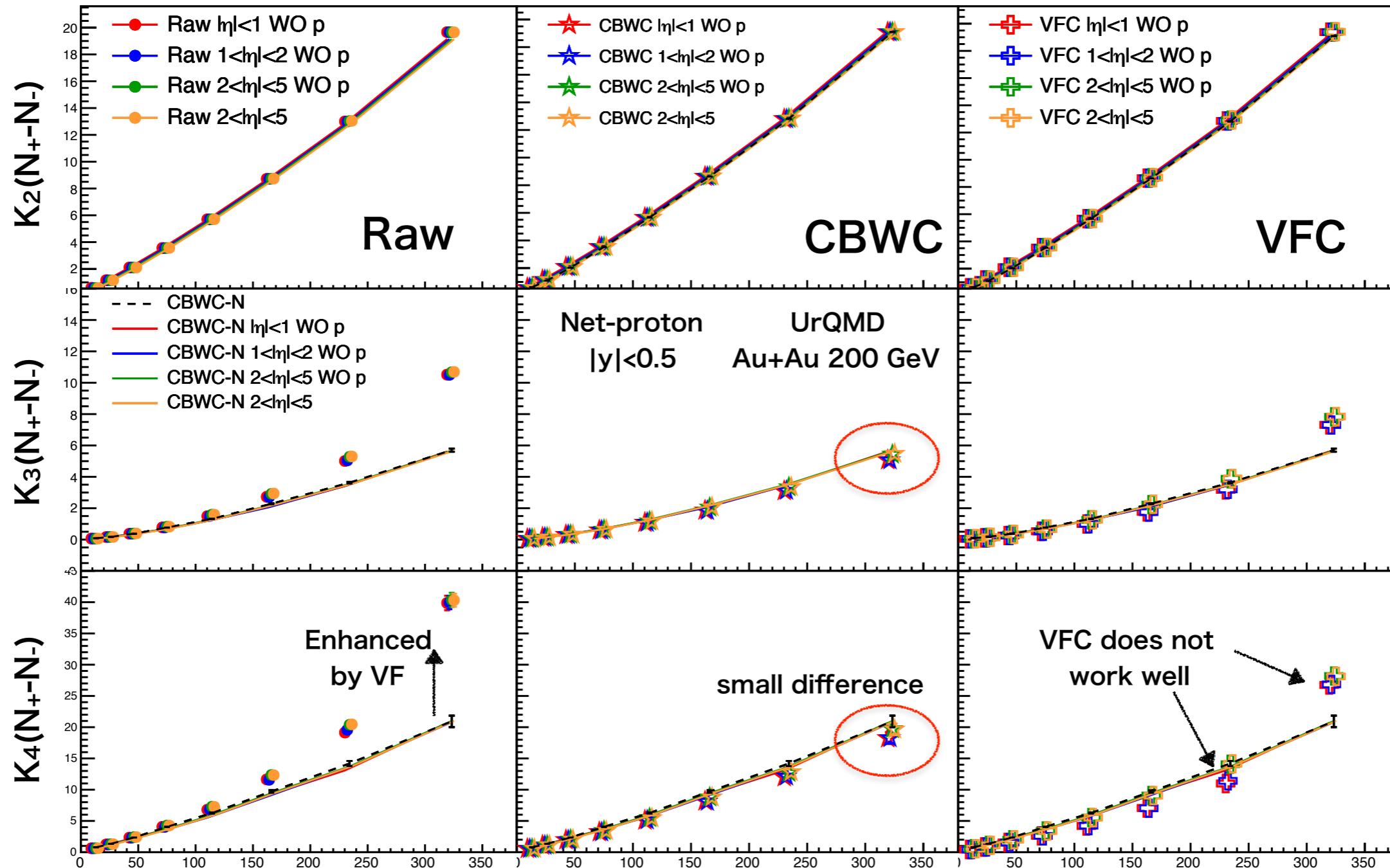


UrQMD with different centrality determination (Net-charge)⁵¹



- Suppressions are observed because of the correlation
 - Largest in $|\eta| < 0.5$ and smallest in EPD region ($2.1 < |\eta| < 5.1$)
- VF are observed in raw results but not observed in CBWC results
 - NOT consistent with Toy model. IPP is broken?
- If CBWC-Nw results are “true”, VFC results are over correction

UrQMD with different centrality determination (Net-proton)⁵²



- Similar studies have been done by net-proton case
- Similar behaviour has been observed as net-charge case but magnitude is smaller than net-charge results.

Summary

C₆ net-charge analysis

- Net-charge cumulants up to 6th-order at $\sqrt{s_{NN}}=200 \text{ GeV}$ and 54 GeV in Au+Au collisions were reported.
- C₆/C₂ were compared to Poisson, NBD, UrQMD and net-proton results.
- Deviation are observed between 200 GeV and 54 GeV in 0-40% centrality at net-proton and 40-50% at net-charge.
 - These results are not conflict with the theoretical prediction of cross-over transition but statistical errors are large.
 - Results of the other centralities are consistent with statistical baseline within the statistical error.
- C₃/C₂ and C₄/C₂ of net-charge at $\sqrt{s_{NN}}=54 \text{ GeV}$ are newly measured in addition to BES energies and results were not conflict with published BES-I results.

Summary

Δη dependence analysis

- **Δη dependence of net-charge cumulants and D-measure** were measured in Au+Au collisions at $\sqrt{s_{NN}}=7.7, 11.5, 14.5, 19.6, 27, 39, 62.4$ and 200 GeV.
- D-measure **decrease with Δη** and this trend is **stronger with collision energies**.
 - These results **do not conflict with ALICE** and the trends of $\Delta\eta$ dependence **might represent the time evolution of phase transition**.
- $C_4/\langle N_{ch} \rangle$ and $C_3/\langle N_{ch} \rangle$ are observed to increase with $\Delta\eta$ in most of centralities which may be the hint of the signal from phase transition.

Summary

VFC

- The validity of the VFC on higher-order net-charge and net-proton cumulants have been studied by using toy model and UrQMD model in Au+Au collisions at $\sqrt{s_{NN}} = 200$ GeV.
- **Even though CBWC has applied, we can't completely remove participant fluctuation under the IPP models.**
- However, **VFC does not works well in UrQMD model so IPP model is expected to be broken** in UrQMD.
- The values of cumulants are also depend on how determine centrality because an correlation between net-quantities and the multiplicity which is used for centrality determination is different.
- EPD is installed from BES-II, and the cumulants **may become larger** than current results when we determine the centrality by EPD instead of TPC.

Ouklook

- BES-II will start from 2019
 - High statistics at lower energy region
 - Access more forward rapidity with new detector (iTPC, eTOF and EPD).
- Sixth-order cumulants can be measured at lower energy region and can be compare to the results at 200 GeV and 54GeV.
- $\Delta\eta$ dependence can be measured from 0 to 1.5 by iTPC and eTOF ($0 < \Delta\eta < 1$ in current analysis).
- It is important to study whether we should apply VFC to experimental results or not for future analysis.

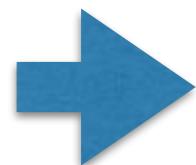
back up

Statistical errors of cumulants

Generally...

- ✓ Statistical error of **net-charge cumulant** is larger than **net-proton cumulants**
- ✓ Statistical error of cumulant **in central** is larger than cumulant in **peripheral**

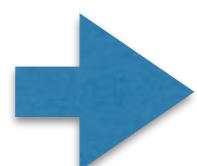
Why?



Cumulants depend on the width of the distribution (σ).

For example

$$\text{Mean value : } C_1 \pm \sigma / \sqrt{N} \longrightarrow \begin{matrix} \text{Standard} \\ \text{Error} \end{matrix}$$



- Statistical error of C_1 is proportional to σ .
- Statistical error of nth-order cumulant is written by

$$\text{Err}(C_n) \sim \sigma^n / \sqrt{N}$$

Rough estimation of statistical errors

Efficiency Uncorrected

$$\text{Err} \left(\frac{C_3}{C_2} \right) \propto \frac{\sigma}{\sqrt{n}}$$

$$\text{Err} \left(\frac{C_4}{C_2} \right) \propto \frac{\sigma^2}{\sqrt{n}}$$

$$\text{Err} \left(\frac{C_6}{C_2} \right) \propto \frac{\sigma^4}{\sqrt{n}}$$

Efficiency Corrected

$$\text{Err} \left(\frac{C_3}{C_2} \right) \propto \frac{\sigma}{\epsilon \sqrt{n}}$$

$$\text{Err} \left(\frac{C_4}{C_2} \right) \propto \frac{\sigma^2}{\epsilon^2 \sqrt{n}}$$

$$\text{Err} \left(\frac{C_6}{C_2} \right) \propto \frac{\sigma^4}{\epsilon^4 \sqrt{n}}$$

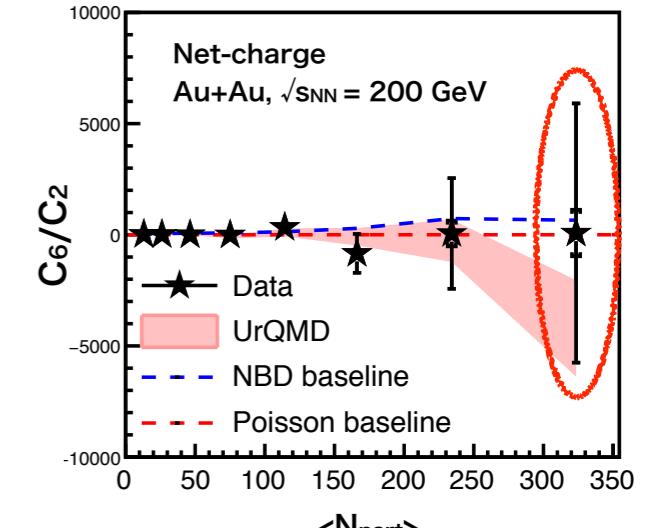
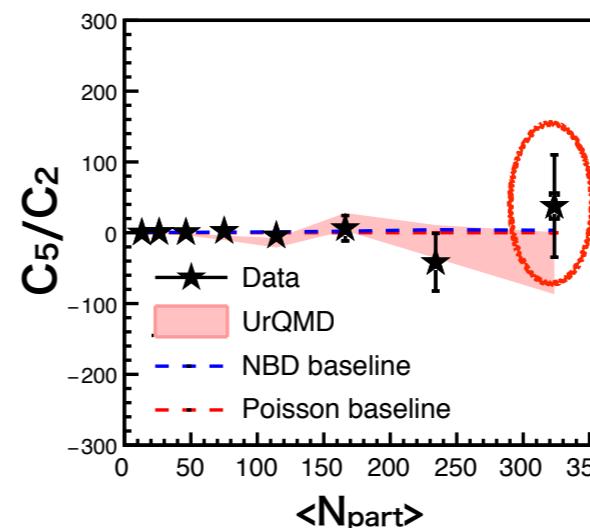
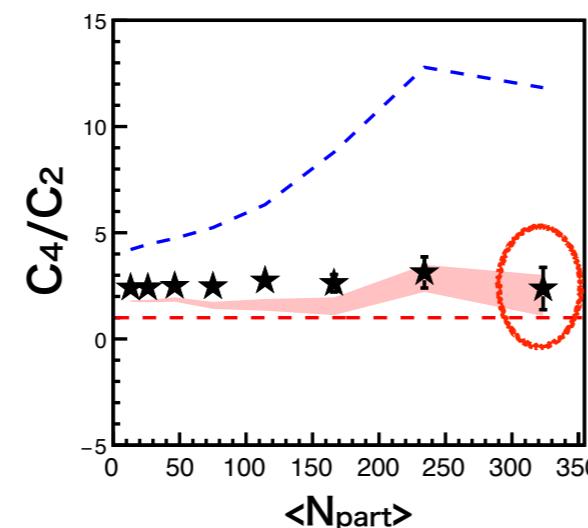
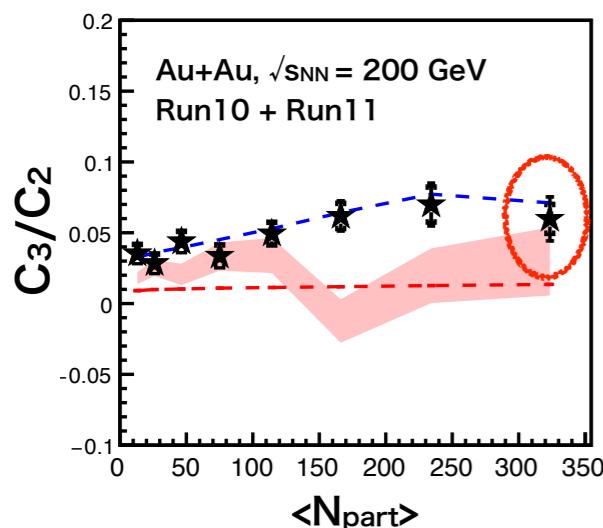
For example: $\sigma^2=350$, $\epsilon=0.4\sim0.5$

(0-10%, 200GeV, Eff corr)



$\sigma / \epsilon = 40 \sim 100$

- Error of C_6/C_2 is more than 1000 time large than that of C_4/C_2 .

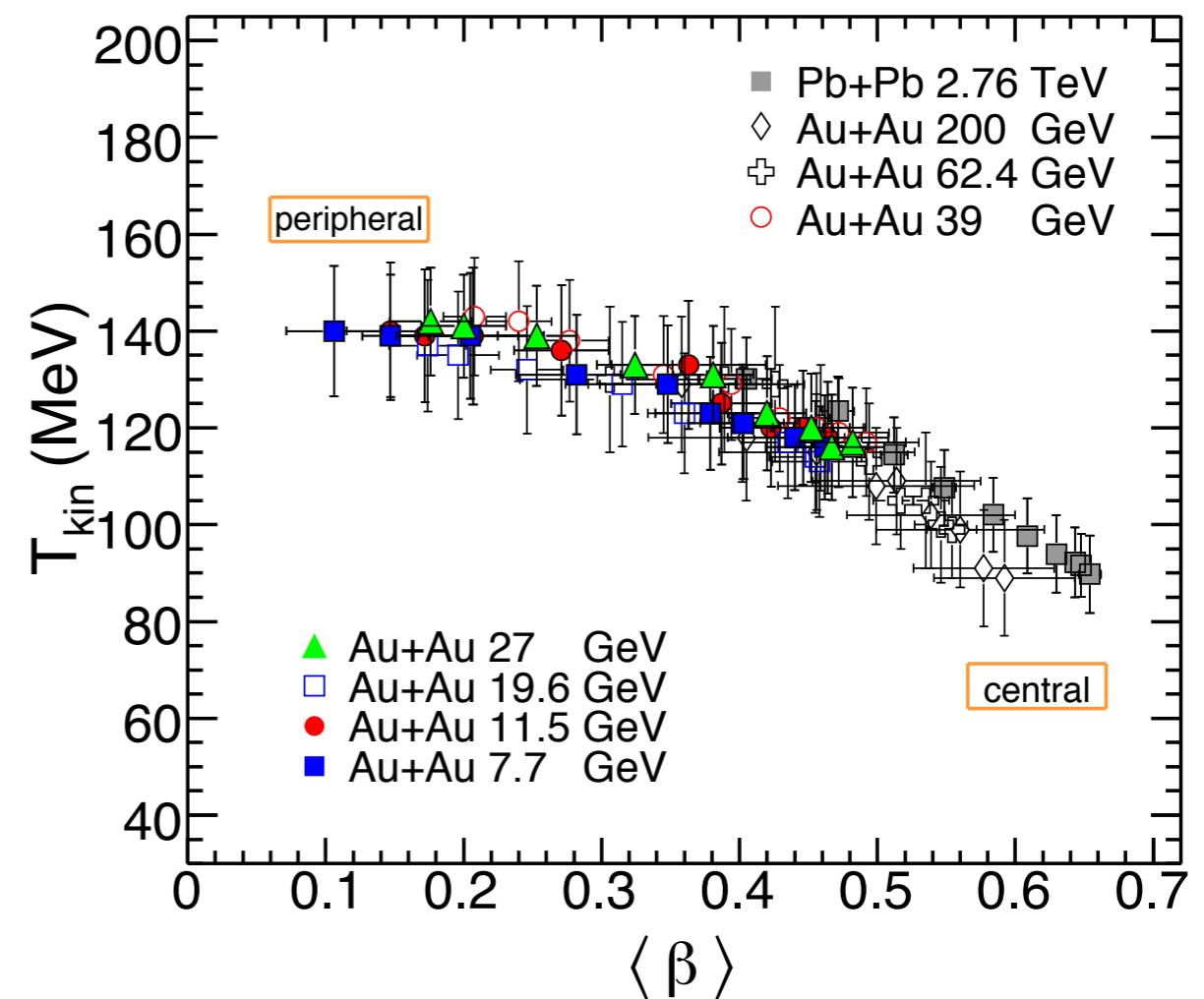
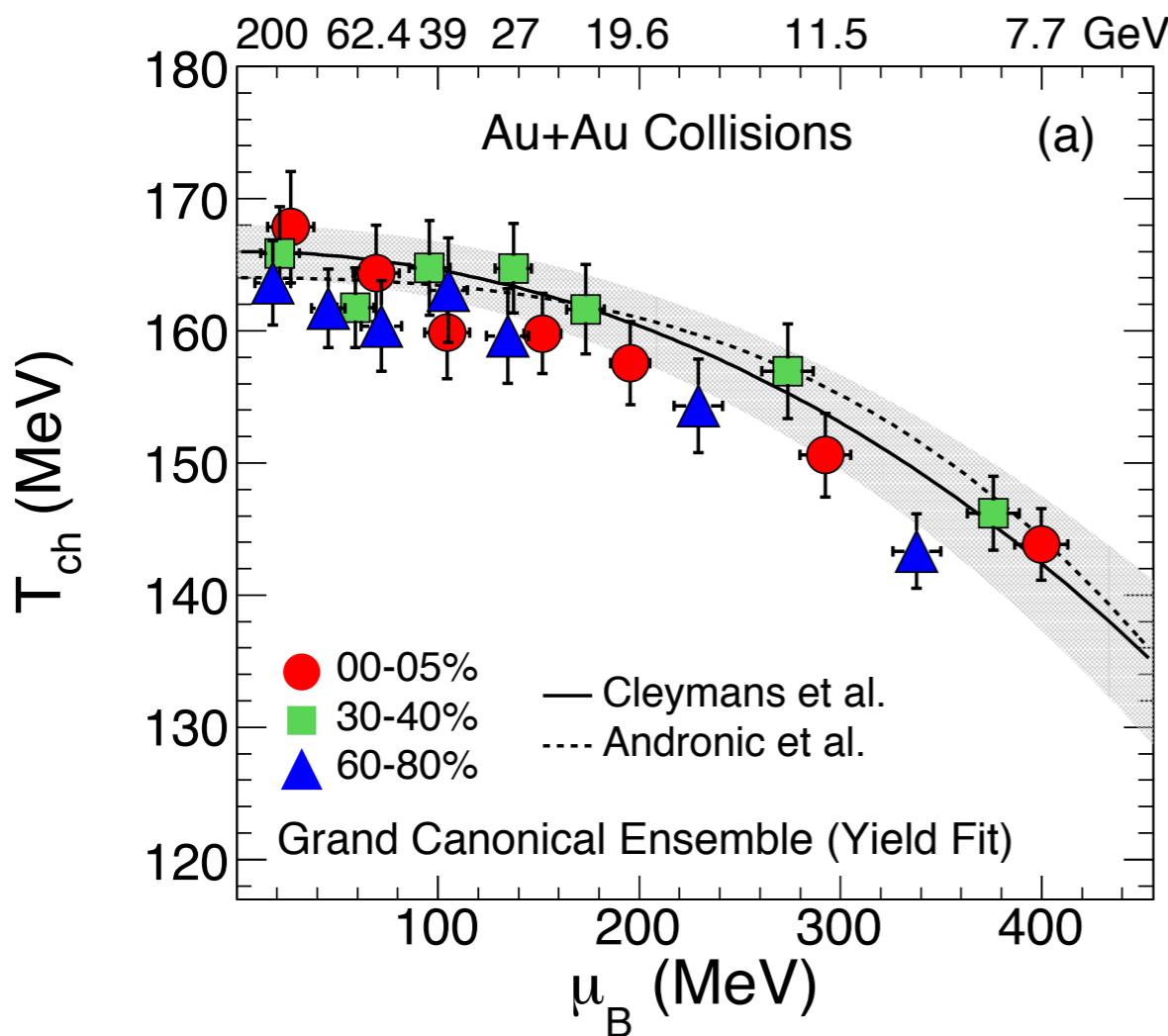


	C3/C2	C4/C2	C5/C2	C6/C2
St.Err	~0.014	~1	~70	~5600

* σ / ϵ

Freeze-out parameters

60



- Chemical and kinetic freeze-out parameters are obtained from particles yield and spectra.

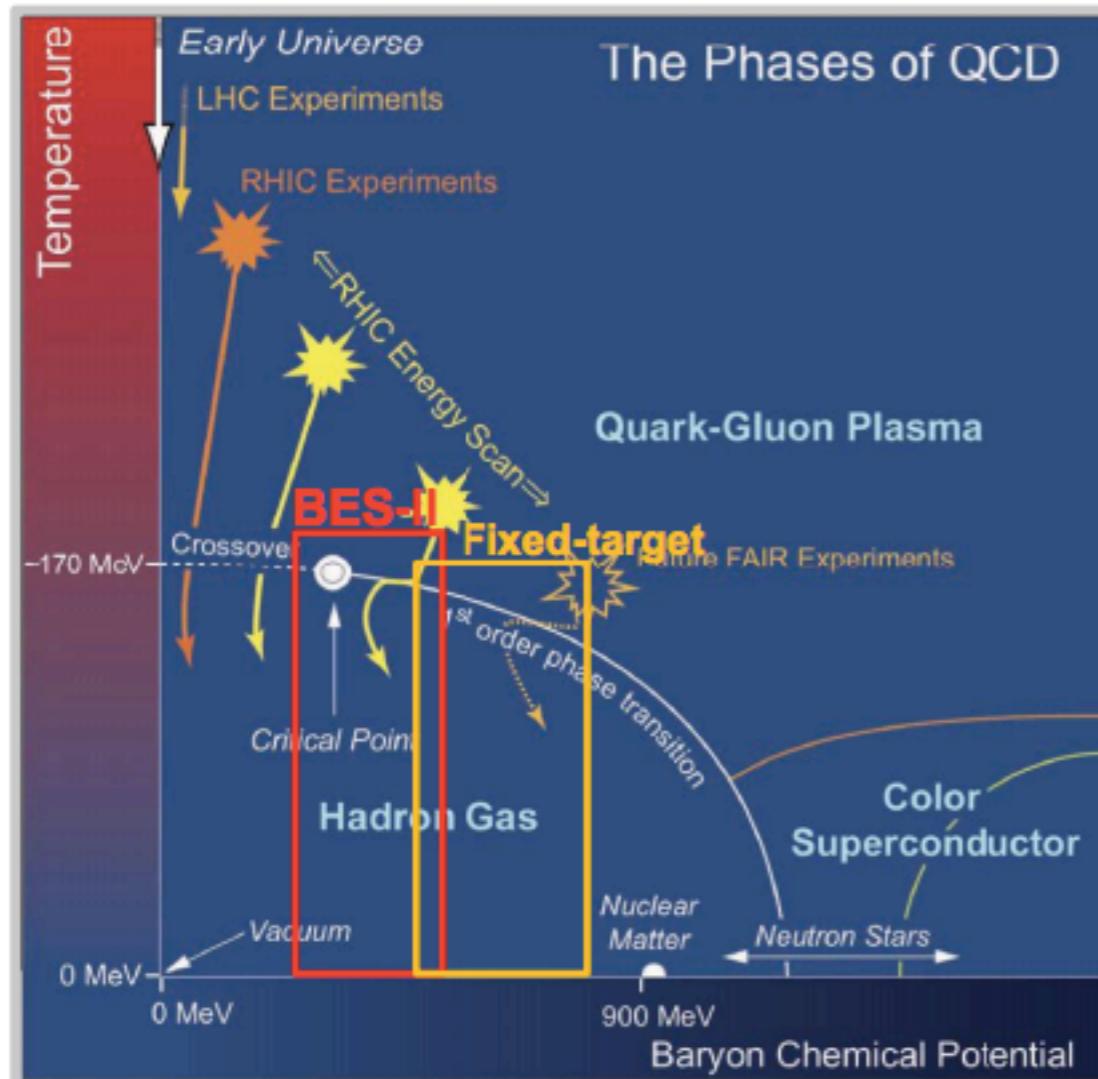
What is the good observables which is sensitive to the phase transition?



Event by event fluctuations are considered as one of the powerful tool to discover the phase diagram.

Beam Energy Scan Phase II (BES-II)

STAR, Q.Yang (talk)

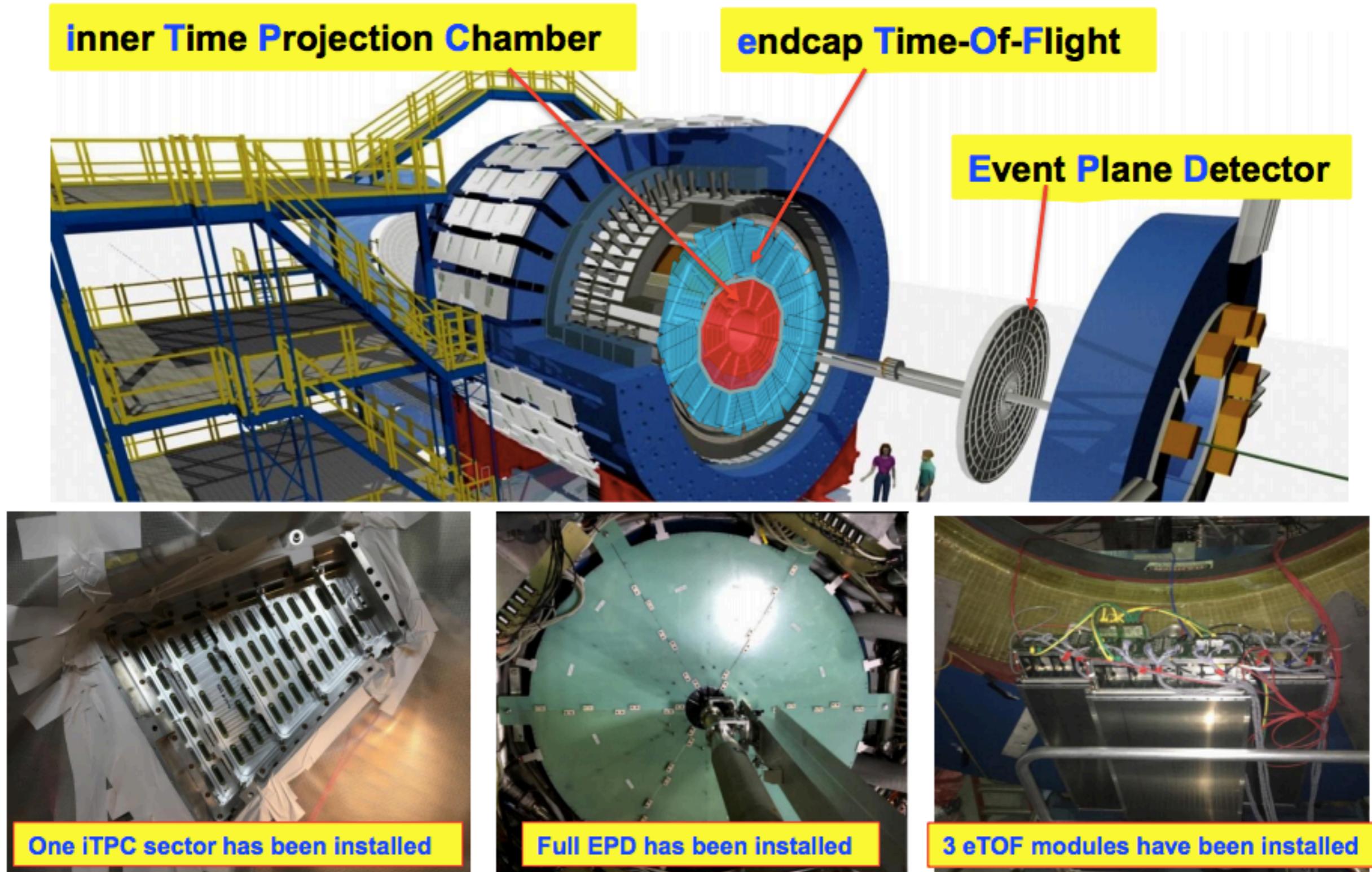


\sqrt{s}_{NN} (GeV)	Proposed Event Goals (M)	BES-I Event (M)
7.7	100	4
9.1	160	N/A
11.5	230	12
14.5	300	20
19.6	400	36
3.0 - 7.7	~100 per energy	N/A

■ Collider mode
■ Fixed-target mode

- RHIC BES II : 10-25 times more statistics and detector upgrade
→ **Dramatically reduce the uncertainties.**
- Precise map the QCD phase diagram $200 < \mu_B < 720$ MeV

Detector upgrades



Q.Yang (STAR collaboration) QM2018

Value of D-measure

63

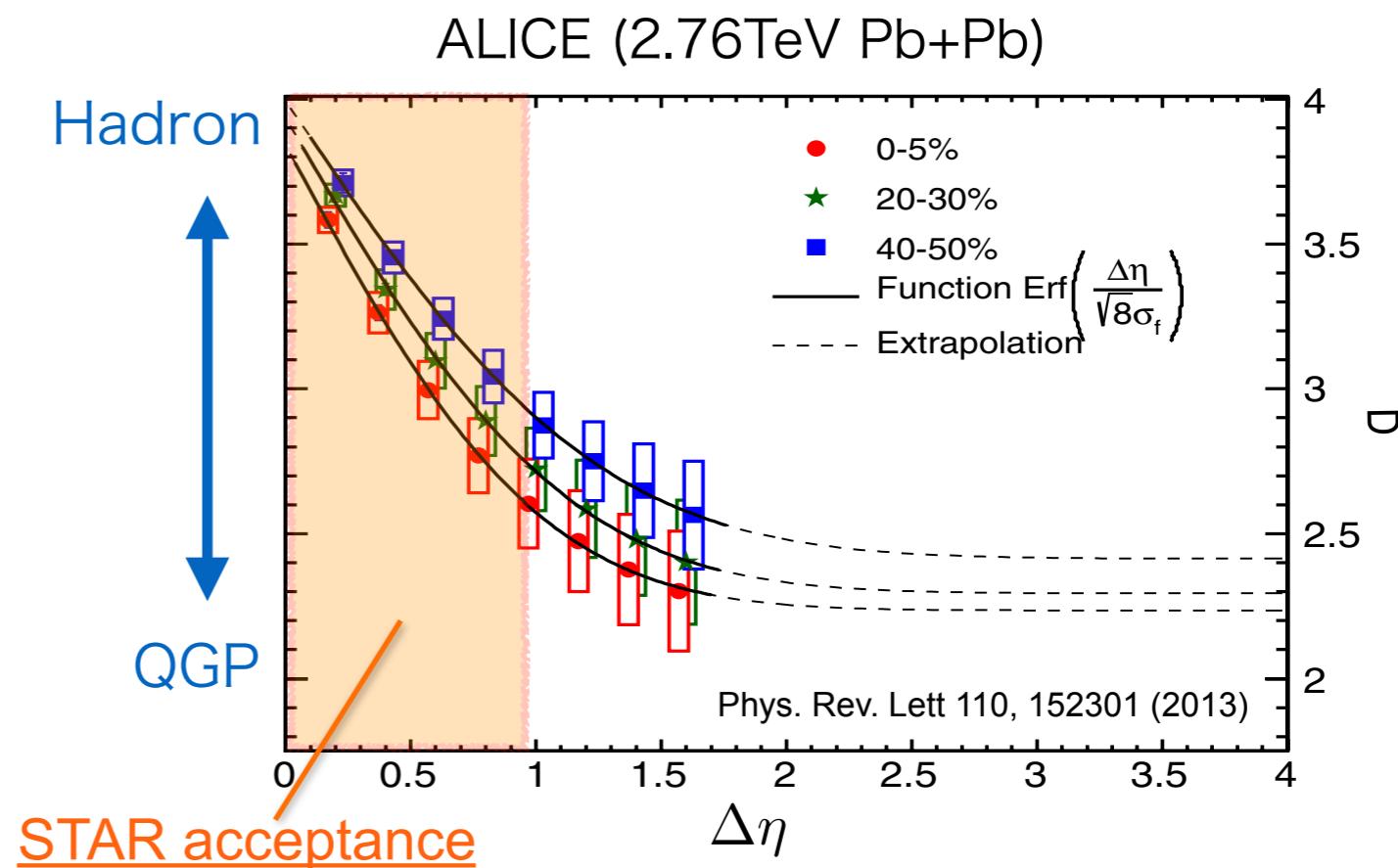
D-measure (D)

$$D = 4 \frac{C_2}{\langle N_{ch} \rangle} \quad N_{ch} = N^+ + N^-$$

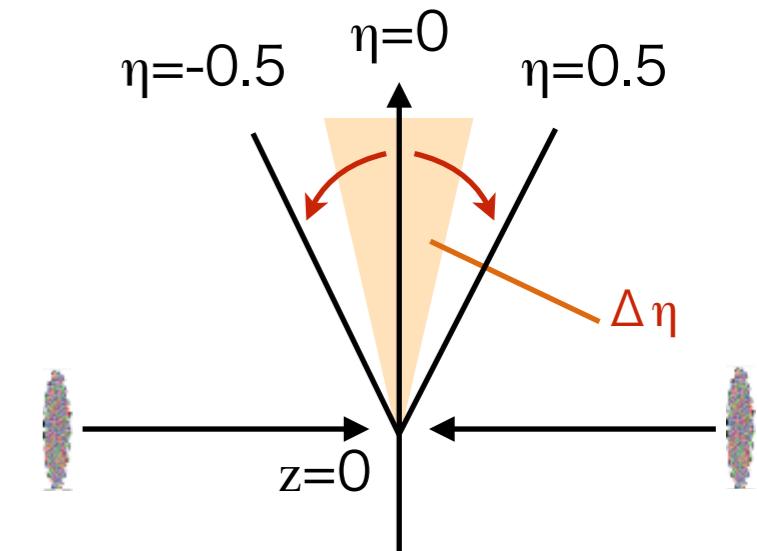
Theoretical predictions

QGP fluctuation : $D = 1-1.5$

Hadron fluctuation : $D = 3-4$



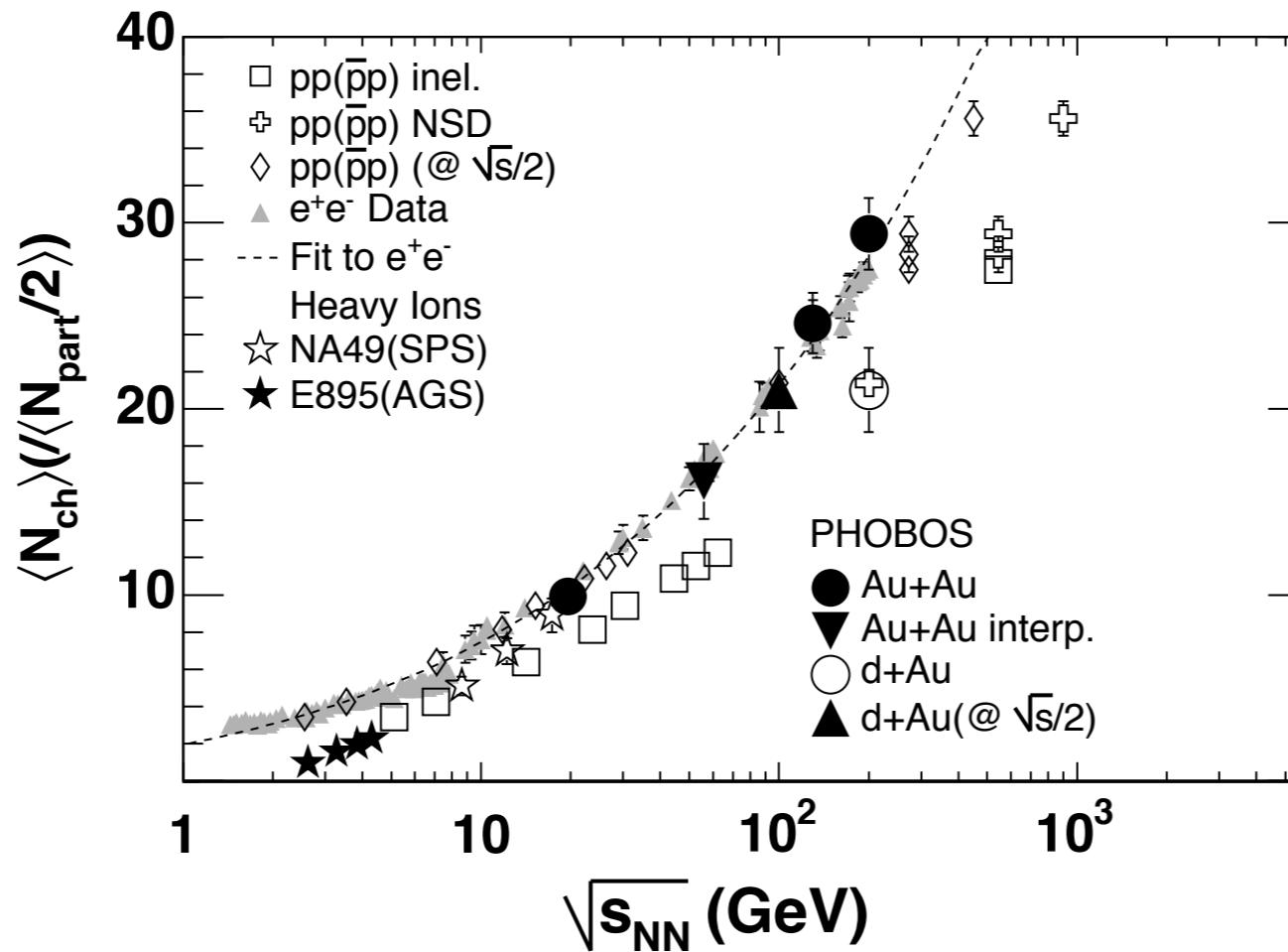
- D-measure is observed to decrease with expanding $\Delta\eta$ in Pb-Pb collisions at 2.76 TeV at ALICE.**
- D-measure also decrease from peripheral to central collisions.



Charge conservation correction on D

64

B. B. Back et al. (PHOBOS Collaboration),
Nuclear Physics A Volume 757, Issues 1–2, 8 August 2005, Pages 28–101

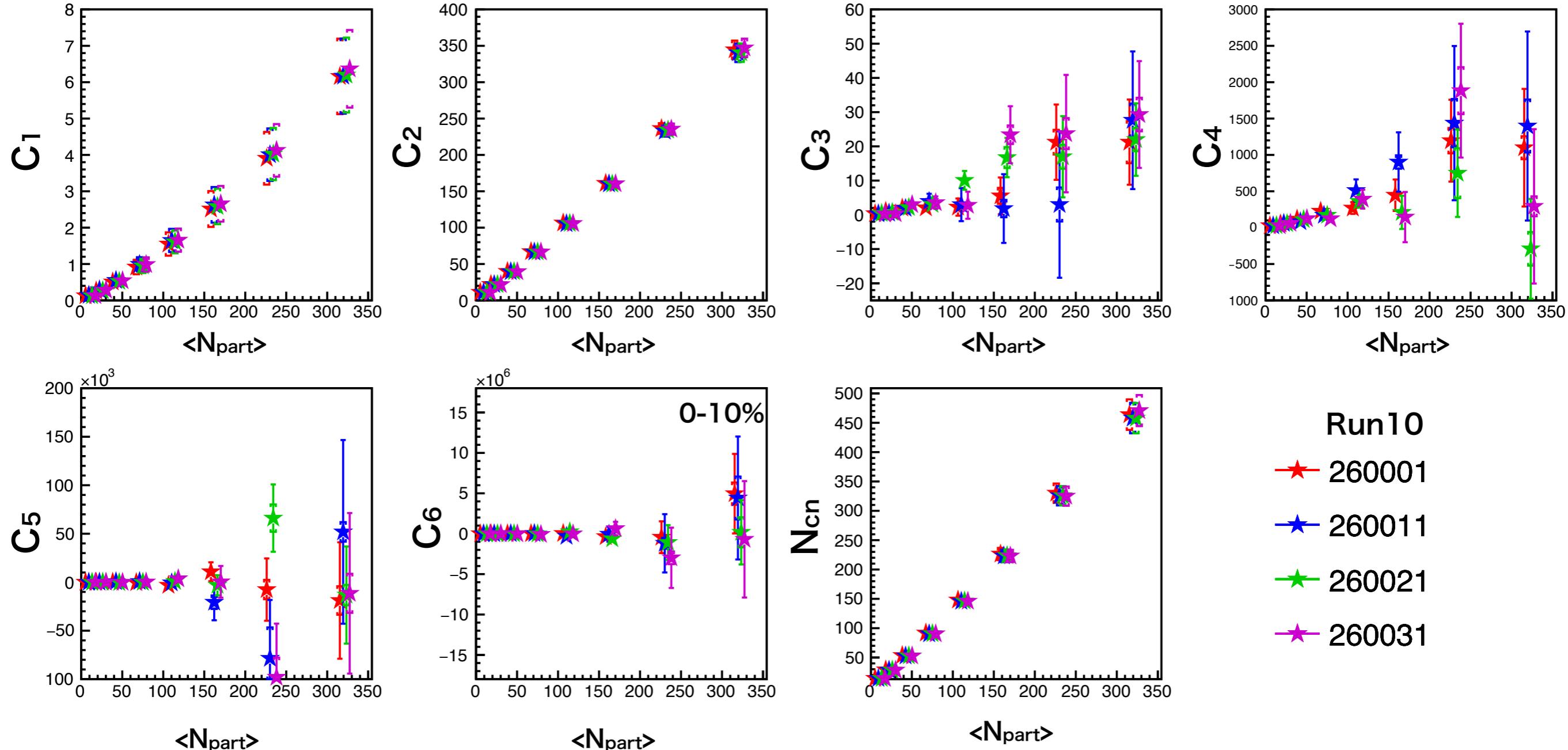


- N_{total} is total multiplicity in full acceptance and estimated from PHOBOS experiment.

back up
(Experimental results)

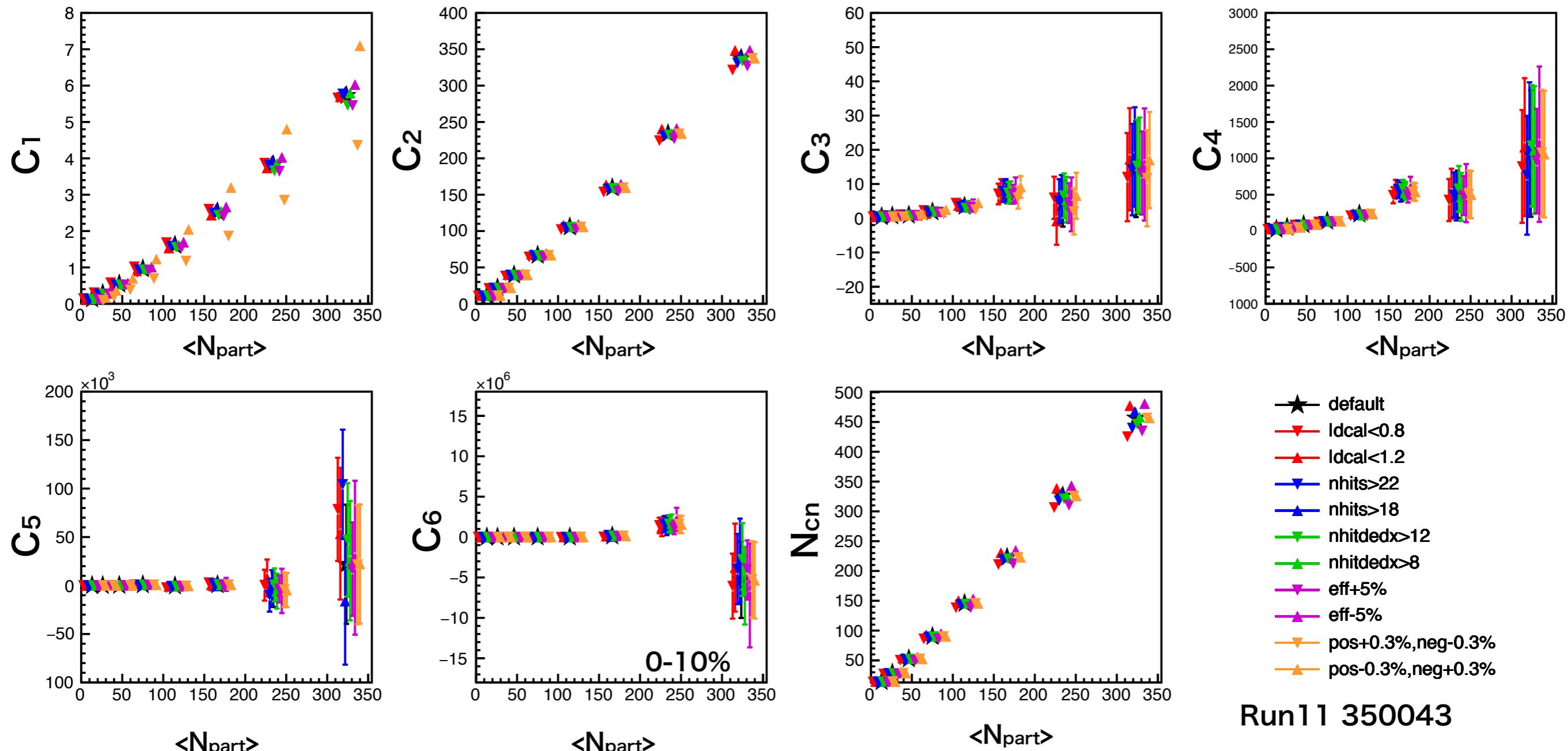
Cumulants (Run10, trigger by trigger)

66



Cumulants (Run11, 350043)

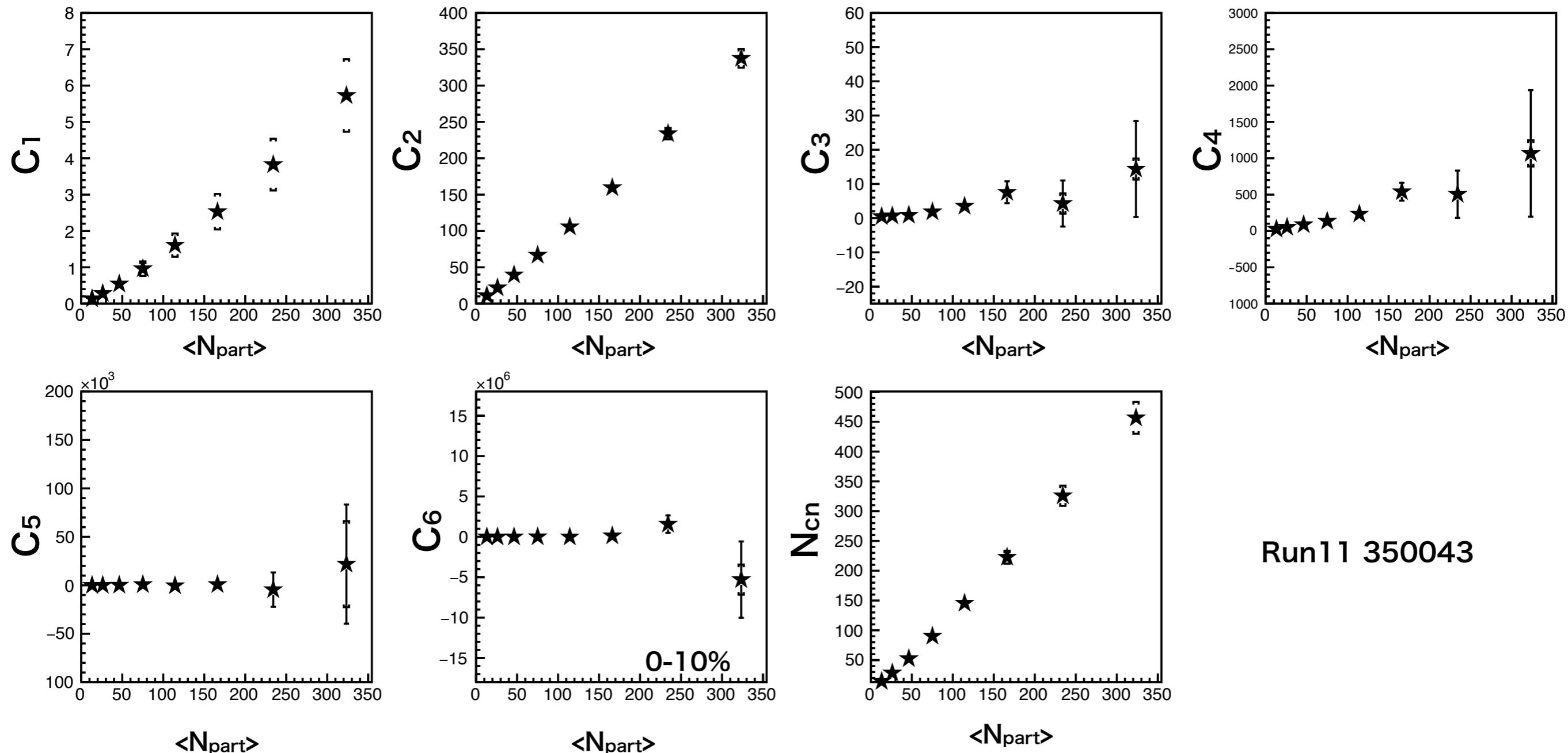
67



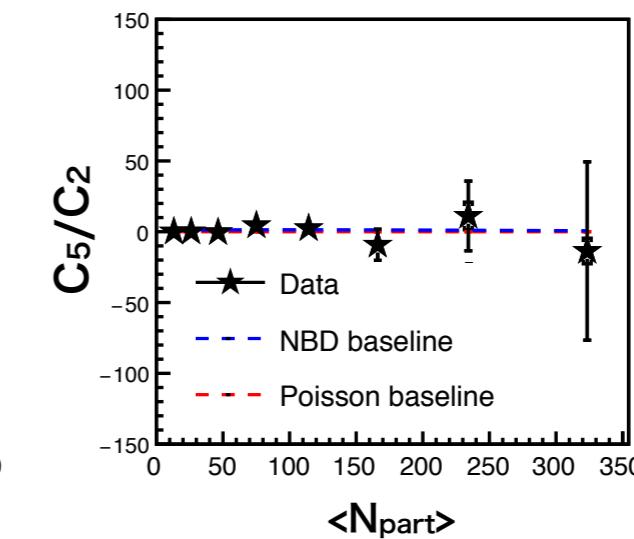
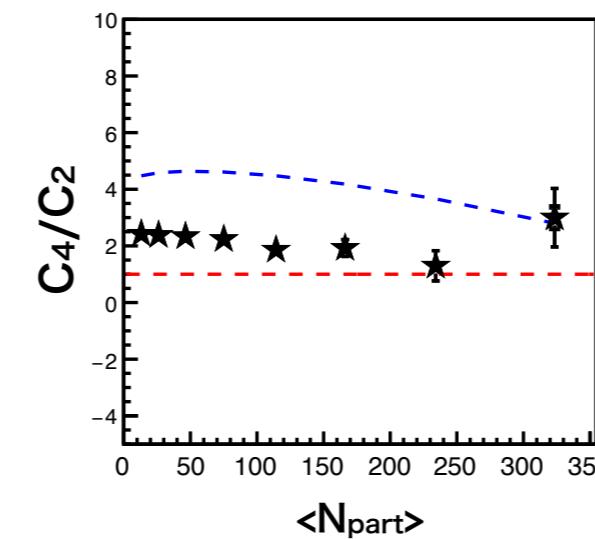
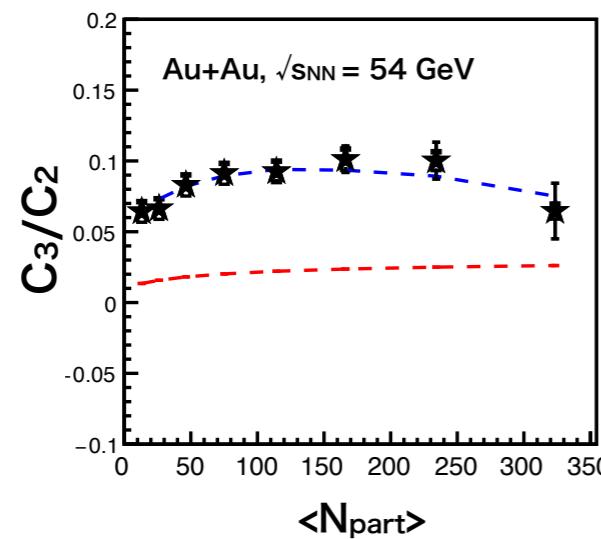
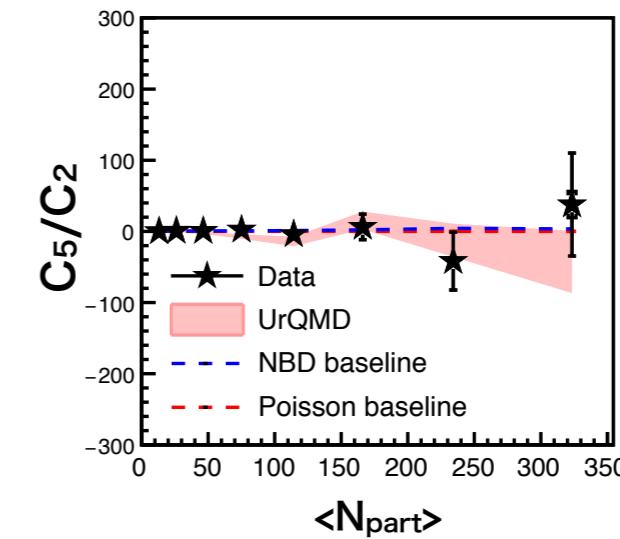
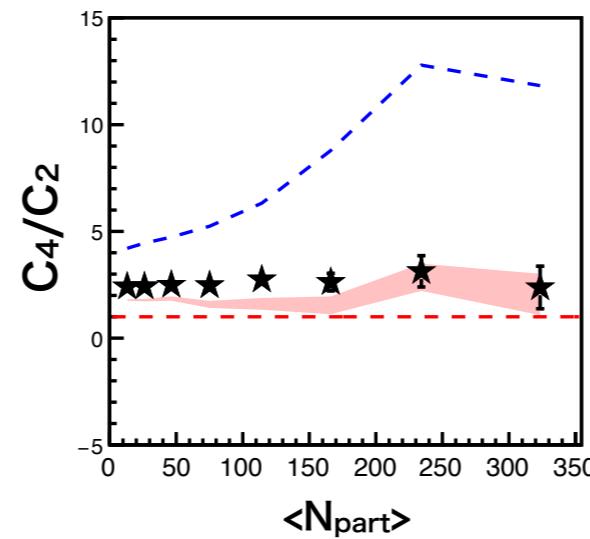
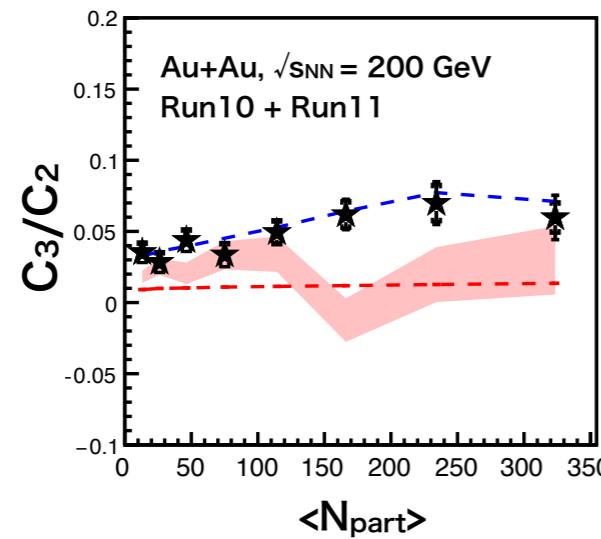
- ($\text{Effpos} * 1.003$, $\text{Effneg} * 0.997$) and ($\text{Effpos} * 0.997$, $\text{Effneg} * 1.003$) were calculated to estimate systematic uncertainty of the deviation between positively and negatively charged particle efficiency.

Cumulants (Run11, 350043)

68

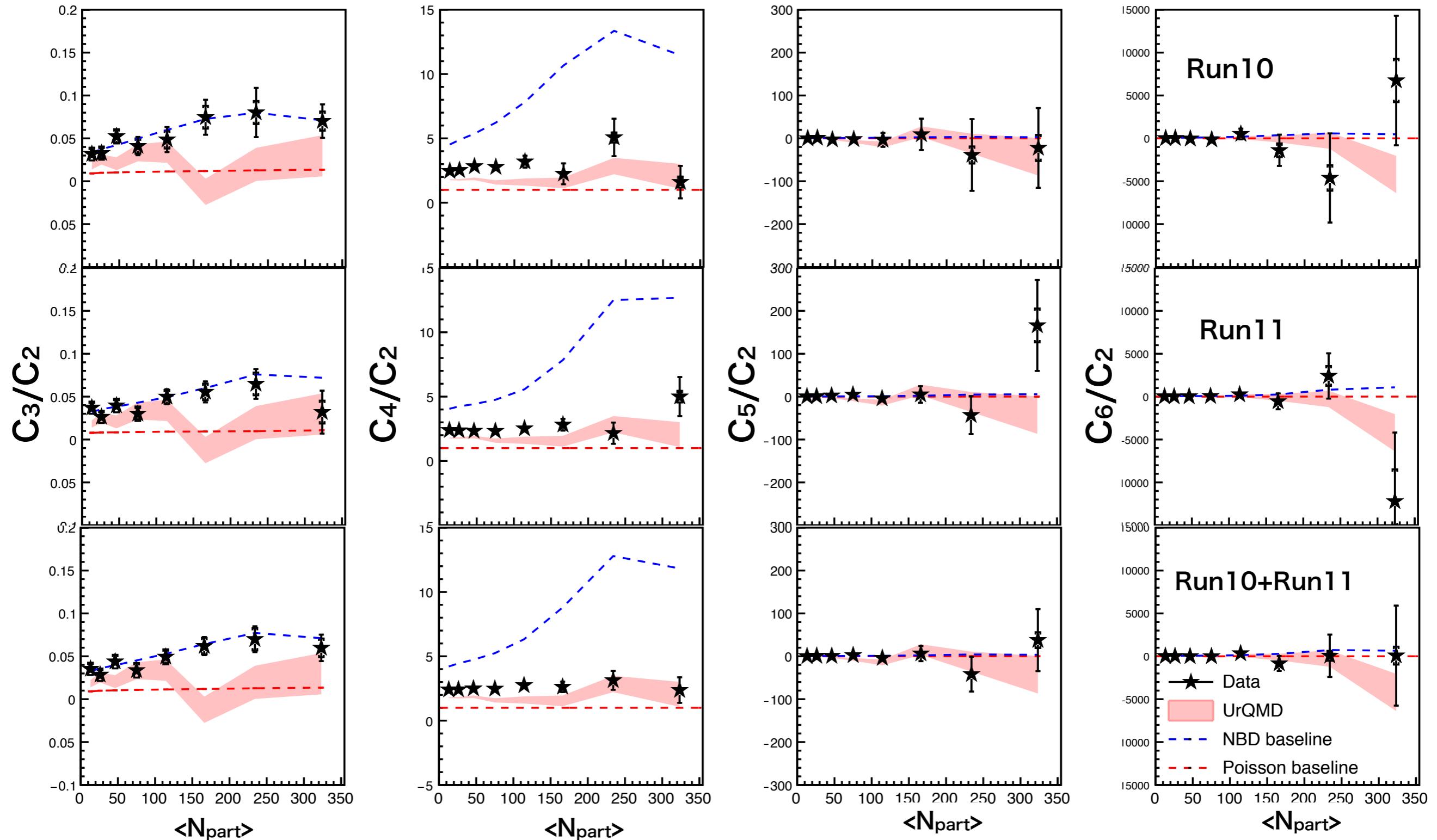


- ($\text{Effpos}^*1.003, \text{Effneg}^*0.997$) and ($\text{Effpos}^*0.997, \text{Effneg}^*1.003$) were calculated to estimate systematic uncertainty of the deviation between positively and negatively charged particle efficiency.



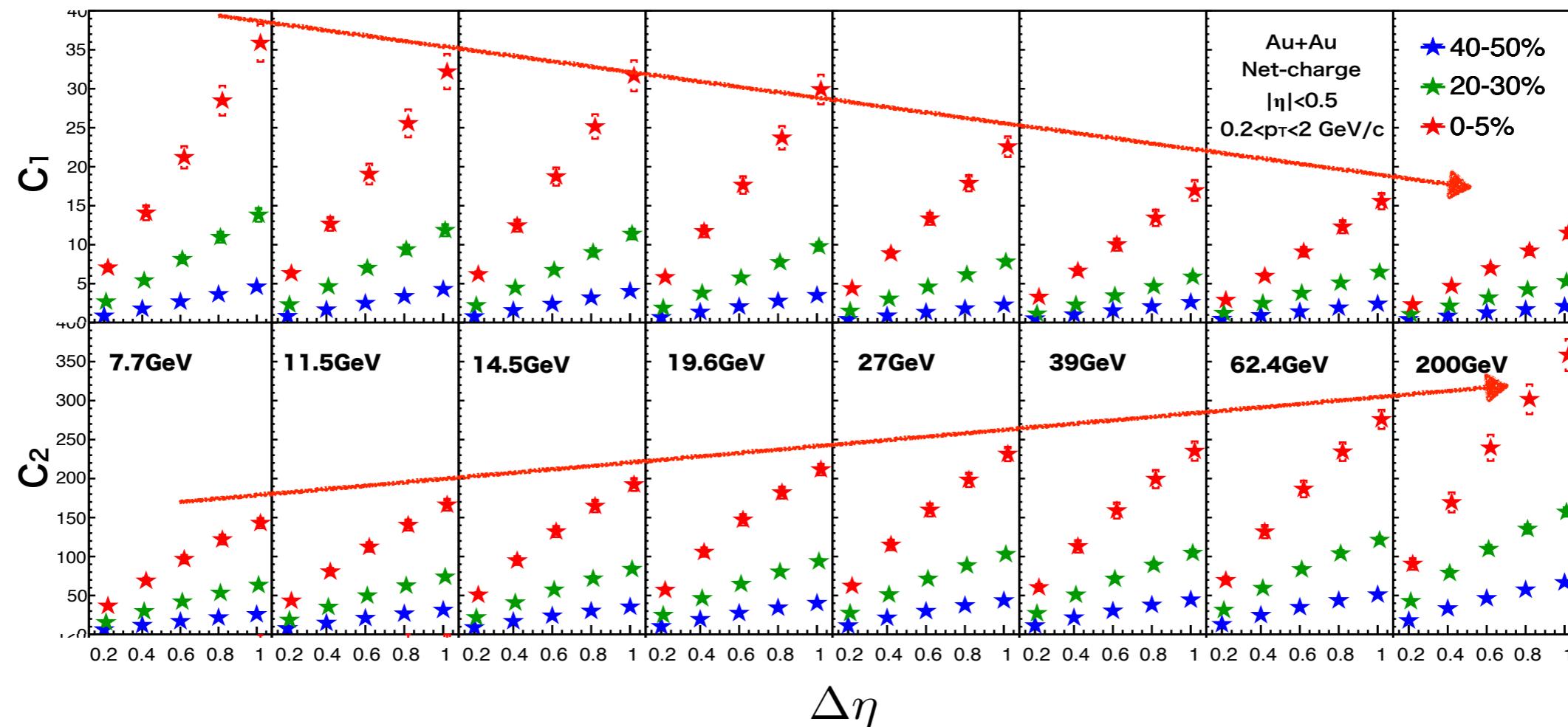
Cumulant ratio (Run10+Run11)

70



$\Delta\eta$ dependence of C_1 and C_2

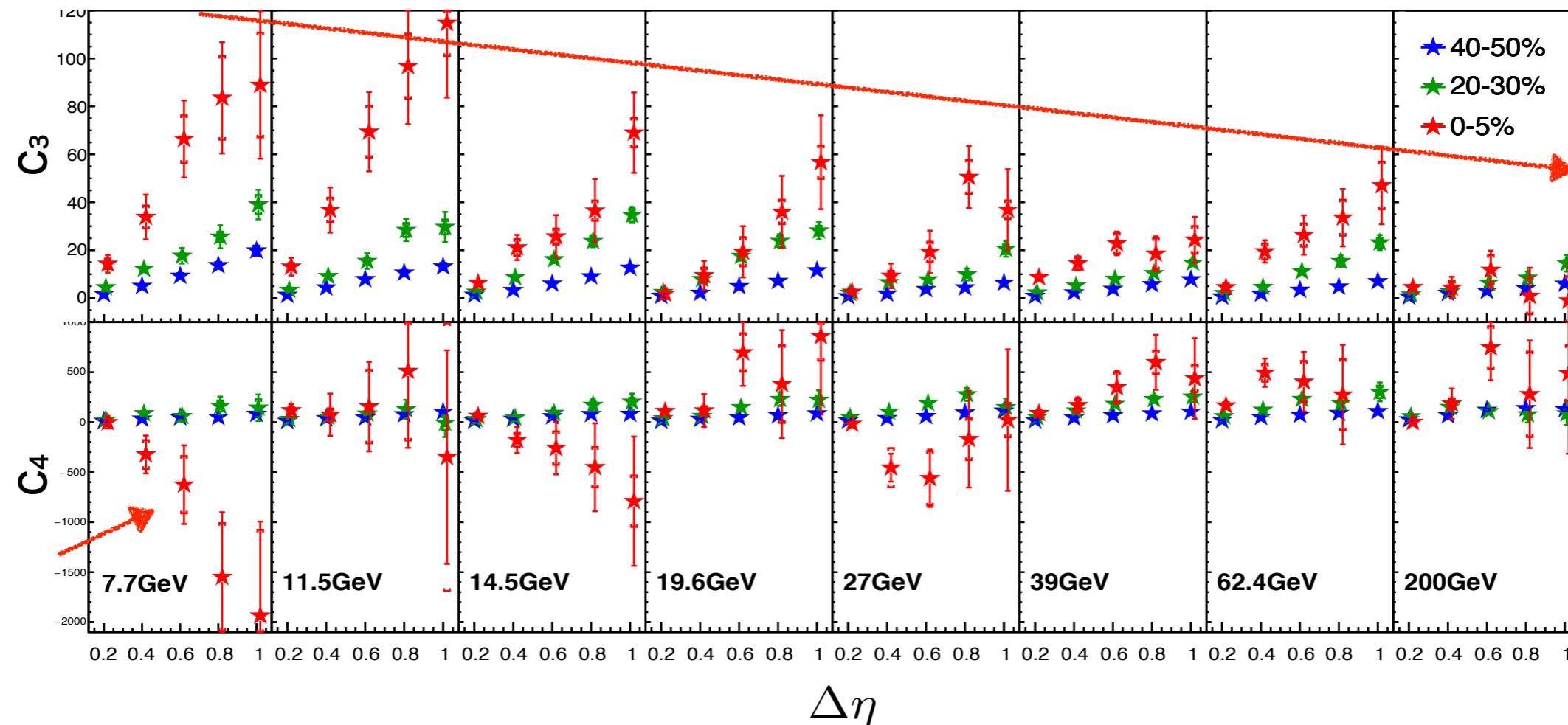
71



- Cumulants are observed to increase with $\Delta\eta$ because of additivity.
- Also increase from peripheral to central collisions.
- C_1 is observed to decrease with collision energies because of the baryon stopping.
- C_2 is observed to increase with collision energies because multiplicities are larger with collision energies.

$\Delta\eta$ dependence of C_3 and C_4

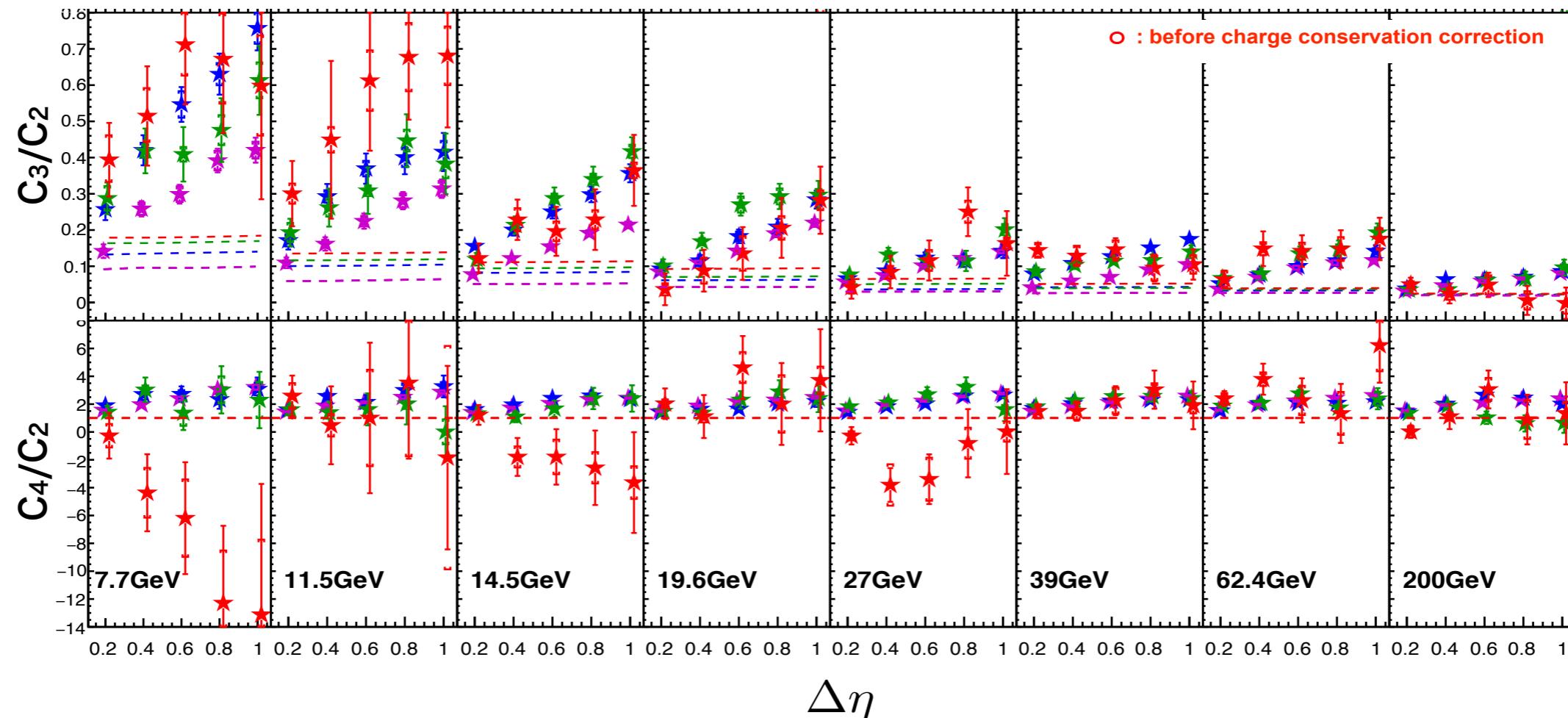
72



- Cumulants are observed to increase with $\Delta\eta$ because of additivity.
- Also increase from peripheral to central collisions.
- C_3 is observed to decrease with collision energies which trend is similar to C_1 .
- Statistical error of C_4 in 0-5% centralities are much larger than that in peripheral collisions.

$\Delta\eta$ dependence of cumulant ratios

73



- At C_3/C_2 and C_4/C_2 , most of the results are observed to increase with $\Delta\eta$ without most central collision at C_4/C_2 GeV.

back up (VFC)

Volume fluctuation correction (5,6-th order)

75

True cumulant

$$\kappa_5(\Delta N) = \langle N_W \rangle \kappa_5(\Delta n) + \{5\kappa_4(\Delta n)\kappa_1(\Delta n) + 10\kappa_3(\Delta n)\kappa_2(\Delta n)\} \kappa_2(N_W)$$

$$+ \{10\kappa_3(\Delta n)\kappa_1^2(\Delta n) + 15\kappa_2^2(\Delta n)\kappa_1(\Delta n)\} \kappa_3(N_W) + 10\kappa_2(\Delta n)\kappa_1^3(\Delta n)\kappa_4(N_W)$$

$$+ \kappa_1^5(\Delta n)\kappa_5(N_W)$$

Additional term caused from N_W fluctuation

T.Nonaka, Doctoral thesis

$$\kappa_6(\Delta N) = \langle N_W \rangle \kappa_6(\Delta n)$$

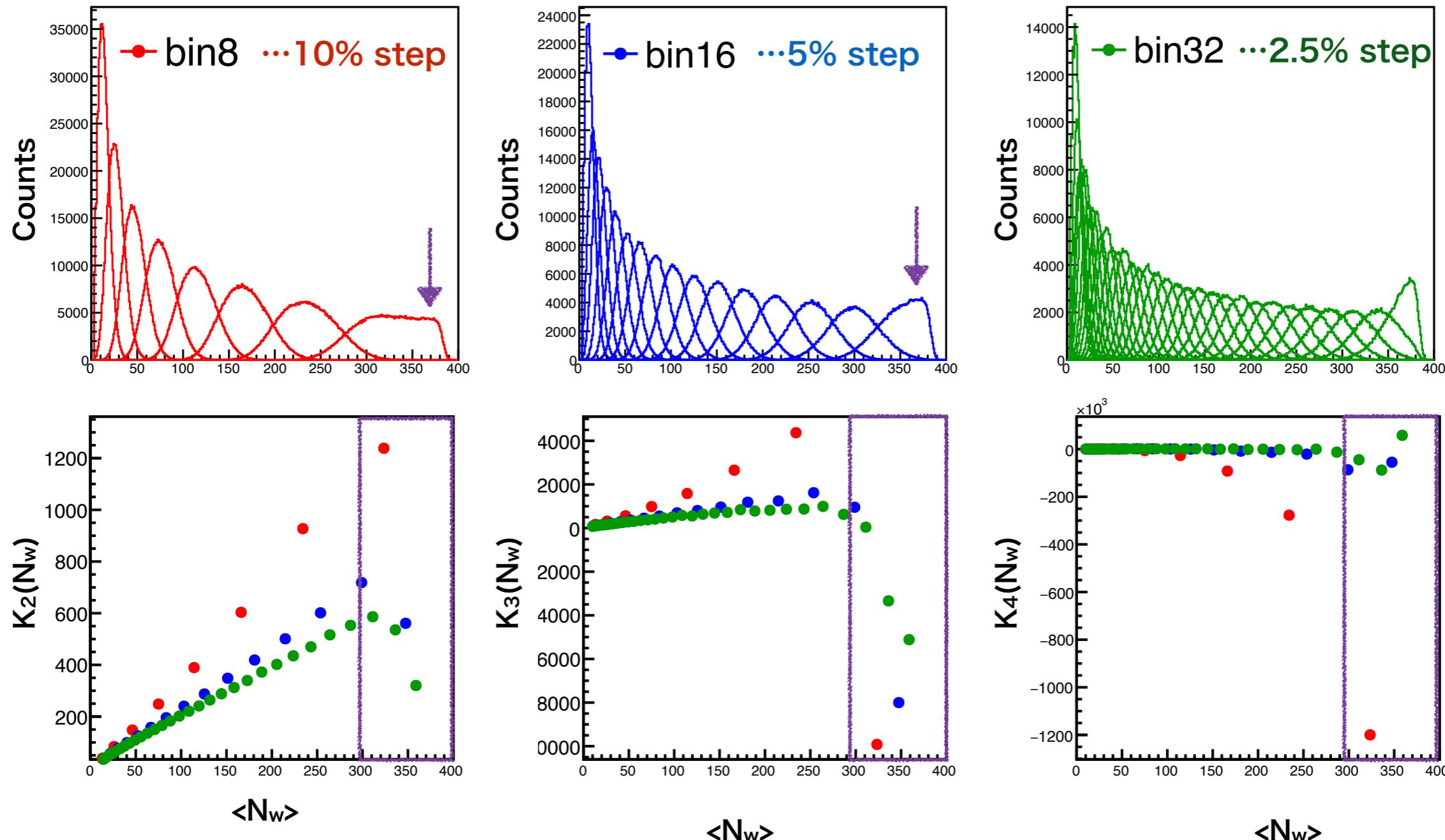
Δn : number of net-particle

Measured cumulant

$$+ \{6\kappa_5(\Delta n)\kappa_1(\Delta n) + 15\kappa_4(\Delta n)\kappa_2(\Delta n) + 10\kappa_3^2(\Delta n)\} \kappa_2(N_W)$$

$$+ \{15\kappa_4(\Delta n)\kappa_1^2(\Delta n) + 60\kappa_3(\Delta n)\kappa_2(\Delta n)\kappa_1(\Delta n) + 15\kappa_2^3(\Delta n)\} \kappa_3(N_W)$$

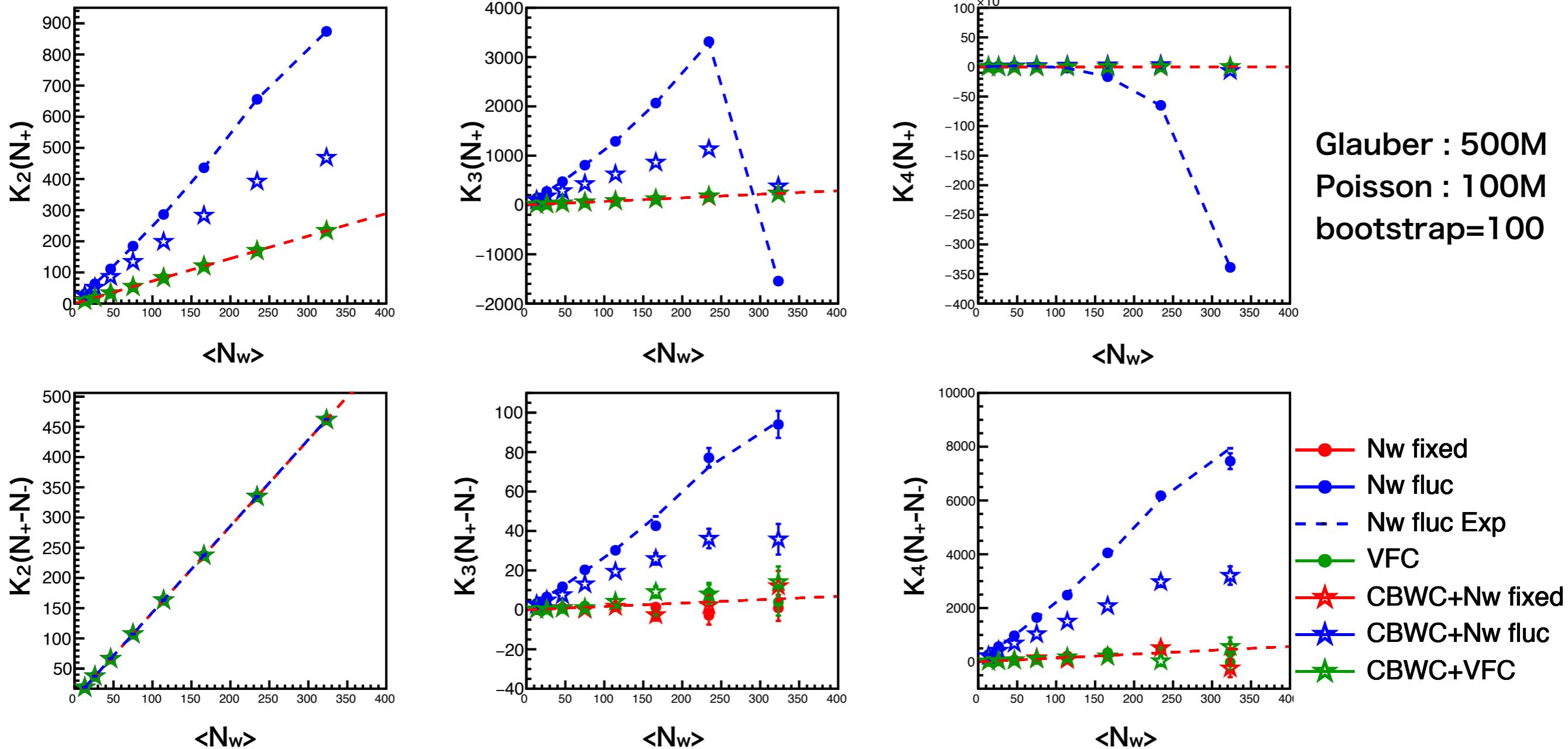
Participant fluctuation



- Trends are changed around central collision because maximum value of N_w is fixed.
- Participant fluctuation become larger with number of bin divisions become small.

Toy model VFC check (8bin)

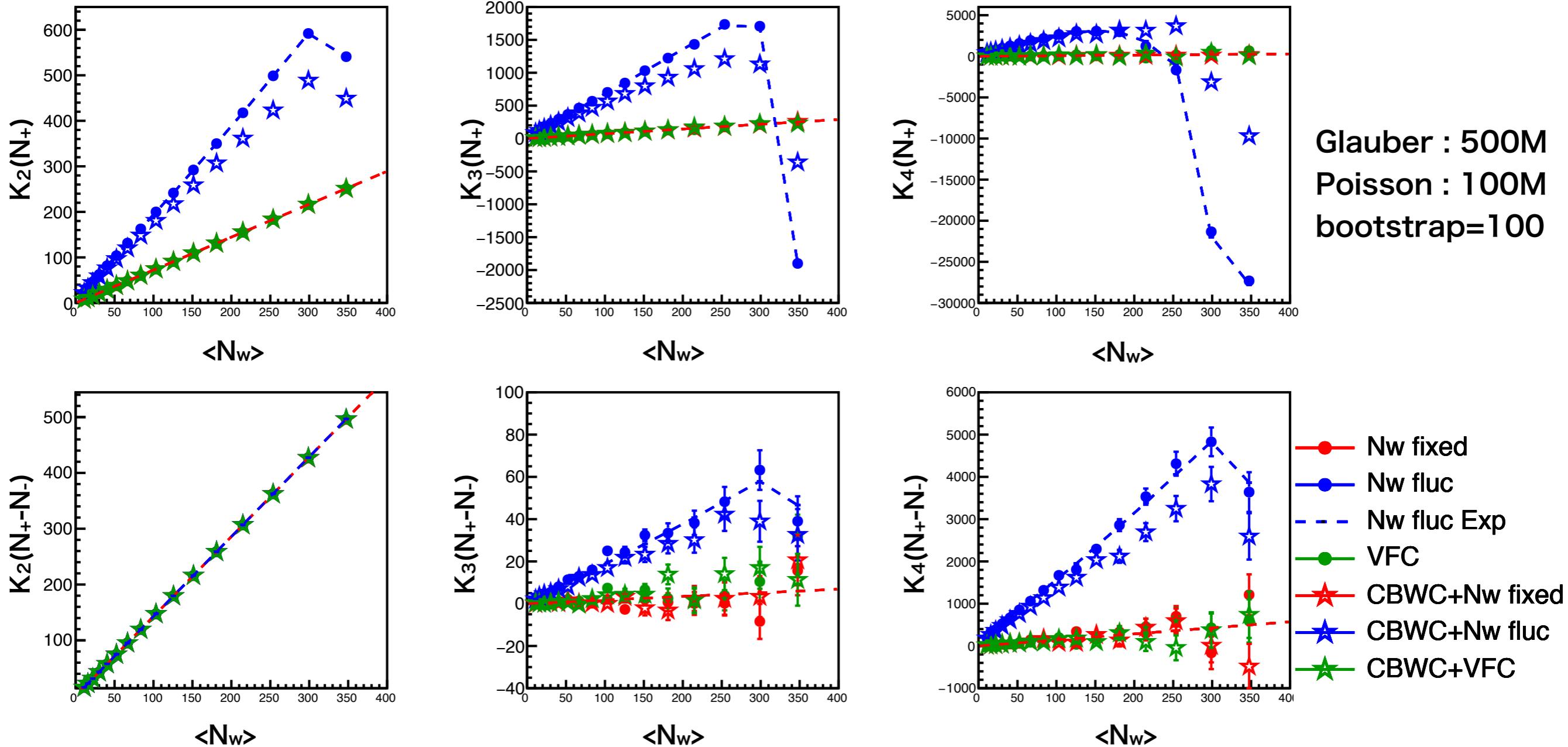
77



- It seems VFC works well.
- K_1 does not affected by participant fluctuation.
- In net-charge fluctuation, effect from participant fluctuation is smaller than N_+ cumulants because Δn of net-charge is smaller than N_+ .

Toy model VFC check (16bin)

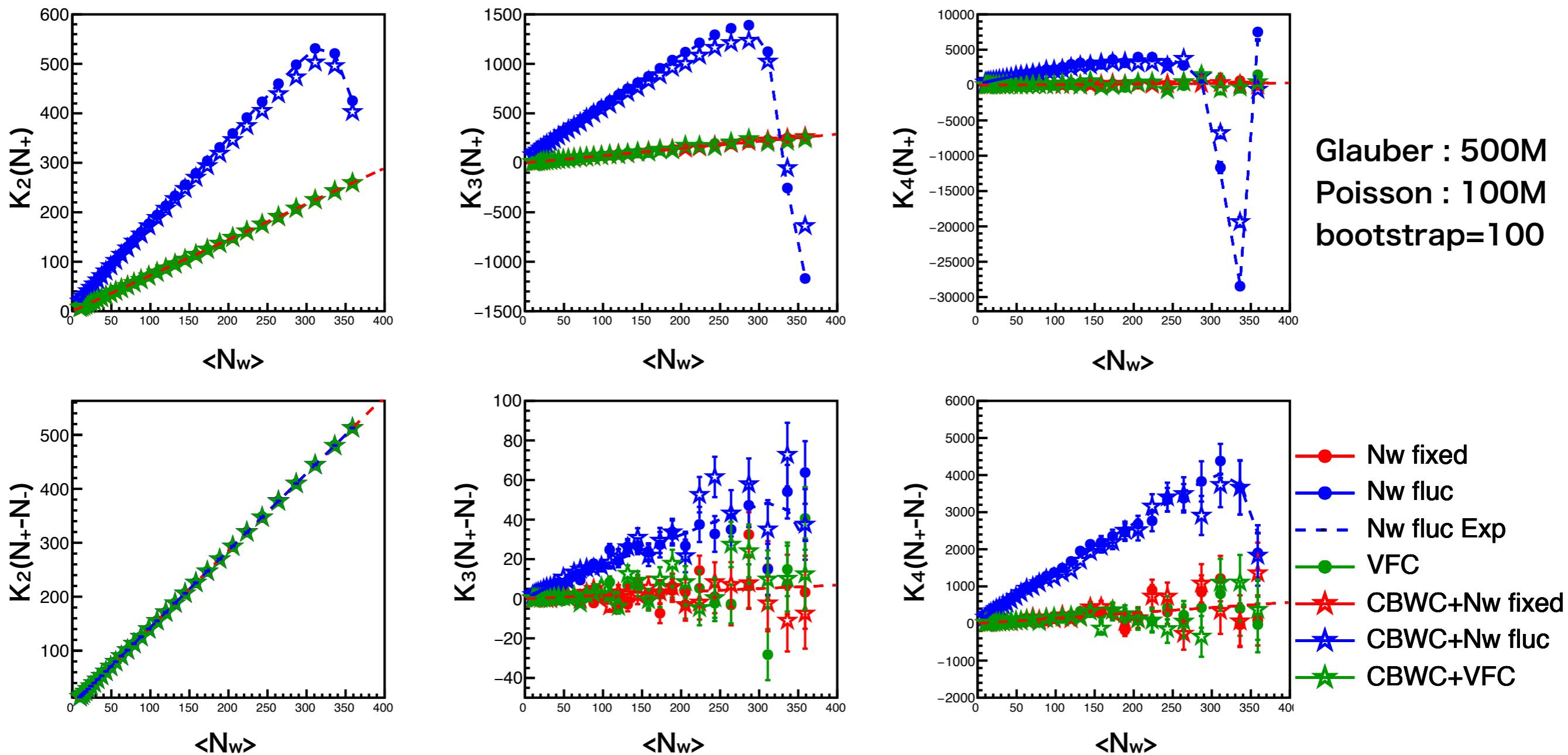
78



- It seems VFC works well.
- K_1 does not affected by participant fluctuation.
- In net-charge fluctuation, effect from participant fluctuation is smaller than N_+ cumulants because Δn of net-charge is smaller than N_+ .

Toy model VFC check (32bin)

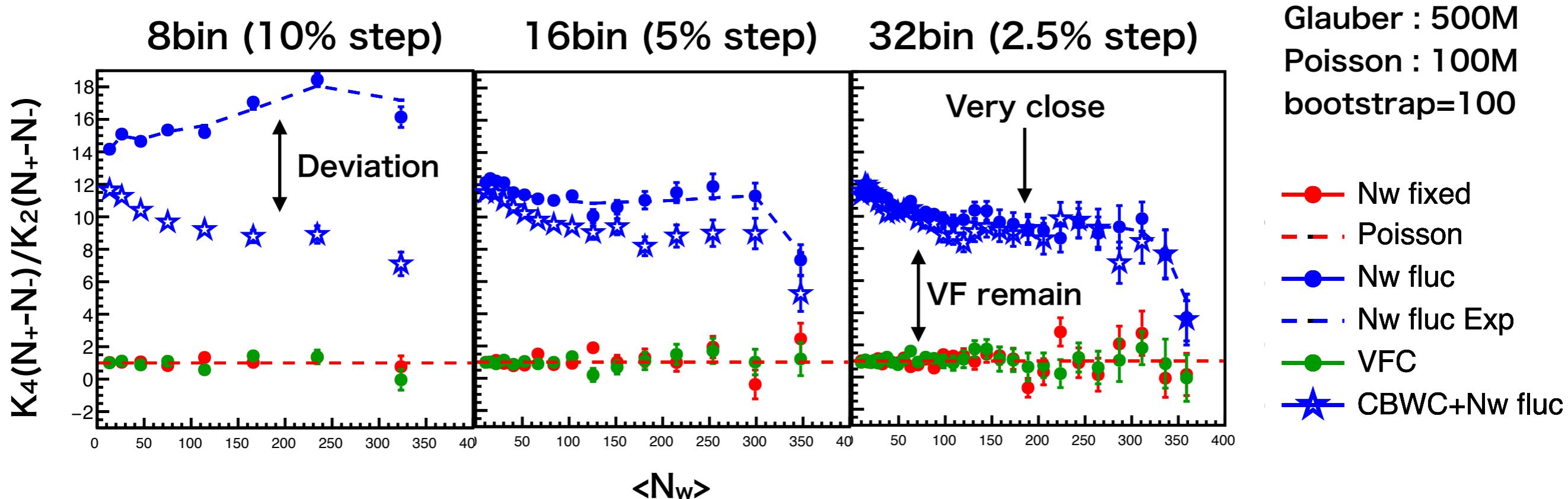
79



- It seems VFC works well.
- K_1 does not affected by participant fluctuation.
- In net-charge fluctuation, effect from participant fluctuation is smaller than N_+ cumulants because Δn of net-charge is smaller than N_+ .

Toy model VFC check (bin dependence)

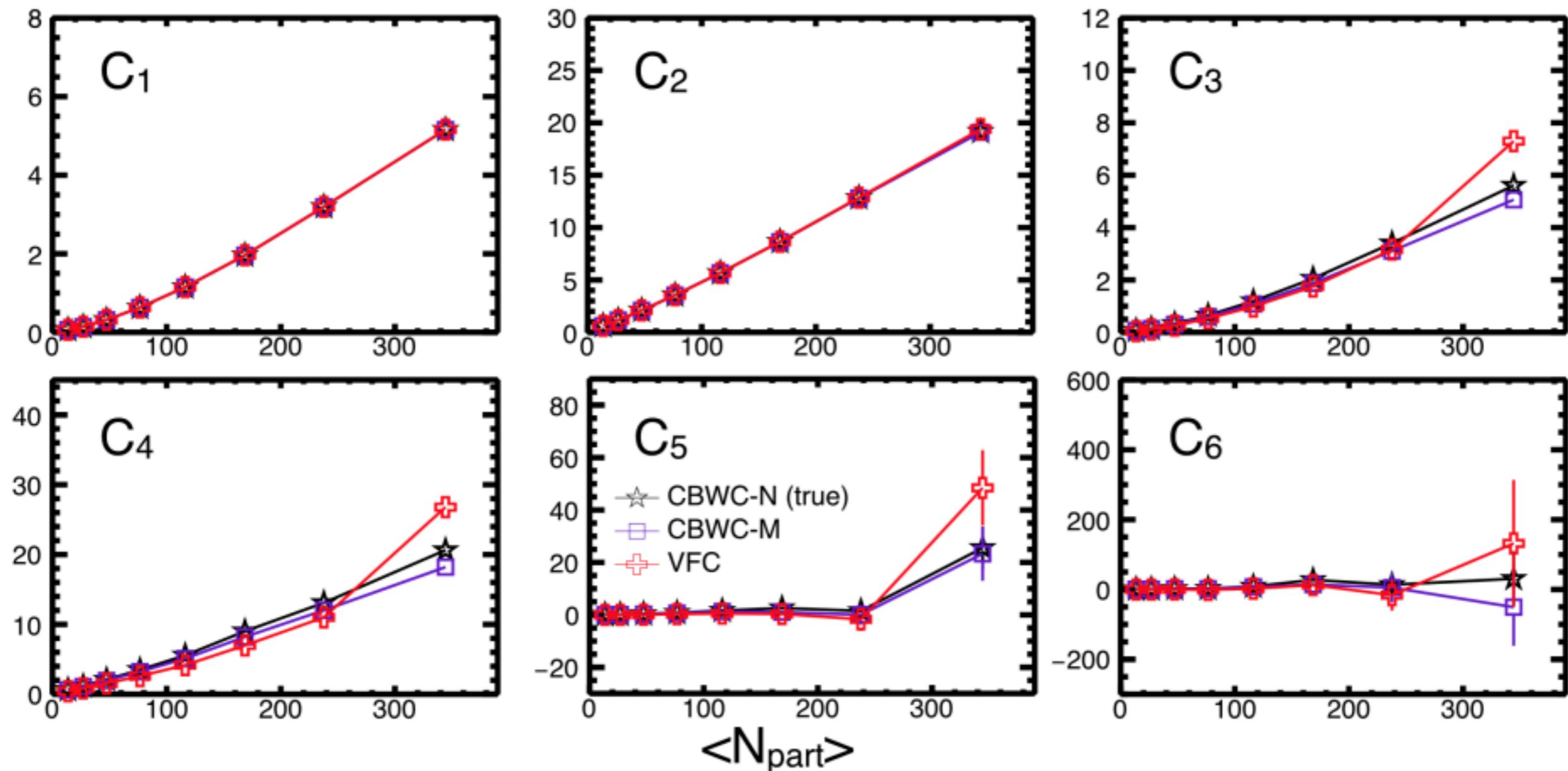
80

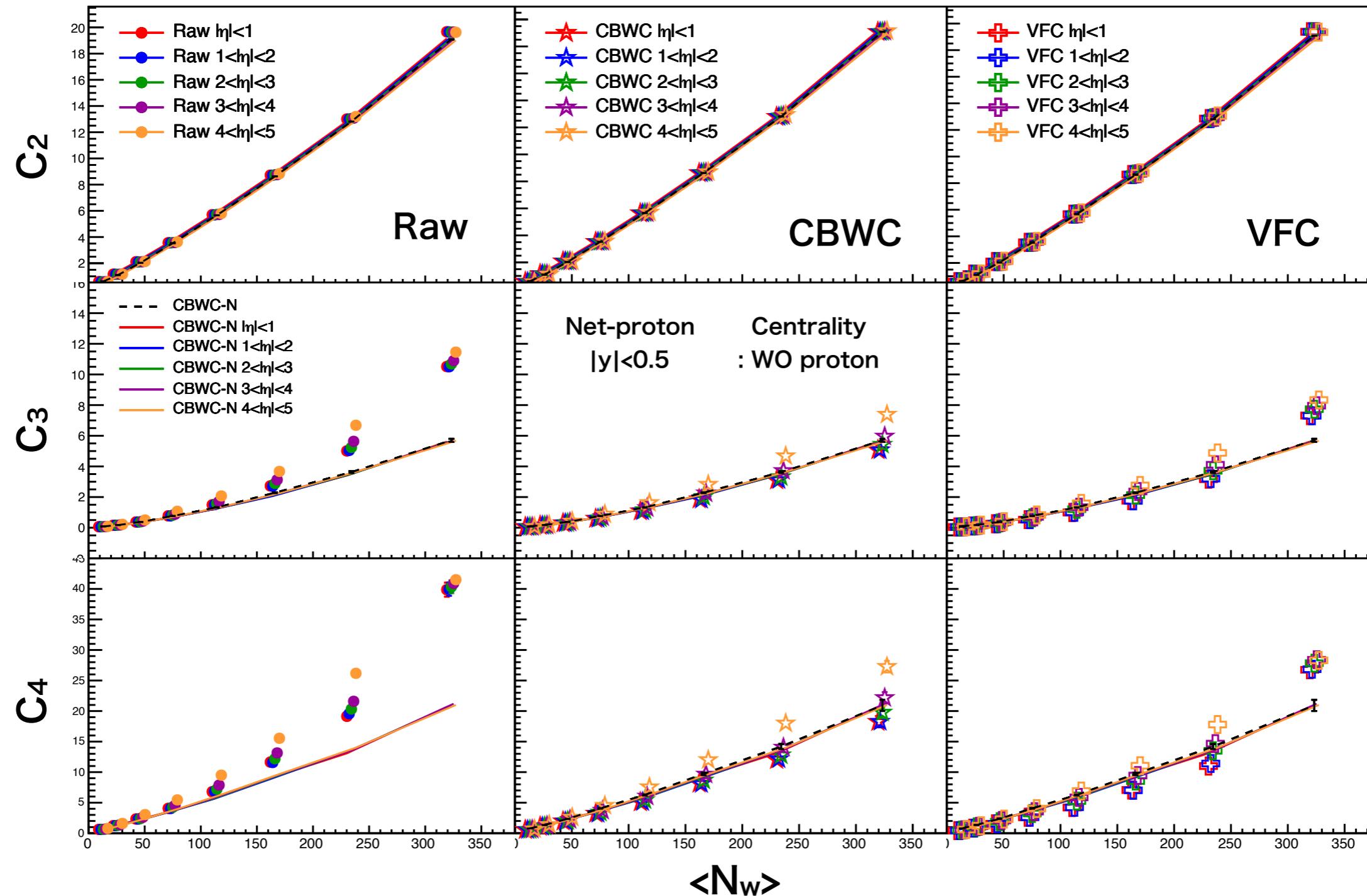


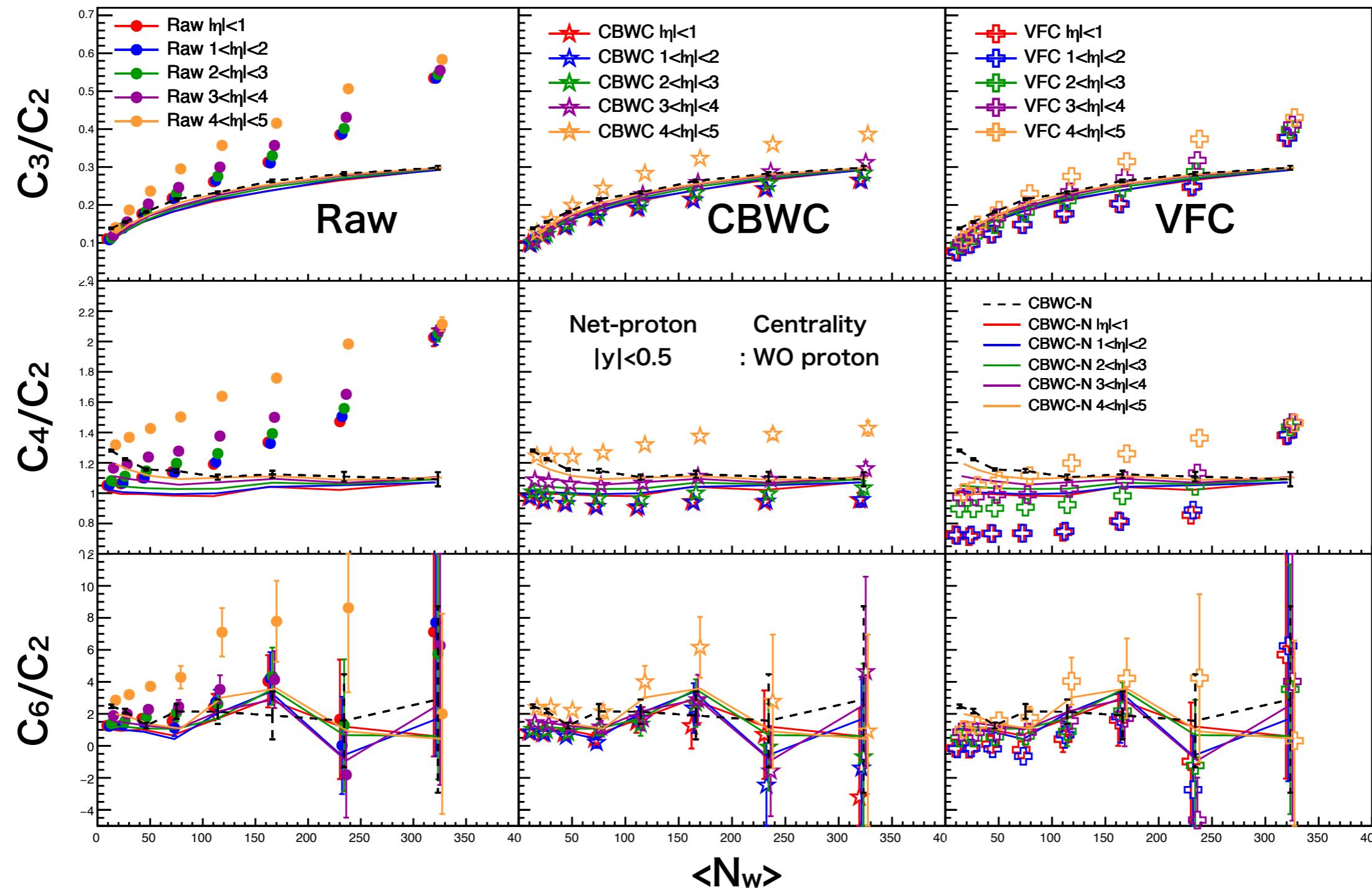
- CBWC cumulants and Without CBWC cumulants become very close if centrality step become 2.5%.

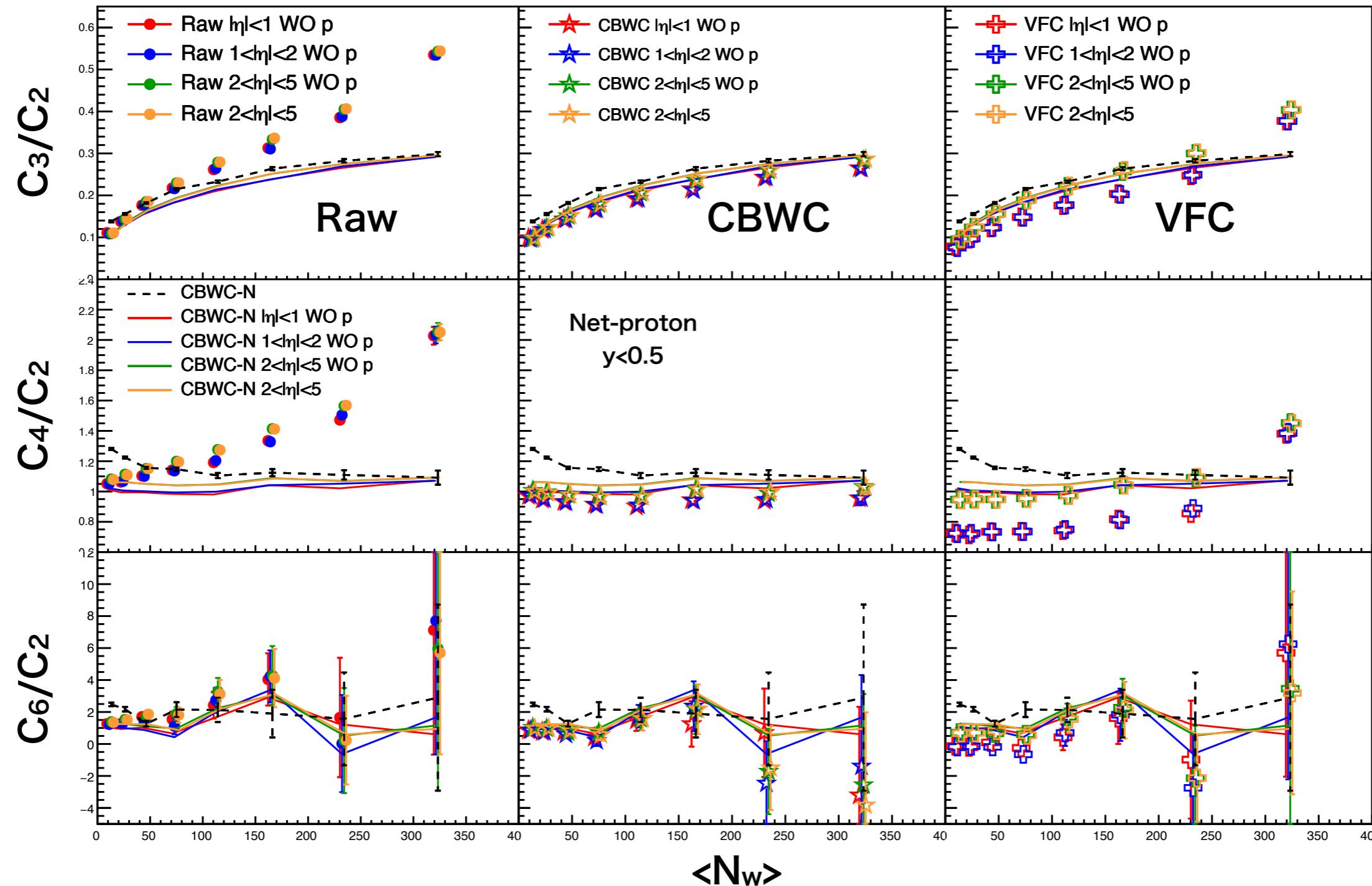
Net-proton case

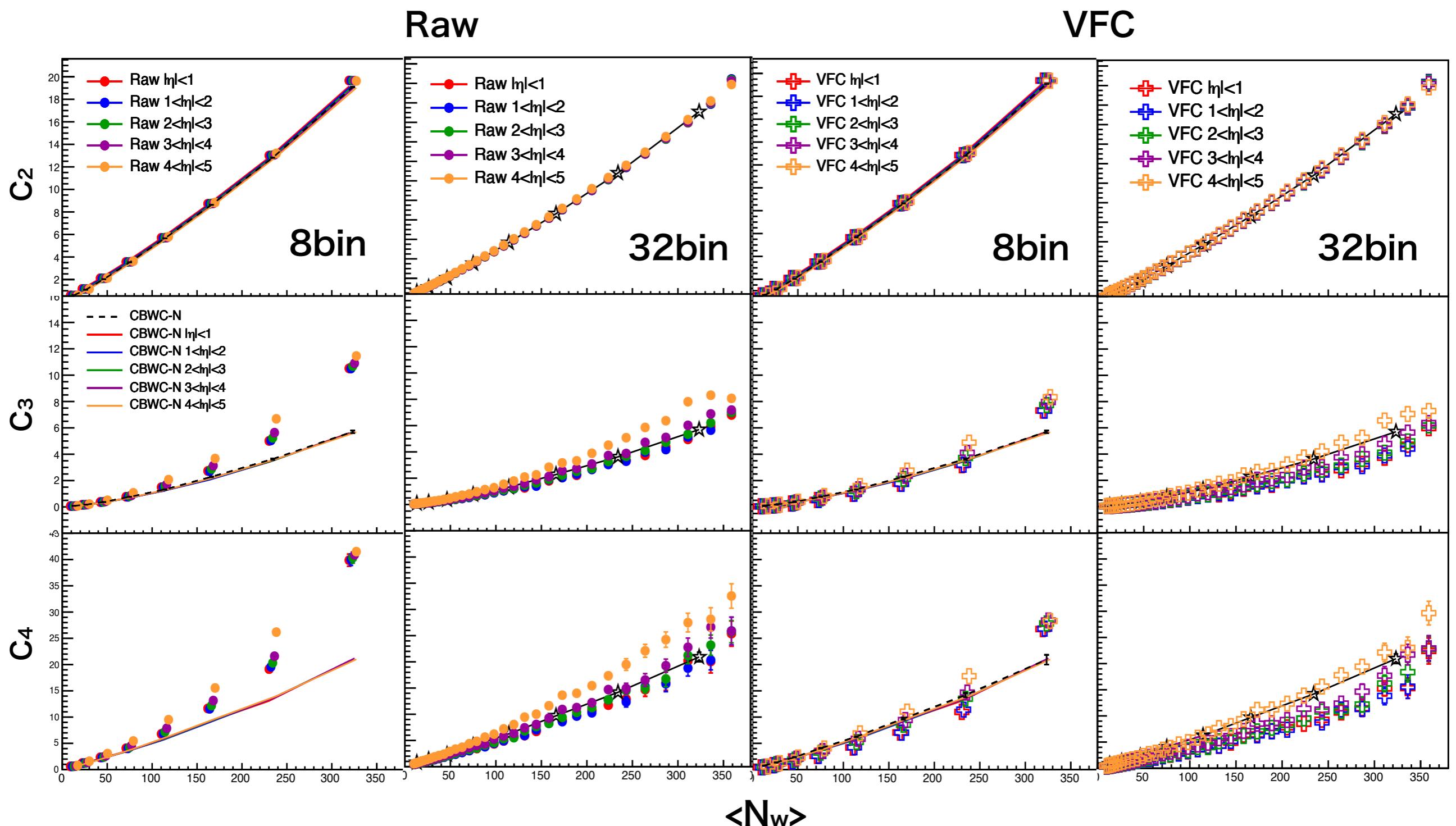
T.Nonaka, Doctoral thesis









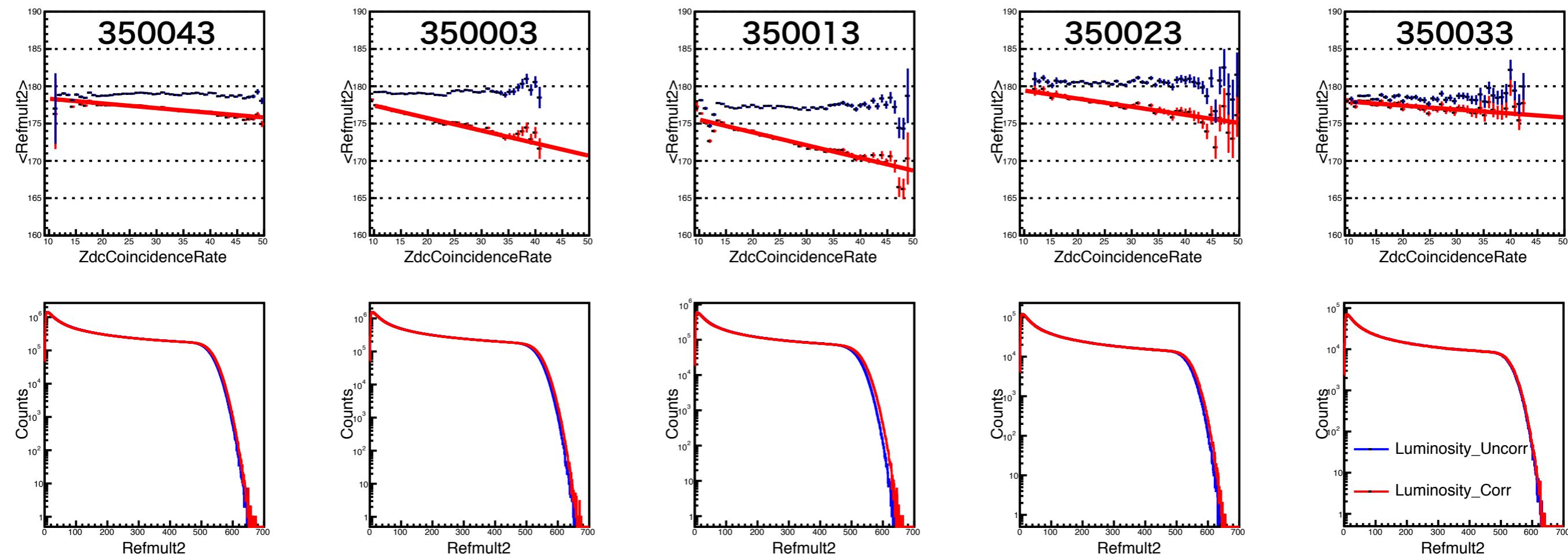


back up
(Centrality)

Correction parameter of Refmult2 (Run 11)

87

- Refmult2 was used to determine centrality at net-charge analysis.
(multiplicity of $0.5 < |\eta| < 1$)
- Refmult2 depend on z-vertex and luminosity and trigger ID.
- Calculate z-vertex and luminosity correction parameter for each trigger ID.



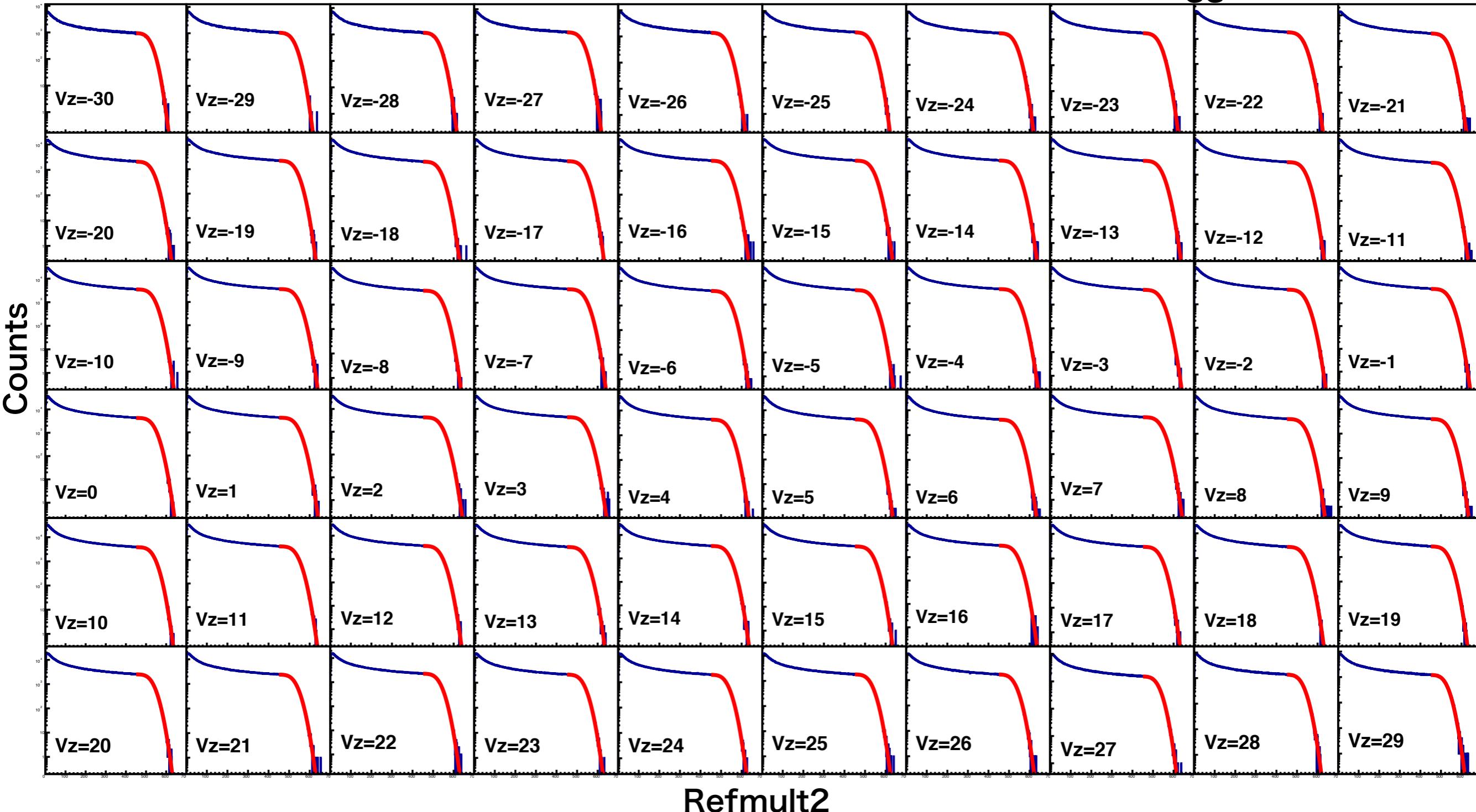
- Fit by pol1 ($[0]+[1]x$)

$$\text{correction_luminosity} = \frac{1}{1 + [1]/[0] * \text{ZdcCoincidenceRate}}$$

luminosity correction factor

Refmult2 distributions of each Vz

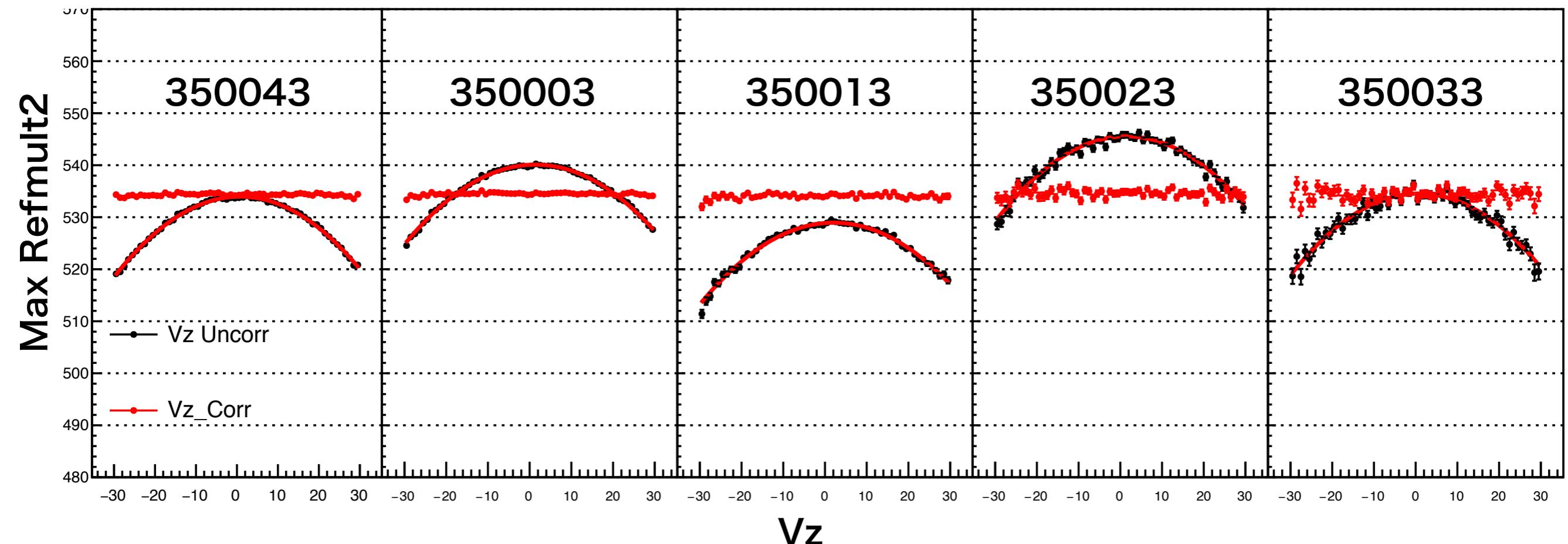
- Maxrefmult2 was measured that is defined by error function for each Vz bins from -30 to 30 after luminosity corrections. **Trigger ID = 350043**



- Fit by $([1]*\text{TMath::Erf}(-[1]*(x-[2]))+[0]) \rightarrow [2] : \text{Maxrefmult2}$

Z-vertex correction (Run11)

8a



- Fit by pol2 ([0]+[1]x+[2]x^2)

$$Hovno = \frac{[0] + \text{par}}{[0] + [1] * x + [2] * x^2}$$

Vz correction factor

([0] of 350043) - [0]
Scaled to 350043

Refmult2Corr = (Refmult2+gRandom->Rndm())*correction_lunimosity*Hovno

Corrected
Refmult2

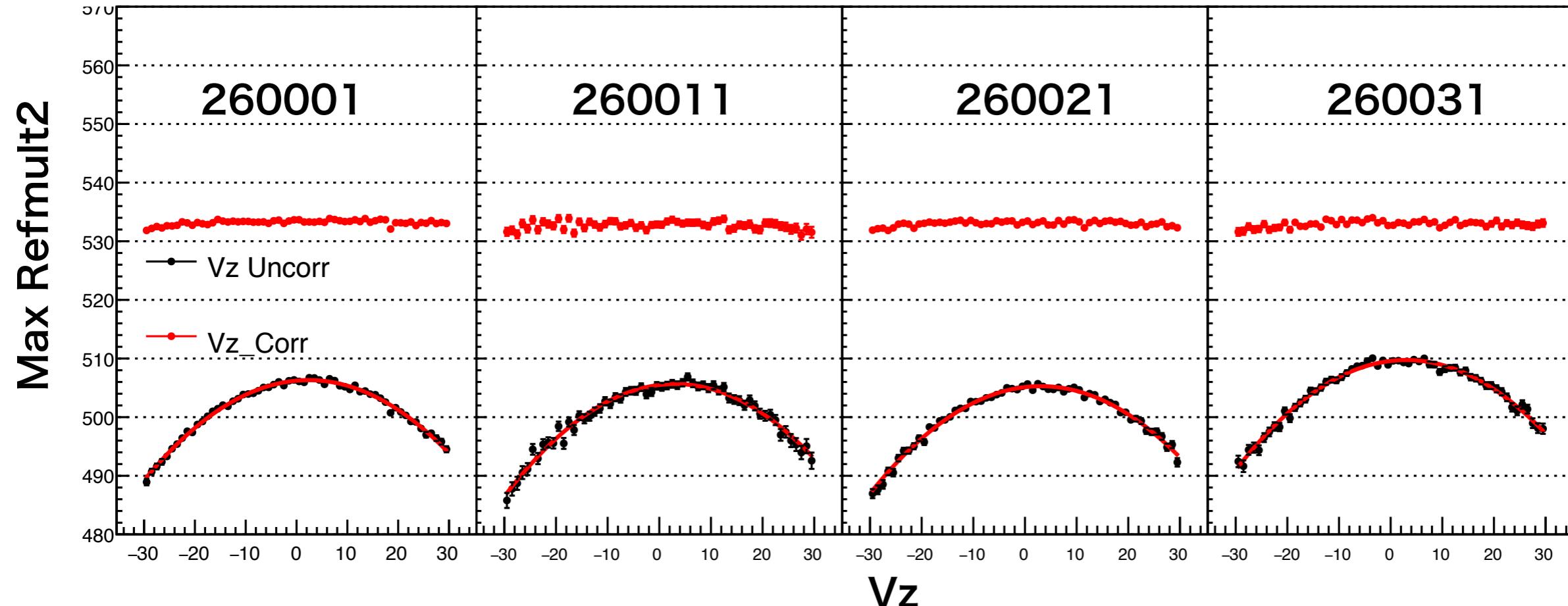
Raw Refmult2

0 to 1 random number
to remove kink structure

Luminosity
correction factor

Z-vertex correction (Run10)

90



- Fit by pol2 ([0]+[1]x+[2]x^2)

$$Hovno = \frac{[0] + par}{[0] + [1] * x + [2] * x^2}$$

Vz correction factor

([0] of 350043) - [0]
Scaled to 350043

Refmult2Corr = (Refmult2+gRandom->Rndm())*correction_lunimosity*Hovno

Corrected
Refmult2

Raw Refmult2

0 to 1 random number
to remove kink structure

Luminosity
correction factor

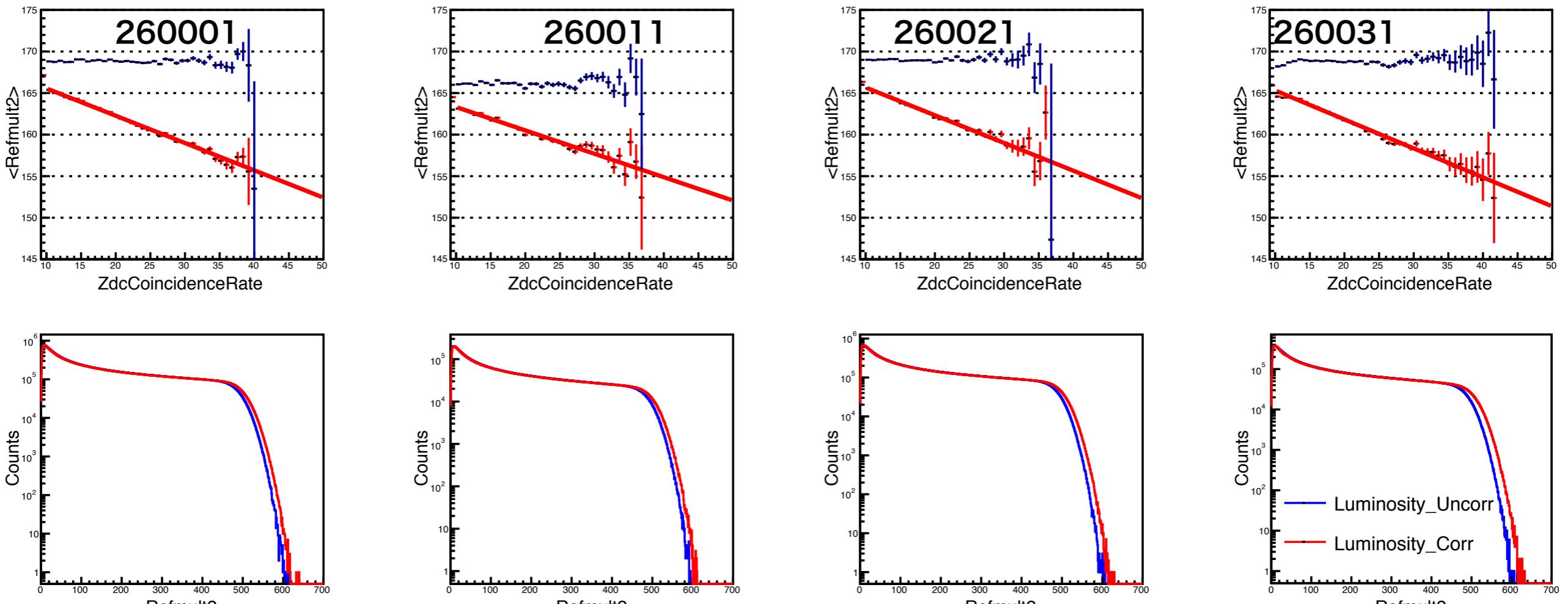
Vz correction
factor

Using StRefmultCorr class

Correction parameter of Refmult2 (Run 10)

91

- Refmult2 was used to determine centrality at net-charge analysis.
(multiplicity of $0.5 < |\eta| < 1$)
- Refmult2 depend on z-vertex and luminosity and trigger ID.
- Calculate z-vertex and luminosity correction parameter for each trigger ID.

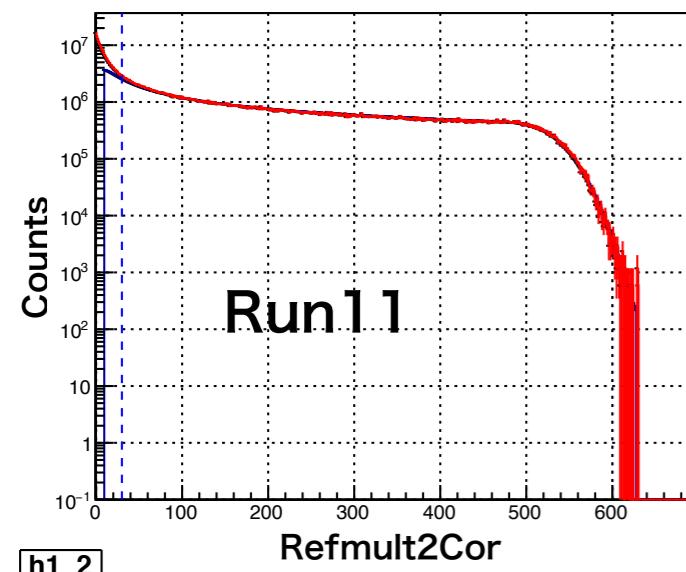


- Fit by pol1 ($[0] + [1]x$)

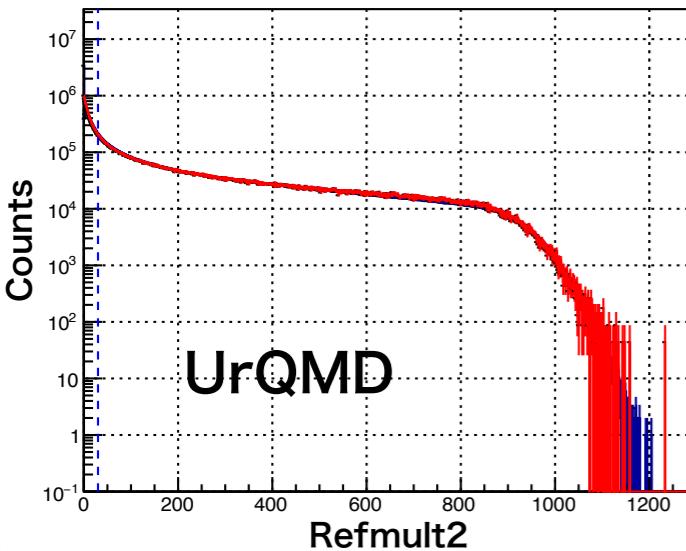
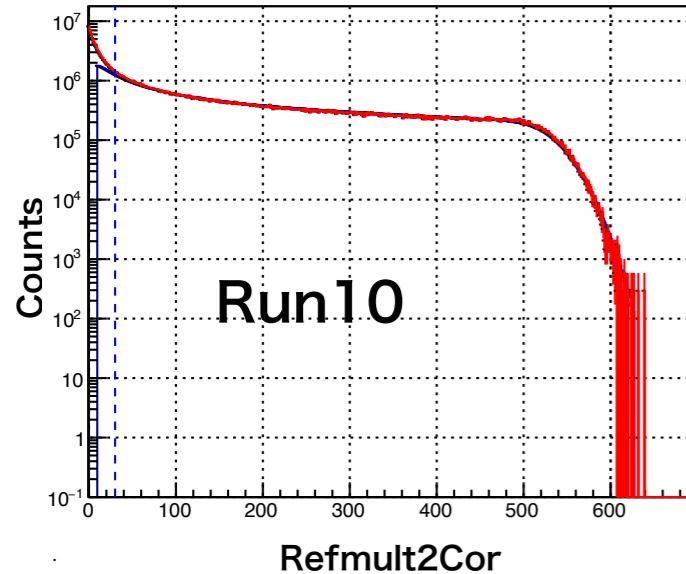
$$\text{correction_luminosity} = \frac{1}{1 + [1]/[0] * \text{ZdcCoincidenceRate}}$$

luminosity correction factor

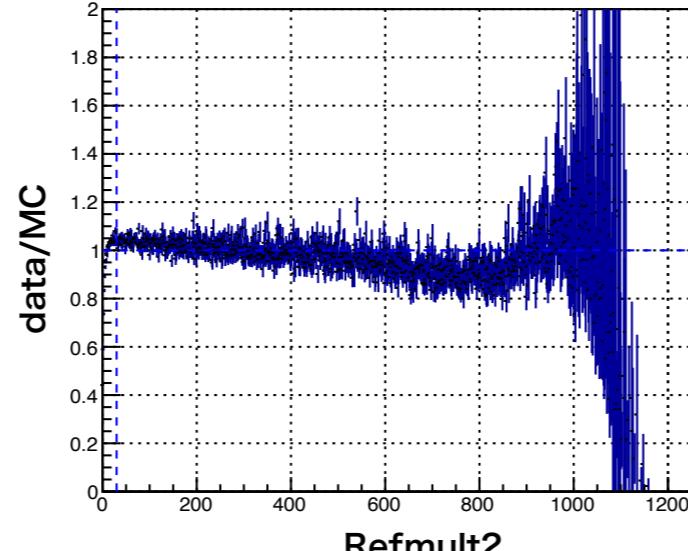
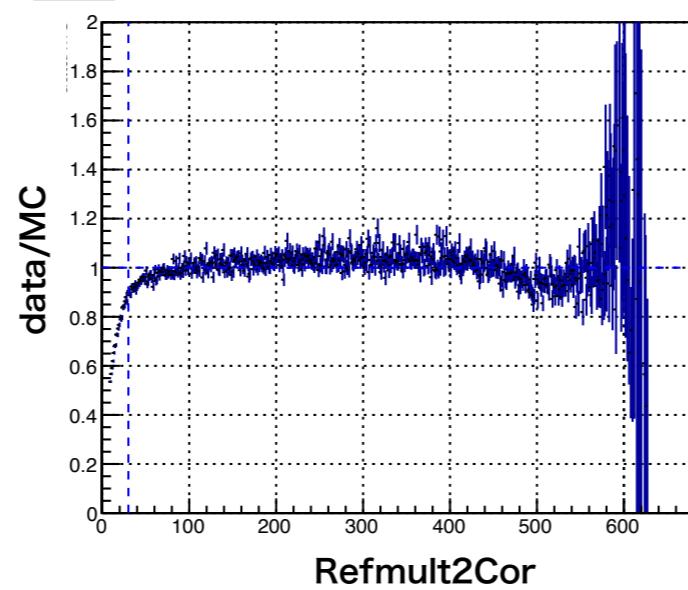
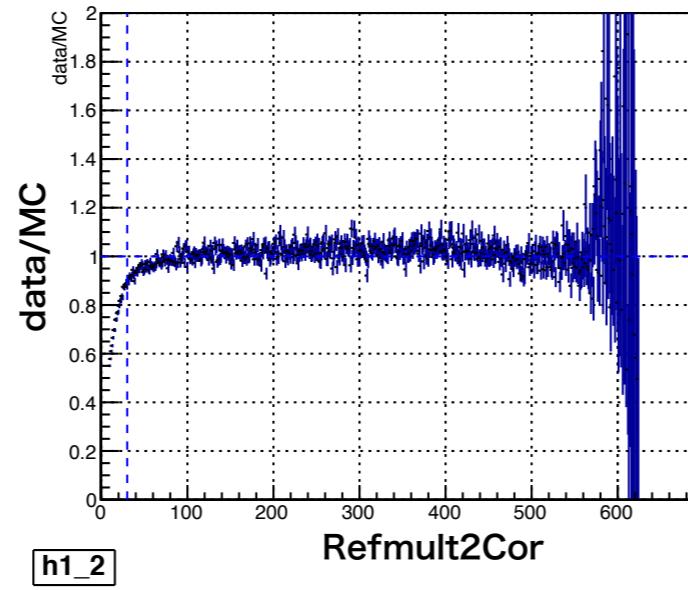
Glauber fit results



h1_2



Defense 2019/2/8



MC : 1M

Experiment

- $n_{pp}=2.35$
 - $x=0.13$
 - $k=0.9$
 - efficiency
- $$=1-\text{Refmult2Cor}^*(0.206/560)$$

n_{pp} : Mean number of generated particles from each source.

k : Parameters of NBD.

x : Parameter of two component model.

Number of source :

$$(1-x)*(N_{part}/2.)+x*N_{coll}$$

UrQMD

- $n_{pp}=2.90$
- $x=0.13$
- $k=0.9$
- efficiency=1

Tetsuro Sugiura Univ. of Tsukuba

Negative binomial distribution

Binomial distribution

r success , given n trials

$$f(r; n, p) = \binom{n}{r} p^r (1-p)^k$$

r : success

k : failures

n : trials

p : probability of success

Negative binomial distribution (NBD)

r success before k failures

$$f(r; k, p) = \binom{r+k-1}{r} p^r (1-p)^k$$

Expand k and r to positive real values.

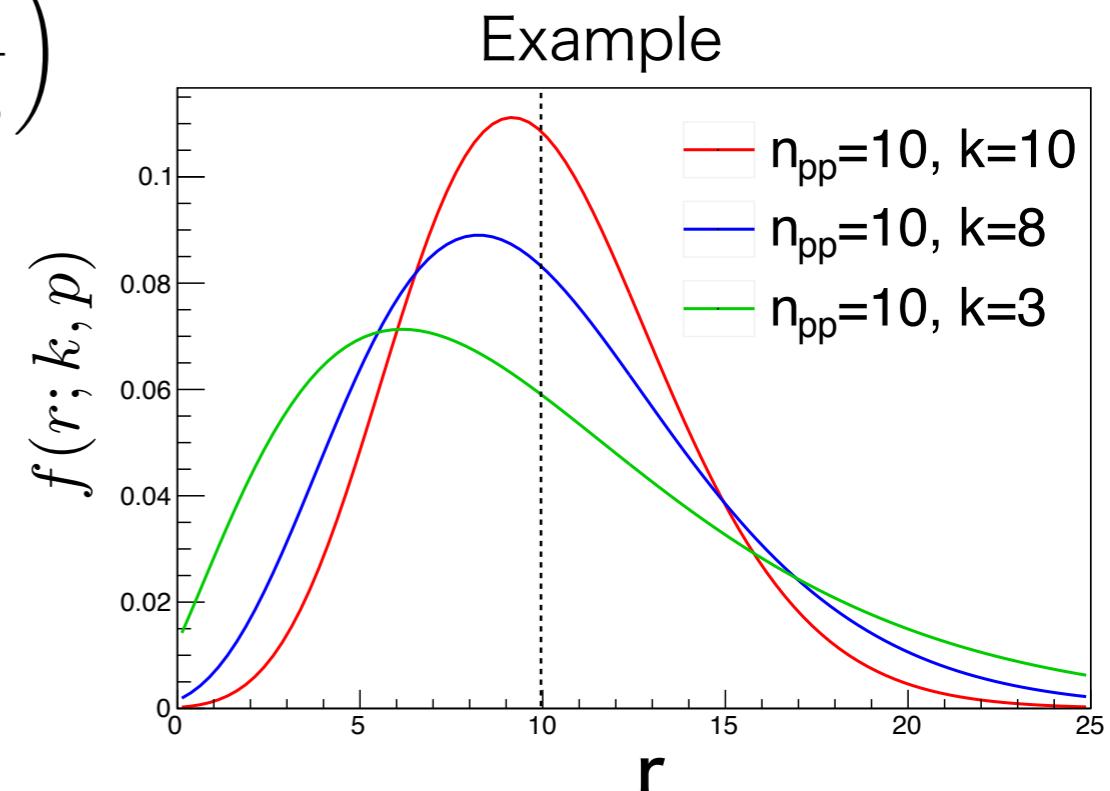
$$f(r; k, p) = \frac{\Gamma(r+k)}{\Gamma(r+1)\Gamma(k)} \left(\frac{n_{pp}}{k+n_{pp}} \right)^r \left(\frac{k}{k+n_{pp}} \right)^k$$

$$\frac{\langle r \rangle}{p} = \frac{k}{1-p} = \text{number of trials}$$

$$\rightarrow \langle r \rangle = \frac{pk}{1-p} \equiv n_{pp}$$

n_{pp} : Mean number of generated particles from each ancestors.

k : Parameters of NBD.



Centrality bins

8 bin

Centrality	bin	$\langle N_{\text{part}} \rangle$
0-10	406	323.41
10-20	295	234.16
20-30	204	166.14
30-40	134	114.3
40-50	82	75.09
50-60	47	46.32
60-70	24	26.14
70-80	11	13.15

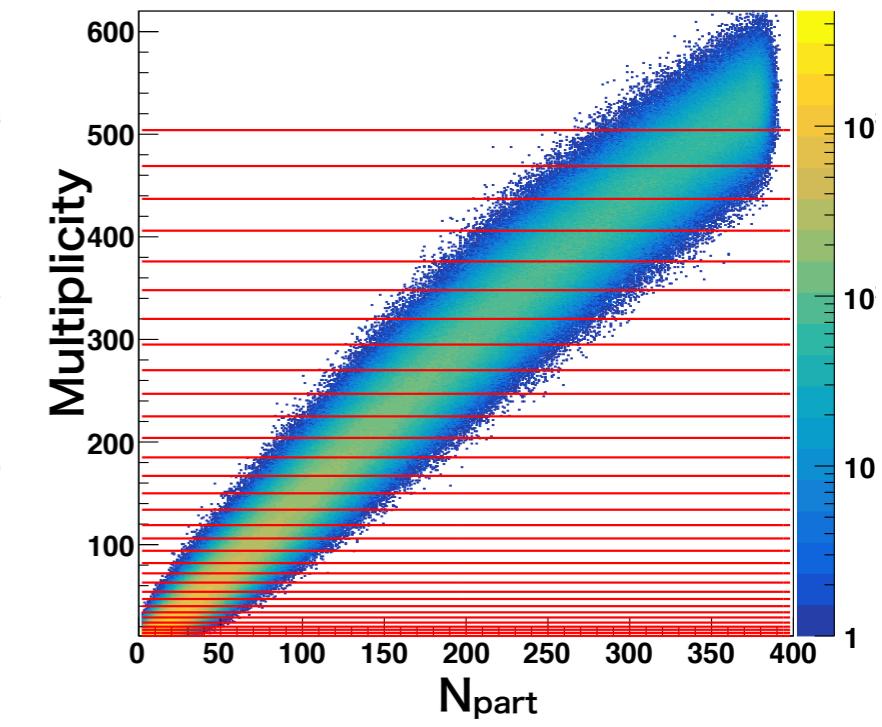
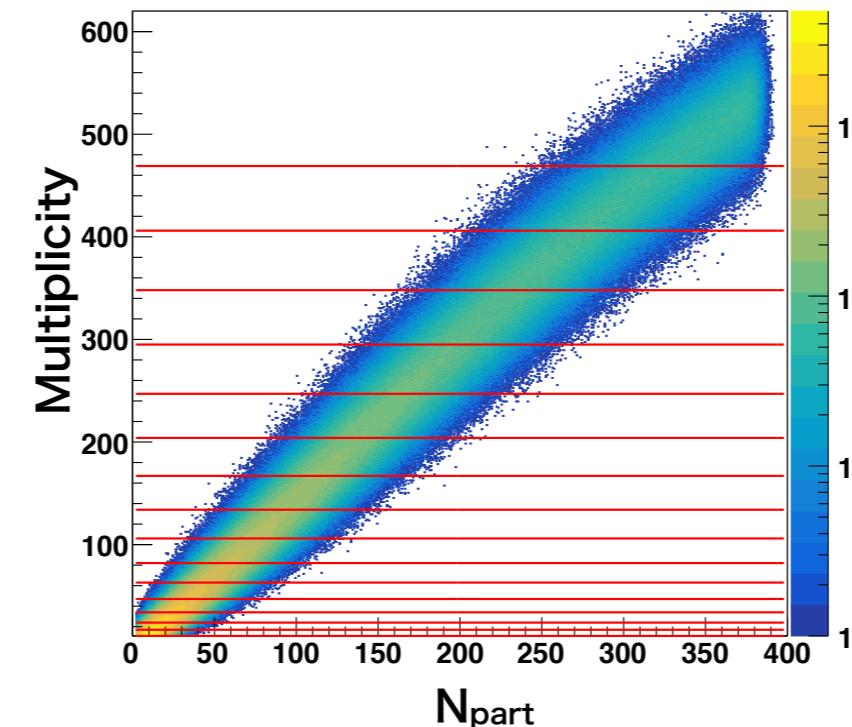
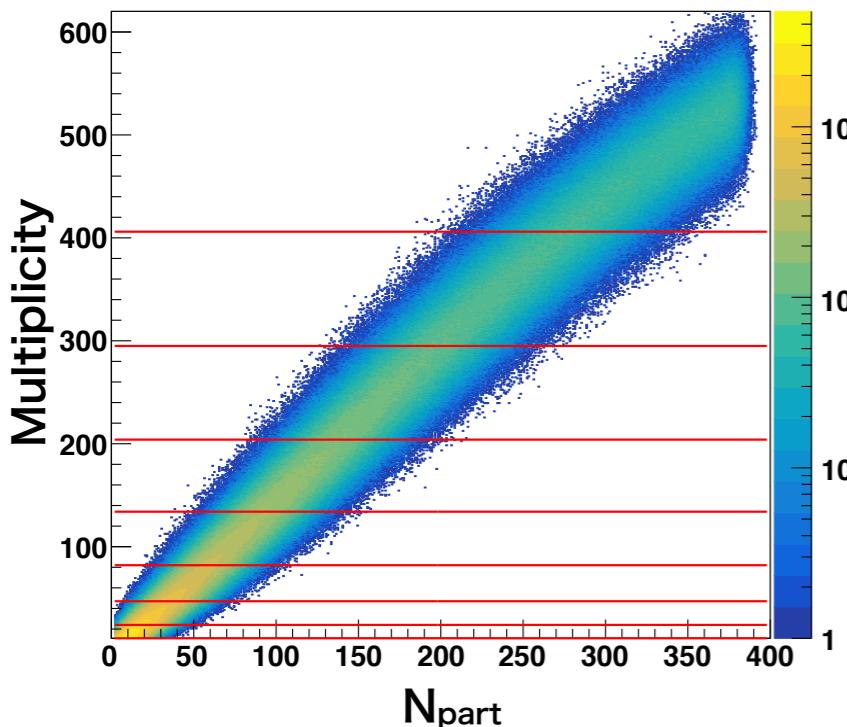
16 bin

Centrality	bin	$\langle N_{\text{part}} \rangle$
0-5	469	347.47
5-10	406	298.9
10-15	348	253.64
15-20	295	214.79
20-25	247	180.99
25-30	204	151.34
30-35	167	125.669
35-40	134	103.2
40-45	106	83.61
45-50	82	66.65
50-55	63	52.38
55-60	47	40.46
60-65	34	30.28
65-70	24	22.0
70-75	17	15.73
75-80	11	10.93

32 bin

Centrality	bin	$\langle N_{\text{part}} \rangle$
0-2.5	504	358.83
2.5-5	469	336.16
5-7.5	437	311.16
7.5-10	406	286.74
10-12.5	376	263.97
12.5-15	348	243.02
15-17.5	320	223.61
17.5-20	295	205.41
20-22.5	270	188.85
22.5-25	247	172.91
25-27.5	225	158.33
27.5-30	204	144.42
30-32.5	185	131.64
32.5-35	167	119.78
35-37.5	150	108.56
37.5-40	134	97.93
40-42.5	119	88.02

⋮



back up
(Data set, QA, Analysis mothod, etc)

Data set (200 GeV)

	Run11(0-80%)	Run10(10-80%)	Run10(0-10%)	UrQMD
NEvent	485M	211M	196M	45M

Event Selection	$ v_z $	<30
	$ v_r $	<2
	Pile up event cut	$0.46 * nRefMult - 10 < nTOFMatch$

Track cut	π	K	p
p_T (TPC only)	0.2 - 0.5 GeV/c	0.2 - 0.4 GeV/c	0.4 - 0.8 GeV/c
p_T (TPC+TOF)	0.5 - 1.6 GeV/c	0.4 - 1.6 GeV/c	0.8 - 2. GeV/c
$ \eta $		< 0.5	
dca		<1	
nhitsfits		>20	Same as net-charge published results
track quality		>0.52	
nhitsdedx		>10	
Centrality		Refmult2 ($ \eta > 0.5$)	
PID (TPC)	nSigmaPion < 2	nSigmaKaon < 2	nSigmaProton < 2
PID (TOF)	-0.15 < m^2 < 0.14	0.14 < $ m^2 $ < 0.4	0.6 < $ m^2 $ < 1.2

Number of Events

97

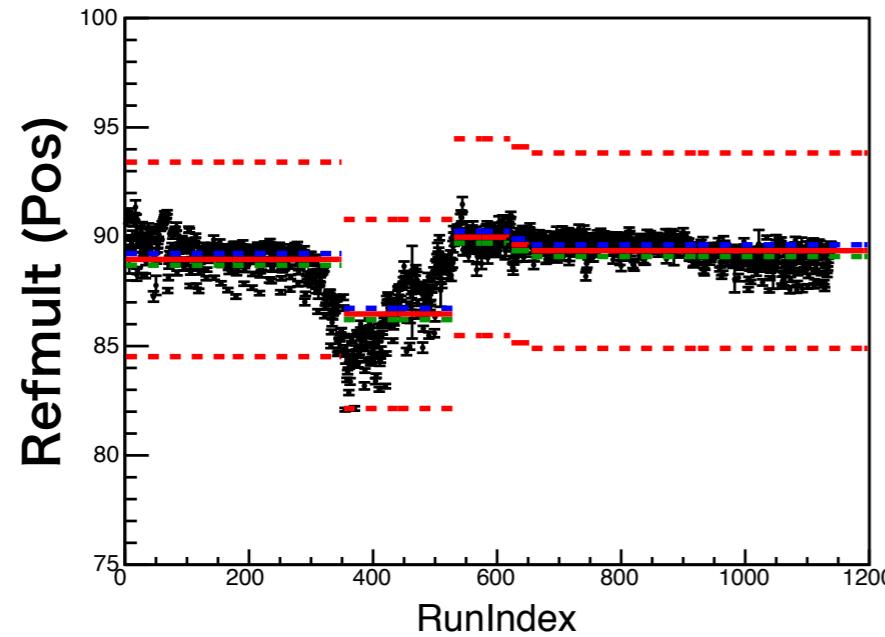
NEvent	350003	350013	350023	350033	350043	Total
Run11 (0-80%)	200M	74M	15M	9M	187M	485M

NEvent	260001	260011	260021	260031	Total
Run10 (10-80%)	80M	21M	70M	40M	211M
Run10 (0-10%)	53M	20M	76M	47M	196M

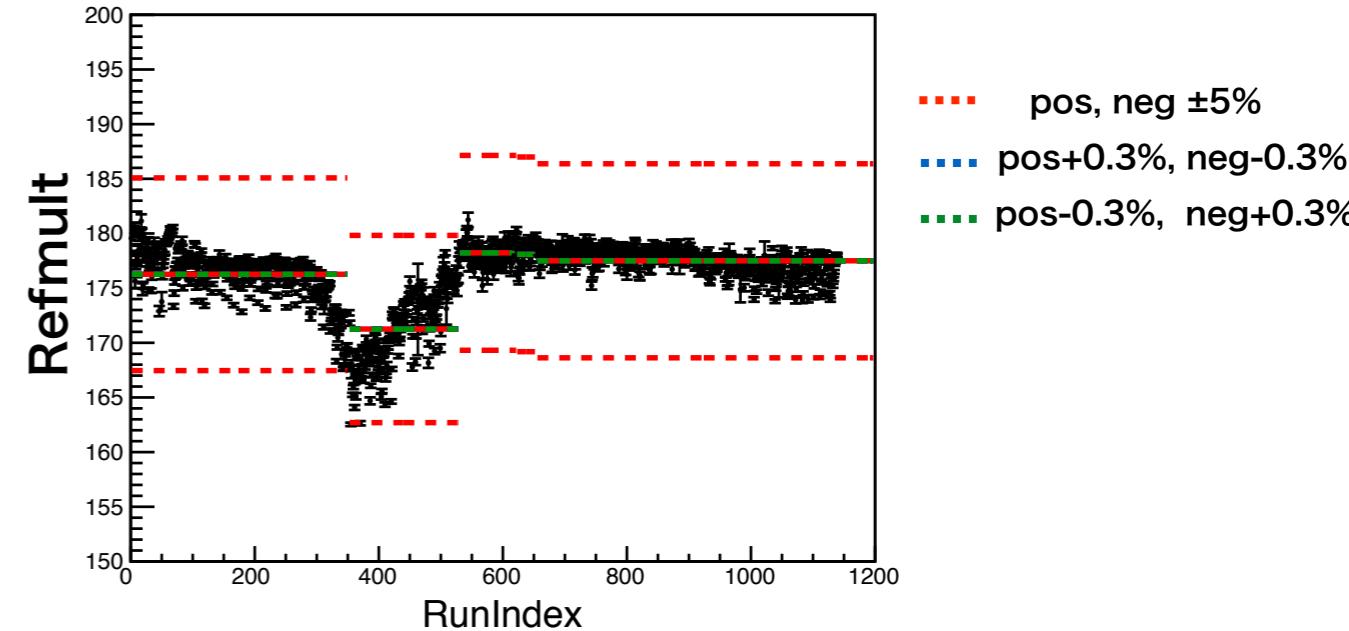
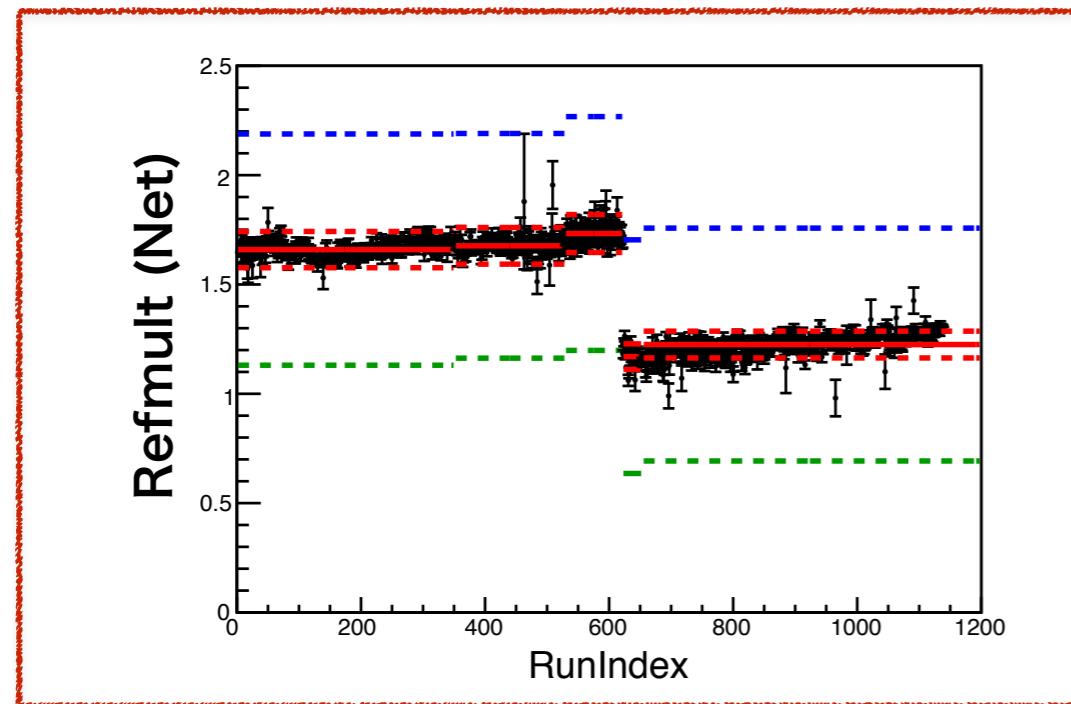
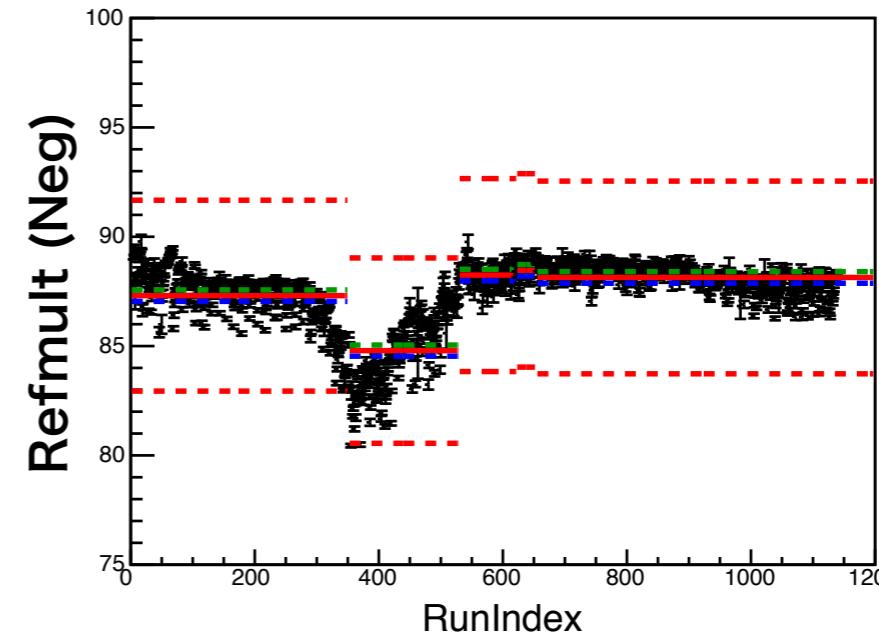
NEvent	Total
UrQMD	45M

Run by run check (Run11, 350043)

Only Good runs are shown



Refmult(Pos)+Refmult(Neg)=Refmult



- When we change $(\text{pos}, \text{neg}) = (+0.3\%, -0.3\%)$ and $(+0.3\%, -0.3\%)$, run by run Refmult (Net) are fluctuate within blue and green line.
- (pos+0.3%, neg-0.3%) and (pos-0.3%, neg+0.3%) seem enough to estimate systematic uncertainty of net-charge.**

Factorial cumulant method

Example (the 1st-order net-charge fluctuation)

Cumulant (true)
 $\langle Q \rangle_c = \langle q_{(1,1)} \rangle_c,$

$$q_{(r,s)} = q(a^r / p^s) = \sum_{i=1}^M (a_i^r / p_i^s) n_i.$$

+1 (positive charge) or -1 (negative charge)

M Number of phase space
n_i Number of particle
a_i^r / p_i^s Efficiency

$$\begin{aligned}
 &= \frac{n(\pi^+, lowpt)}{p(\pi^+, lowpt)} + \frac{n(\pi^+, highpt)}{p(\pi^+, highpt)} - \frac{n(K^+, lowpt)}{p(K^+, lowpt)} + \frac{n(K^-, highpt)}{p(K^-, highpt)} + \frac{n(pbar, lowpt)}{p(pbar, lowpt)} + \frac{n(p, highpt)}{p(p, highpt)} \\
 &\quad - \frac{n(\pi^-, lowpt)}{p(\pi^-, lowpt)} - \frac{n(\pi^-, highpt)}{p(\pi^-, highpt)} - \frac{n(K^-, lowpt)}{p(K^-, lowpt)} - \frac{n(K^-, highpt)}{p(K^-, highpt)} - \frac{n(pbar, lowpt)}{p(pbar, lowpt)} - \frac{n(pbar, highpt)}{p(pbar, highpt)}
 \end{aligned}$$

- $2(\text{charge}) * 3(\pi K p) * 2(\text{highpt}, \text{lowpt}) = 12$ efficiency bins are used at net-charge fluctuation analysis.

NBD check

100

$$C_1 = \mu$$

$$C_2 = \mu\epsilon$$

$$C_3 = \mu\epsilon(2\epsilon - 1)$$

$$C_4 = \mu\epsilon(6\epsilon^2 - 6\epsilon + 1)$$

$$C_5 = \mu\epsilon(2\epsilon - 1)(12\epsilon^2 - 12\epsilon + 1)$$

$$C_6 = \mu\epsilon(120\epsilon^4 - 240\epsilon^3 + 150\epsilon^2 - 30\epsilon + 1)$$

For example : 350003, Refmult2=300,

Efficiency Uncorrected case

$$C_1^+ = 79.53$$

$$C_1^- = 77.81$$

$$C_2^+ = 122.1$$

$$C_2^- = 117.93$$

$$\varepsilon^+ = C_2/C_1 = 1.53$$

$$\varepsilon^- = 1.51$$

used as NBD parameters.

$$C_3^+ = 253$$

$$C_3^- = 240$$

$$C_4^+ = 724$$

$$C_4^- = 671$$

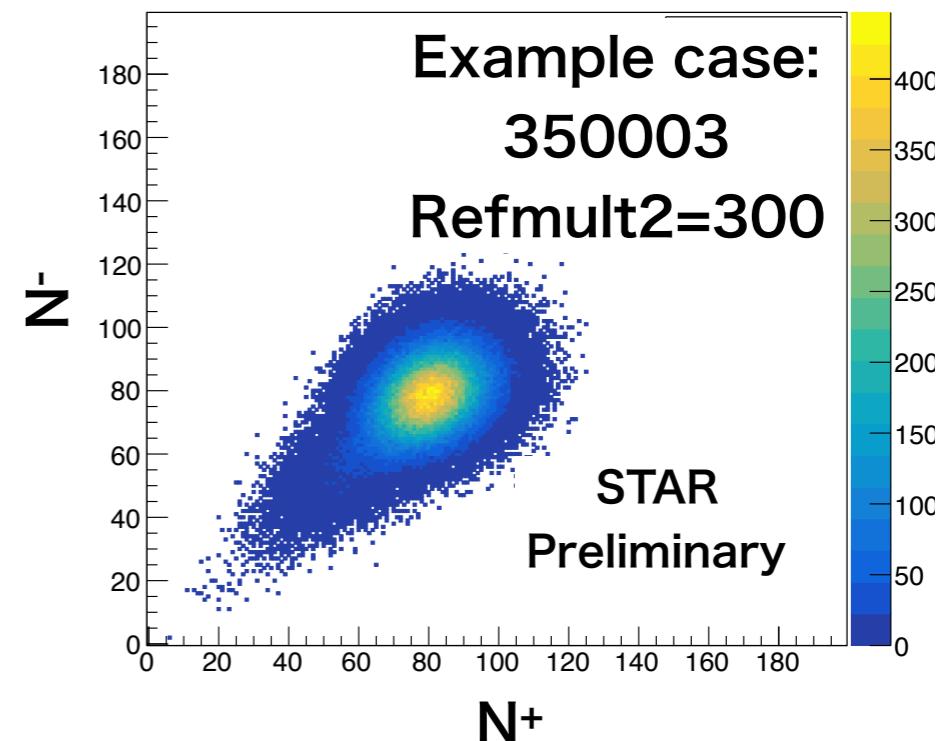
$$C_5^+ = 2746$$

$$C_5^- = 2486$$

$$C_6^+ = 13027$$

$$C_6^- = 11525$$

$$C_n^{net-c} = C_n^+ + (-1)^n C_n^-$$



NBD

$$C_1^{net} = 1.72$$

$$C_2^{net} = 240$$

$$C_3^{net} = 13.3$$

$$C_4^{net} = 1395$$

$$C_5^{net} = 260$$

$$C_6^{net} = 24552$$

Poisson

$$C_1^{net} = 1.72$$

$$C_2^{net} = 157$$

$$C_3^{net} = 1.72$$

$$C_4^{net} = 157$$

$$C_5^{net} = 1.72$$

$$C_6^{net} = 157$$

Odd: $C_1^+ - C_1^-$
Even: $C_1^+ + C_1^-$

- NBD baseline is larger than Poisson because of large ε at 200GeV.
- If $\varepsilon = 1$, NBD and Poisson baseline are exactly the same.

NBD check

101

For example : 350003, Refmult2=300,

Efficiency Corrected case

$C_1^+ = 139.44$	$C_1^- = 136.54$
$C_2^+ = 265.5$	$C_2^- = 256.26$
$\varepsilon^+ = C_2/C_1 = 1.90$	$\varepsilon^- = 1.87$

used as NBD parameters.

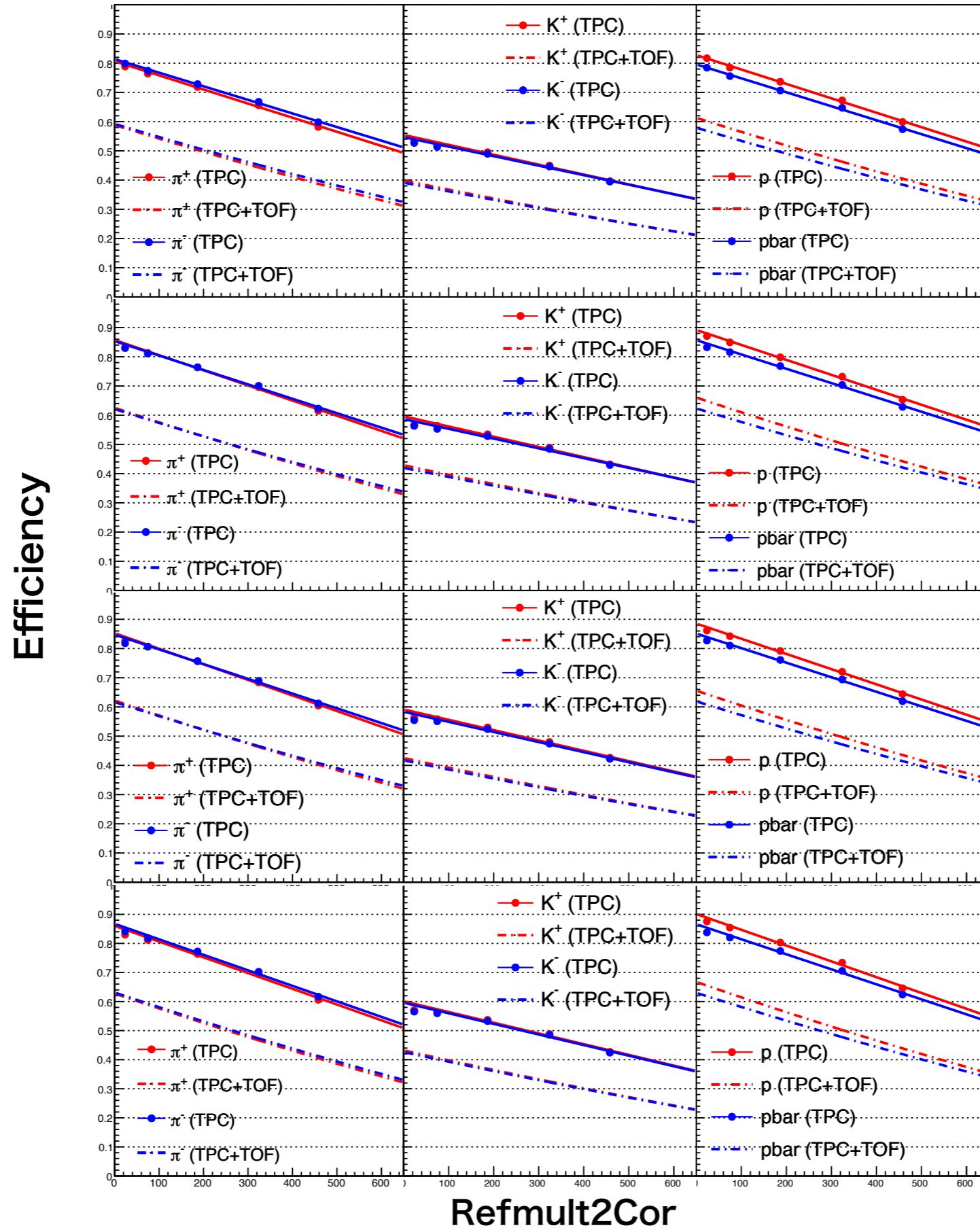
$C_3^+ = 745$	$C_3^- = 705$
$C_4^+ = 3007$	$C_4^- = 2786$
$C_5^+ = 16145$	$C_5^- = 14640$
$C_6^+ = 108378$	$C_6^- = 96183$

	NBD	Poisson	N ⁺
$C_1^{\text{net}} = 1.72$	$C_1^{\text{net}} = 1.72$		
$C_2^{\text{net}} = 521$	$C_2^{\text{net}} = 157$		
$C_3^{\text{net}} = 40.0$	$C_3^{\text{net}} = 1.72$		
$C_4^{\text{net}} = 5794$	$C_4^{\text{net}} = 157$		
$C_5^{\text{net}} = 1505$	$C_5^{\text{net}} = 1.72$		
$C_6^{\text{net}} = 204561$	$C_6^{\text{net}} = 157$		

Odd: $C_1^+ - C_1^-$
Even: $C_1^+ + C_1^-$

- NBD baseline is larger than Poisson because of large ε .
- If $\varepsilon = 1$, NBD and Poisson baseline are exactly the same.

Tracking efficiency



Run10

350003

350013

350043